

Topological Luttinger-semimetal CoAs₃ dye-sensitized photocatalyst for efficient solar hydrogen evolution

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Solar-driven photocatalytic water splitting is a process for hydrogen production from a renewable source. The practical implementation of this technology is limited by the low conversion efficiency of the hydrogen evolution reaction under visible light and the insufficient long-term stability of photocatalysts. Here we demonstrate a dye (Eosin Y)-sensitized photocatalyst for efficient hydrogen production. The topological semimetal CoAs₃ achieves a hydrogen production rate of 2688 $\mu\text{mol h}^{-1} \text{g}^{-1}$ ($\lambda \geq 420 \text{ nm}$) and an apparent quantum efficiency of 15.2% at $\lambda = 500 \text{ nm}$. Efficient photocatalytic activity is attributed to the electronic properties of CoAs₃, which facilitate electron transfer at the Eosin Y/CoAs₃ interface, determined by transient absorption spectroscopy. Density functional theory calculations predict that CoAs₃ is a Luttinger semimetal, exhibiting a quadratic band touching point near the Fermi level and an associated topological insulator gap. The carrier mobility of the material facilitates the transfer of injected electrons from the dye to active sites. Herein, we report a topological photocatalyst that exhibits enhanced stability and efficiency for solar hydrogen production.

The amount of solar energy incident on the Earth in a single hour exceeds the total annual energy consumption of the global population. Consequently, solar energy represents a viable alternative to fossil fuels^{1,2}. In order to facilitate the wider implementation of solar energy, it is essential to develop a framework for the sustainable and environmentally friendly storage of chemical fuels. Hydrogen is a candidate renewable fuel; it is an energy-dense carrier that produces only water upon combustion with oxygen^{3–6}.

The electrocatalytic hydrogen evolution reaction (HER) requires an external electrical input, the sustainability of which

depends on its primary generation source. On the other hand, dye-sensitized photocatalysis (DSP) has been proven to be a promising method for directly converting solar energy into H₂ without using electric power. In our specific system for photocatalytic HER, tetrabromofluorescein (EY) is employed as the photosensitizer to absorb visible light^{7,8}. For a standard DSP process, photoexcitation of the sensitizer is followed by electron-hole separation and transfer to the catalysts and sacrificial agents to initiate reductive and oxidative reactions, respectively^{9,10}. The charge separation and transfer, which are crucial for the performance of DSP systems,

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require suitable catalysts with fast charge transfer kinetics, both internally and interfacially^{11–15}.

Until now, noble metals (and noble-metal-containing materials), such as Pt, Ir, and Pd, as well as the state-of-the-art HER catalysts showing high electrocatalytic HER performance^{16–20}, have been well studied, and some well-known concepts such as *d*-band theory have been proposed^{21–25}. While noble metals exhibit favorable electron conductivity, their scarcity and cost limit large-scale application^{26–28}. Consequently, a central challenge is the development of non-noble metal catalysts that exhibit comparable activity. Various low-cost catalysts based on earth-abundant materials have been explored for HER, such as transition metal chalcogenides^{29,30}, transition metal phosphides^{31,32}, transition metal carbides^{33,34}, transition metal nitrides^{35,36}, and so on. Non-noble metal photocatalysts with high carrier mobility and density are worth investigating for photocatalytic hydrogen production under visible light.

The possibilities for developing high-performance catalysts have been expanded after the discovery of topological materials^{37,38}. Topological materials, which include topological insulators (TIs), topological semimetals (e.g., Dirac, Weyl, and nodal-line semimetals), and topological superconductors, are characterized by electronic boundary states protected by their band topology³⁹. The protected surface states in topological materials host mobile charge carriers, a property expected to promote efficient charge transport and long-term stability, ultimately enhancing photocatalytic activity^{37–40}. Several topological semimetals like (Nb,Ta)/(As,P) and *M*IrTe₄ (*M* = Nb,Ta) have been investigated for dye-sensitized HER, and the hydrogen production rate of NbIrTe₄ was reported to be 18,000 $\mu\text{mol g}^{-1} (10\text{ h})^{-1}$ ⁴¹. However, the range of topological materials applied to photocatalytic systems is still limited. Moreover, the relationship between topological properties and the photocatalytic HER reaction process requires further investigation^{41,42}. Further progress in photocatalytic HER is limited by two key challenges: inefficient utilization of solar light and the recombination of charge carriers. These challenges often cause low solar-to-hydrogen conversion efficiency and necessitate the exploration of photocatalysis with earth-abundant elements.

Our work aims to design an effective photocatalyst by exploring the properties of topological materials. The transition metal pnictide CoAs₃ is one of the promising candidates, as its predicted topological electronic structure is expected to yield favorable charge transport properties. In this work, we prepared the topological semimetal CoAs₃ for dye-sensitized HER, as well as other topological materials, including the semimetals WTe₂, HfP₂, and the insulators Mo₂C, Bi₂Te₃, and Bi₂Se₃. We evaluated the photocatalytic activity of these materials, found that the noble-metal-free photocatalyst CoAs₃ achieves a hydrogen production rate of 2688 $\mu\text{mol h}^{-1} \text{g}^{-1}$ under visible light irradiation ($\lambda \geq 420\text{ nm}$). Furthermore, transient absorption (TA) spectroscopy and density functional theory (DFT) calculations indicate that this high catalytic activity can be attributed to effective electron transfer at the dye/CoAs₃ interface and the material's topological semimetal nature. Our results suggest that the carrier dynamics in such materials can be utilized to address the existing challenges, providing a foundation for designing improved photocatalysts.

Results and discussion

Characterizations of CoAs₃ compounds

The cubic crystalline structure of the skutterudite-type compounds CoAs₃ in the *Im* $\bar{3}$ symmetry is shown in the insert image in Fig. 1a, where the Co (deep blue) and As (green) atoms occupy the 8c and 24g Wyckoff positions, respectively. The six-coordinated octahedron is formed by one cobalt atom and six arsenic atoms, with a Co-As bond length of 2.345 Å. The powder X-ray diffraction (XRD) pattern (Fig. 1a) shows sharp diffraction peaks, all of which can be indexed to the cubic skutterudite phase, indicating a high crystalline quality and phase-pure

sample. The diffraction patterns were indexed to the cubic space group *Im* $\bar{3}$ (Fig. 1a and Table S2). Rietveld refinement of the XRD data for CoAs₃ yielded lattice parameters consistent with previously reported values⁴³. Scanning electron microscopy images (Fig. S13) show that the CoAs₃ crystals have an average size of 1.2 μm , with individual crystals ranging from 0.5 to 3 μm . Energy dispersive X-ray (EDX) spectroscopy mapping of a CoAs₃ microparticle (Fig. 1c–e) reveals a spatially uniform distribution of Co and As, with no other elements detected. The atomic structure of the CoAs₃ microcrystals was further investigated using high-resolution scanning transmission electron microscopy (HR-STEM) and selected-area electron diffraction (SAED). The HR-STEM images (Fig. 1h, i) and SAED pattern (Fig. 1g) confirm the crystal structure, and the measured lattice constants are consistent with reported values⁴⁴. X-ray photoelectron spectroscopy (XPS) was used to analyze the chemical states of the CoAs₃ polycrystals (Fig. 1j, k). The Co 2p spectrum shows two main peaks at 779.3 eV (2p_{3/2}) and 794.3 eV (2p_{1/2}), which correspond to a mixed valence state of Co²⁺ and Co³⁺. The As 3d spectrum exhibits peaks at 42.2 and 41.4 eV, corresponding to Co-As bonds. These binding energies are all consistent with previously reported values^{45,46}.

Photocatalytic activity of CoAs₃

To identify the topological properties beneficial for photocatalytic HER, we compared the dye-sensitized HER performance of CoAs₃ with that of other topological materials, including WTe₂, HfP₂, Mo₂C, Bi₂Te₃, and Bi₂Se₃. These materials were selected to represent a range of topological properties and band structures near the Fermi level (*E*_F) and were tested under identical experimental conditions. EY and triethanolamine (TEOA) were used as the photosensitizer and the sacrificial electron donor, respectively. The photocatalytic activity of these materials remains stable over 5 h (Fig. 2a); such stability may be attributed to their protected topological surface states^{47,48}. Among these materials, CoAs₃ achieves a photocatalytic HER rate of 684 $\mu\text{mol h}^{-1} \text{g}^{-1}$, 14.3 times that of Bi₂Se₃. The electrochemically active surface area (ECSA) of each photocatalyst was determined from cyclic voltammetry measurements (Supplementary Figs. S1–S3). Although the ECSA of CoAs₃ ($1.61 \times 10^{-4} \text{ mF cm}^{-2}$) is approximately half that of Mo₂C, its photocatalytic activity exceeds that of other materials tested. This discrepancy points to the contribution of the electronic property of CoAs₃ to its catalytic activity, which we investigate in the following sections.

The photocatalytic activity of CoAs₃ was investigated by varying the catalyst mass and photosensitizer concentration. To determine the effect of catalyst mass, the concentrations of the EY and the TEOA were kept constant at 10 mg and 15 mL, respectively. Figure S4 shows that H₂ can be produced in the absence of CoAs₃ catalysts, and the activity increased with the addition of CoAs₃. Therefore, the hydrogen evolution rate attributed to CoAs₃ can be calculated by subtracting the rate of the control experiment from the total measured rate. As shown in Fig. 2b, the normalized activity of CoAs₃ remains stable between 1 and 25 mg, which well reflects the intrinsic activity of CoAs₃. When the catalyst mass was increased beyond 30 mg, the hydrogen evolution rate decreased. This rate decrease is attributed to the large amount of catalyst masking the active sites and blocking the light absorption⁴⁹. Therefore, in subsequent experiments, the amount of CoAs₃ was fixed at 25 mg.

The dosage of photosensitizer was then adjusted from 10 to 80 mg, while keeping the catalyst and sacrificial agent constant (Fig. 2c). The hydrogen evolution rate initially increased as the amount of photosensitizer was raised. The rate reached a maximum value of 2697 $\mu\text{mol h}^{-1} \text{g}^{-1}$ at a photosensitizer amount of 70 mg, and subsequently decreased when the amount was increased to 80 mg (Fig. 2c). This decrease at higher concentrations is attributed to two factors: (1) a shading effect caused by excess photosensitizer, and (2) collisional quenching of the excited photosensitizer (*EY) by ground-state EY

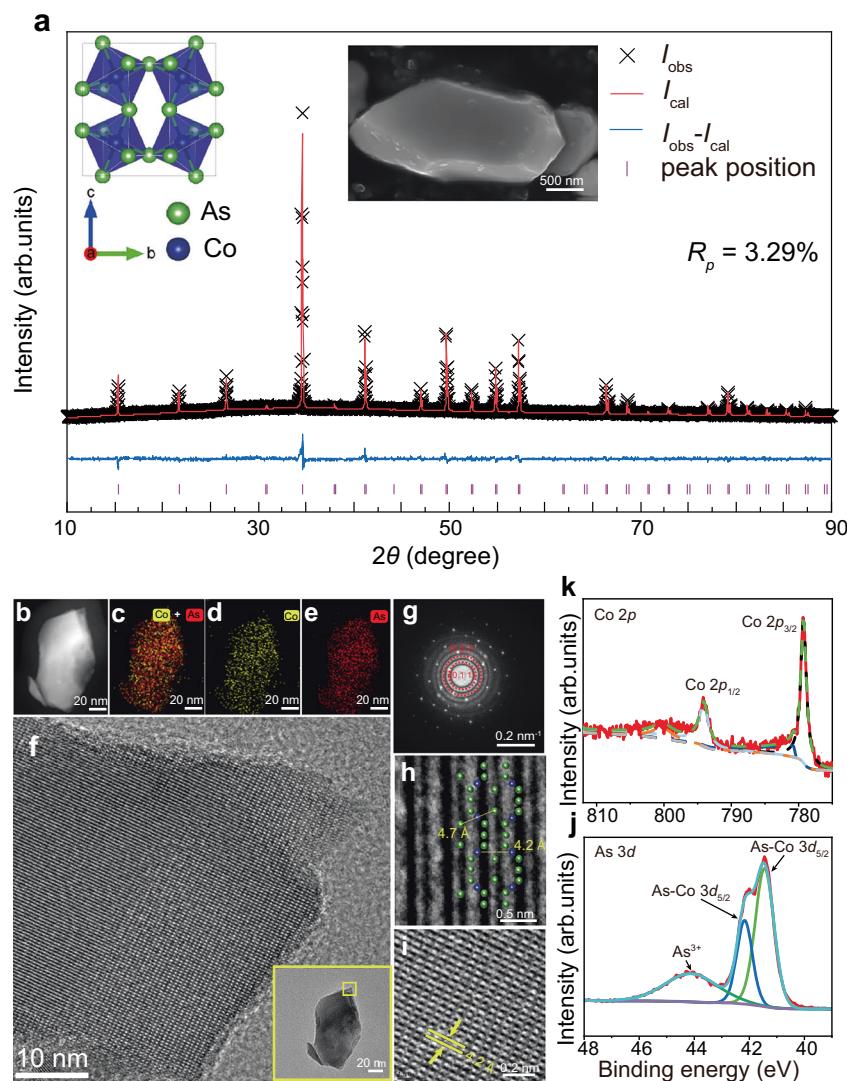


Fig. 1 | Crystalline structure, X-ray diffraction, SEM images, element mappings, and XPS characterization of CoAs₃ microcrystal. **a** The final Rietveld refinement of XRD spectrum of CoAs₃ polycrystals. (left insert) Structural model of CoAs₃ (space group *Im*3). Blue and green balls represent Co and As atoms, respectively. (right insert) SEM image of CoAs₃ microcrystal. **b–e** SEM images (**b**) and element mappings of CoAs₃ flake, showing a uniform distribution of Co + As (**c**), Co (**d**), and As (**e**) elements. The average Co:As atomic ratio determined by the energy-dispersive X-ray spectroscopy (EDX) spectra is close to 1:3, and no foreign elements were detected within the detection limit of the instrument; **f–i** Atomic-level TEM

image and low-magnification scanning transmission electron microscopy (STEM) image. (insert) of CoAs₃ microcrystals. The yellow rectangle at the top right of the insert graph indicates the area where the atomic-level TEM (**f**) and the selected-area electron diffraction (SAED) pattern (**g**) were measured. Red circles in (**g**) mark the primary diffraction spots. **h, i** Atomic-level TEM images of a CoAs₃ crystal taken at different locations, where 4.2 Å corresponds to the distance to the nearest neighboring Co. **j, k** X-ray photoemission spectroscopy (XPS) spectra of **j** Co 2p and **k** As 3d for freshly prepared CoAs₃ polycrystals. Source data for Fig. 1 are provided as a Source data file.

molecules, which can deactivate *EY before interfacial electron transfer occurs^{11,50,51}.

Control experiments show that hydrogen evolution is negligible in the absence of any of the three components: the CoAs₃ catalyst, the EY photosensitizer, or the TEOA sacrificial agent (Fig. 2d). Under full-spectrum irradiation, the hydrogen evolution rate increased to 3232 μmol h⁻¹ g⁻¹. In contrast, graphene with super electron transfer efficiency was also used for comparison, which exhibits a lower activity compared with CoAs₃. These results suggest that CoAs₃ not only serves as an effective charge transfer material but also provides abundant HER active sites in this system. It is important to note that the rate of photocatalysis is likely to vary substantially depending on the effects of a given temperature, pressure, light intensity, and other factors. The ultraviolet–visible absorption spectra of CoAs₃ and Bi₂Se₃ exhibit minimal absorption in the visible range (Fig. S12). This property suggests they do not function as the primary light absorbers in this

photosensitized system. The absorption peaks of pure EY solution are at 480 and 518 nm due to π–π* excitation⁴¹, which is consistent with previous results⁵². When EY powder is mixed with CoAs₃ and Bi₂Se₃ powders, respectively, absorption peaks of 487 and 526 nm appear near the characteristic peaks of the dye, indicating a redshift. This redshift suggests an electronic interaction between EY and the CoAs₃ or Bi₂Se₃ materials. Such an interaction may facilitate the transfer of photogenerated electrons from the excited photosensitizer to the catalyst⁵³. Additionally, a color change in the solution from red–orange to saffron was observed during the photocatalytic process, suggesting degradation of the EY photosensitizer over the reaction time.

The apparent quantum efficiency (AQE) was measured at several incident light wavelengths (Fig. 2e) (See the “Methods” section for detailed experimental and calculation procedures.). The CoAs₃-EY system exhibited a maximum AQE of 15.2% at an incident wavelength of 500 nm. Relatively high AQE values can also be obtained at 450 and

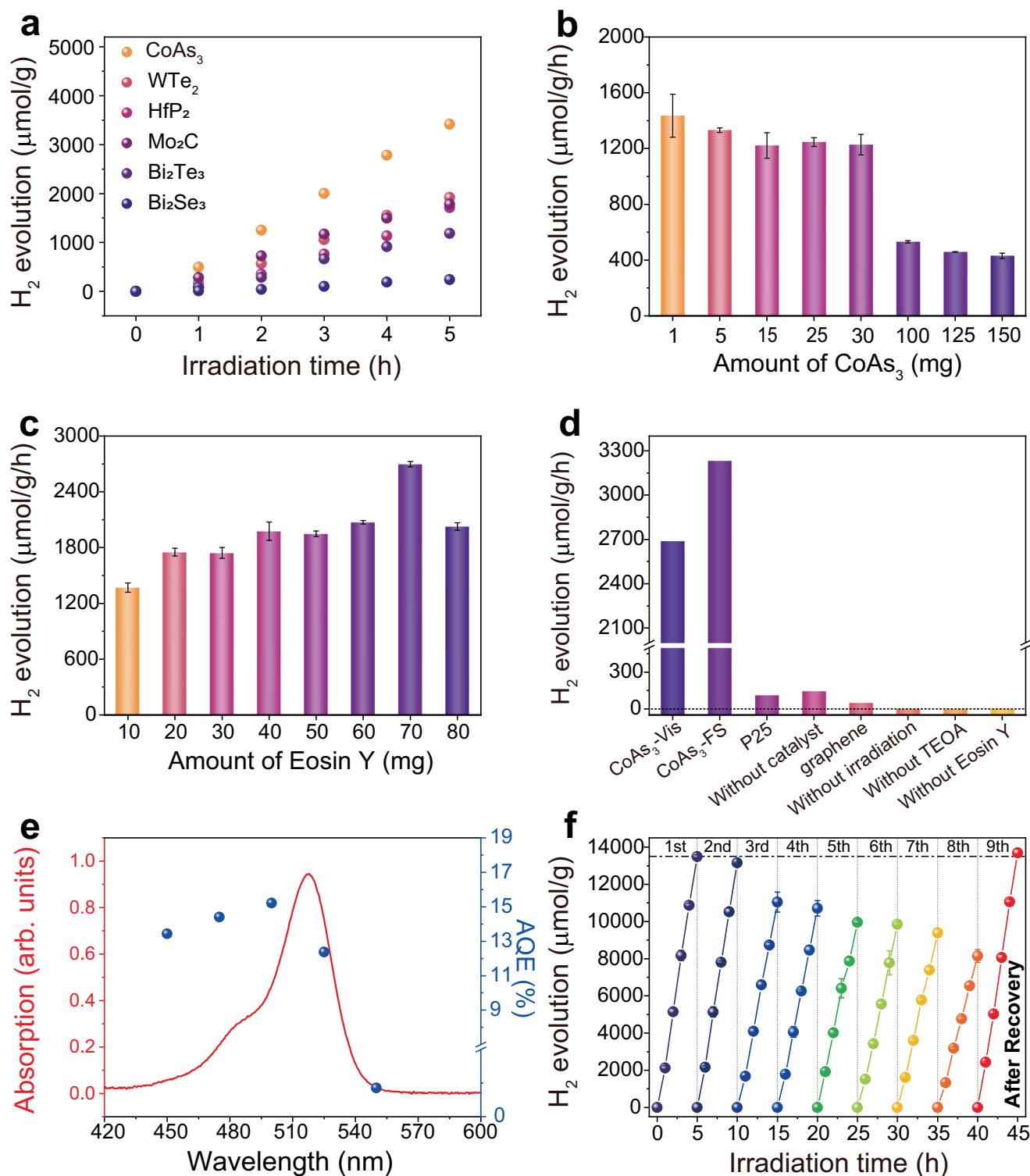


Fig. 2 | Photocatalytic activity of various topological materials. **a** Hydrogen production rates of CoAs_3 , WTe_2 , HfP_2 , Mo_2C , Bi_2Te_3 and Bi_2Se_3 . **b** Eosin Y dye-sensitized photocatalytic HER activity of CoAs_3 after 5 h of irradiation, normalized with respect to the catalyst amount in grams, with different amounts of catalyst. **c** Hydrogen production rate of CoAs_3 at different amounts of dye. **d** Control experiments of hydrogen production in the dye sensitization system. P25 represents for TiO_2 nanoparticles photocatalyst of ca. 25 nm. CoAs_3 -Vis represents

photocatalysis under visible light, while CoAs_3 -FS refers to photocatalysis under a full-spectrum light. **e** The variation of the quantum apparent efficiency (AQE) of photocatalyst CoAs_3 vs. radiation wavelength (blue dots) and the UV-vis absorption spectrum of EY (red curve). **f** The reproducibility of the EY dye-sensitized photocatalytic hydrogen evolution reaction (HER) activity of CoAs_3 for the optimum dye/catalyst ratio. In (**b**, **c**, **f**), data are presented as mean \pm s.d. from three independent experiments ($n=3$). Source data for Fig. 2 are provided as a Source data file.

480 nm with lower EY absorption intensity. The system remains active at shorter wavelengths despite the reduced light absorption by the photosensitizer.

To investigate the stability of the CoAs₃ photocatalyst at the optimal ratio of 25 mg CoAs₃ and 70 mg photosensitizer, a set of cycling experiments with a unit cycle of 5 h was conducted. Over eight consecutive 5-h cycles, the hydrogen evolution rate showed a minor decrease in each cycle (Fig. 2f). In the recovery experiment (ninth cycle), the catalysts were cleaned and centrifuged before being reused in a fresh experimental solution. Upon reuse, the catalyst's activity was restored to its initial level, yielding 13,692 $\mu\text{mol g}^{-1}$ of H₂ in 5 h. Due to the inferior stability of EY⁵⁴, the decreased activity of the DSP system may be ascribed to the reduced concentration of EY. Therefore, after the recovery of EY concentration, the activity can be well-restored, indicating a stable performance of CoAs₃ catalyst. The EDX and XPS analyses (Supplementary Fig. S5) of the catalysts performed after the cycling experiments showed no observable changes in their components or morphology. Collectively, these findings indicate that CoAs₃ is a structurally and functionally stable catalyst for photocatalytic hydrogen evolution.

Electron transfer efficiency at the EY/CoAs₃ interface

To investigate the electron-transfer process during DSP photocatalytic HER, amperometric current density *I*-*t* curves were measured (Supplementary Fig. S6) in a 0.1 M Na₂SO₄ and TEOA aqueous solution. In the absence of a catalyst, the EY-sensitized electrode exhibited a transient anodic photocurrent spike upon illumination. This feature is characteristic of charge accumulation at the electrode-sensitizer interface followed by rapid recombination, as the photogenerated electrons are not transferred to an acceptor^{55–57}. This behavior indicates fast charge recombination when no effective electron acceptor presents^{58,59}.

In contrast, the introduction of CoAs₃ catalyst changes the carrier dynamics. For the EY/CoAs₃ system, the photocurrent increases throughout the illumination period. This photoactivation is attributed to light-induced effects that occurred at the dye-catalyst interface, such as a stable dye adsorption and a dynamic adsorption-desorption equilibrium, which enhance photosensitization efficiency⁶⁰. Furthermore, the slow current decay after illumination ceases indicates that CoAs₃ possesses charge-storage capacity, allowing trapped electrons to subsequently participate in the HER⁵⁸. Notably, while the EY/Bi₂Se₃ system shows a similar photoactivation trend, its total photocurrent density is much lower. This lower photocurrent is observed despite Bi₂Se₃ being a TI with surface states reported to support efficient charge transport. Therefore, this result strongly suggests that the rate-determining step is not charge transport but the slow kinetics of the surface HER. These comparative results indicate that the overall photoelectrochemical efficiency depends on a combination of interfacial charge separation, bulk transport, and surface reaction kinetics. The high performance of the EY/CoAs₃ system should be due to a better balance of many factors, particularly the higher intrinsic HER activity and favorable photoactivation properties of the CoAs₃ surface.

The kinetics of interfacial charge transfer were investigated using electrochemical impedance spectroscopy. The charge transfer resistance (*R*_{ct}) of the EY-sensitized CoAs₃ and Bi₂Se₃ systems was compared. Both the EY/CoAs₃ and EY/Bi₂Se₃ systems exhibited their lowest *R*_{ct} at 520 nm, as shown in Fig. S11. This wavelength corresponds to the primary absorption peak of the EY photosensitizer, suggesting that the rate of electron injection from the excited dye to the catalyst is dependent on the amount of light absorbed. Moreover, under all comparable illumination wavelengths, the EY/CoAs₃ system consistently displayed a smaller *R*_{ct} than the EY/Bi₂Se₃ system (specific values summarized in Table S5). This lower *R*_{ct} value for the CoAs₃ system indicates more rapid interfacial electron transfer kinetics

compared to the Bi₂Se₃ system. This process includes both electron injection from the dye to the catalyst and the subsequent utilization of charge in the HER.

The interfacial charge transfer mechanism between photoexcited EY and the catalysts was investigated using in situ electron paramagnetic resonance (EPR) spectroscopy measured under illumination. In contrast to the weak intrinsic signal observed for pure EY under illumination, attributed to the formation of the excited state, the EY/catalyst mixtures exhibited a dramatic increase in signal intensity at *g* = 2.001 upon illumination (Fig. S15). This prominent *g*-value is characteristic of an organic radical, consistent with the formation of the EY radical cation (EY^{•+}). Quantitatively, the EPR signal intensity for the EY/Bi₂Se₃ and EY/CoAs₃ mixtures was enhanced by factors of 6.1 and 26, respectively, relative to that of pure EY under the same illumination. This signal enhancement provides the spectroscopic evidence for efficient electron transfer from photoexcited EY to the catalyst, leading to the generation of the EY^{•+} radical cation. The larger signal enhancement factor for CoAs₃ (26) compared to that for Bi₂Se₃ (6.1) indicates that CoAs₃ more readily accepts an electron from photoexcited EY, leading to more effective charge separation.

Femtosecond TA spectroscopy was used to investigate the ultrafast charge transfer dynamics between the EY photosensitizer and the prepared catalysts with an excitation wavelength of 570 nm. By systematically investigating the ultrafast carrier dynamic processes of pure EY solution and the EY/Cat composites, we have obtained direct spectroscopic evidence of photoinduced electron transfer (PET) and resolved the subsequent multi-step kinetic pathways. The two-dimensional pseudo-color maps for pristine EY/CoAs₃ and EY composite are presented in Fig. 3a, d, respectively, with the data for EY/Bi₂Se₃ provided in the Supporting Information (Fig. S7a).

To establish a benchmark, the transient behavior of pristine EY was characterized to understand its intrinsic photophysical processes. The TA spectrum of pristine EY (Fig. 3e) exhibits a positive photo-induced absorption (PIA) band centered at 420 nm and a negative bleach signal centered at 510 nm. According to the previous report, the PIA signal is assigned to the absorption of singlet ¹EY, while the bleach signal comprises the broad ground-state bleach and stimulated emission of EY^{61,62}. In stark contrast, the introduction of the catalysts in EY/Cat composites leads to an obvious change in their TA spectra. Besides the PIA band centered at 420 nm (for ¹EY) and the negative bleach signal at 510 nm present in pure EY, both the TA spectra of EY/CoAs₃ and EY/Bi₂Se₃ (Figs. 3b and S7a) still show an additional PIA signal with its peak at 470 nm, which is the definitive signature of the EY radical cation (EY^{•+}). This emergence of EY^{•+} provides direct spectroscopic evidence for electron transfer from the excited EY to the catalyst, which can also be clearly reflected in their respective TA kinetics.

In pristine EY, the similar TA kinetics probed, respectively, at 420 and 445 nm are observed (Fig. 3f), while in EY/Cat composites, their TA kinetics probed at 420 and 470 nm exhibit completely different behaviors. The PIA signal at 420 nm shows a fast decay in the initial 2 ps delay, accompanied by a clear rising process of PIA signal at 470 nm in the same time scale (Figs. 3c and S7c), again verifying the PET from EY to the catalyst. Fig. 3g compares the PIA kinetics probed at 420 nm in pure EY and the EY/Cat composites. Compared with pure EY, both the excited states in the two EY/Cat composites exhibit evidently faster decay kinetics, which can be well fitted by a tri-exponential function (Table 1). Obviously, the fastest process (with τ_1 = 2.1 ps and 3.2 ps for EY/CoAs₃ and EY/Bi₂Se₃, respectively) is attributable to the PET process, also consistent with the rise time observed in the PIA kinetics at 470 nm, while the slowest time constant exactly represents the lifetime of the charge-separated state. We also note that the faster process with time constant (τ_2) about tens of ps, likely originates from the accelerated decay of EY excited states in EY/Cat composites⁶¹. The above kinetics comparison between EY/CoAs₃ and EY/Bi₂Se₃ suggests the faster PET rate and longer lifetime of charge-separated state at the EY/

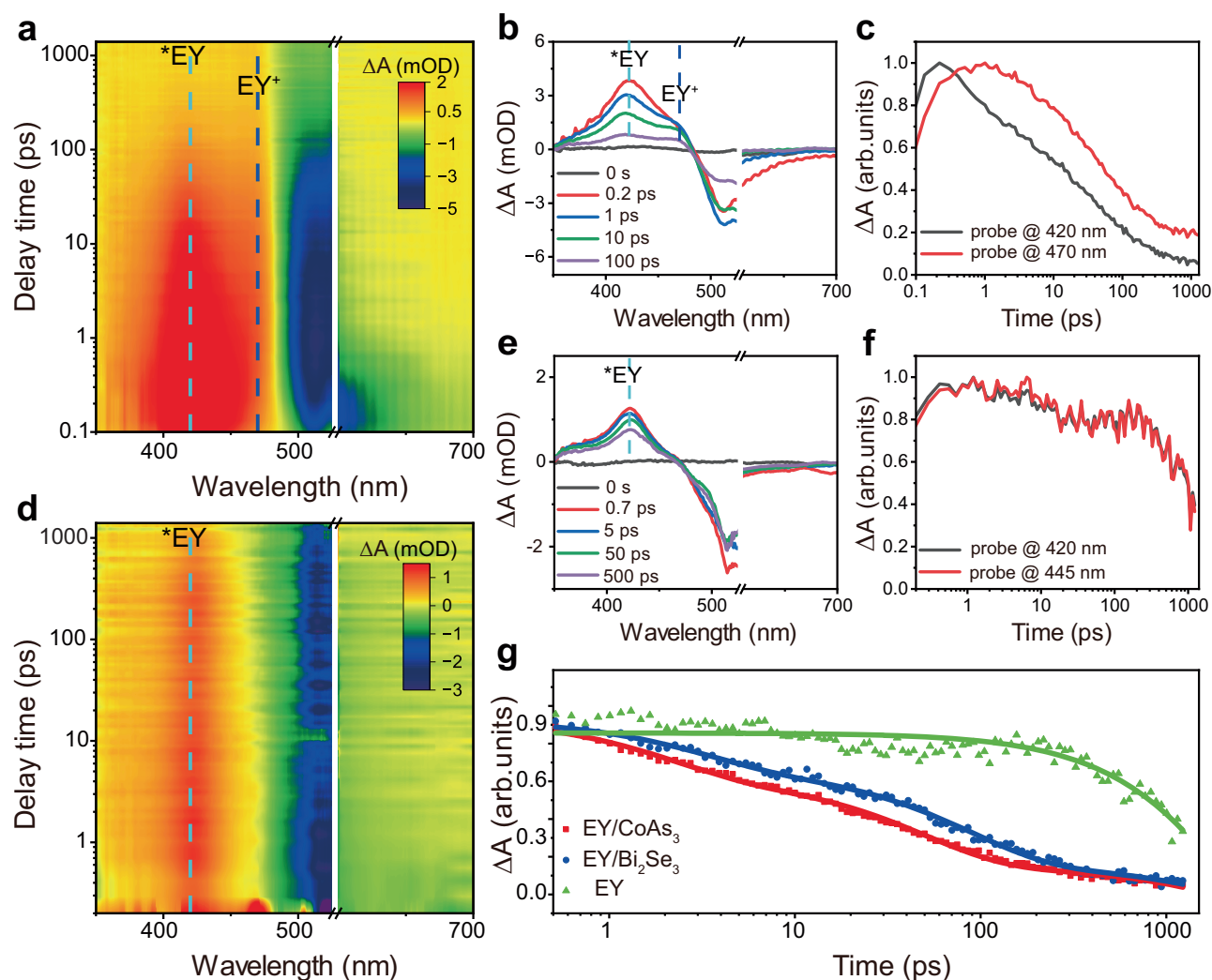


Fig. 3 | Ultrafast carrier dynamics at the EY/catalyst interface. Time slices of the TA spectra of **a**, **b** EY/CoAs₃ composite and **d**, **e** pure EY solution under excitation at 570 nm. Comparison of the TA kinetic probed at different wavelengths in EY/CoAs₃

(c) and EY (f). **g** Comparison of the TA kinetic probed at 420 nm between in EY solution (green), EY/CoAs₃ (red), and EY/Bi₂Se₃ (blue). Solid lines are their tri-exponential decay fitting. Source data for Fig. 3 are provided as a Source data file.

Table 1 | Summary of kinetic parameters obtained from multi-exponential fitting of the transient absorption decays at 420 nm

	A1 (%)	T1 (ps)	A2 (%)	T2 (ps)	A3 (%)	T3 (ps)
EY/CoAs ₃	28.0	2.13	41.2	49.67	30.8	3245.22
EY/Bi ₂ Se ₃	24.5	3.24	46.9	86.12	28.6	2989.68
EY	>>1.5ns					

CoAs₃ interface over the Bi₂Se₃ interface, which can account for the higher photocatalytic hydrogen production activity of the CoAs₃ catalyst.

DFT-calculated electronic structure of CoAs₃

Following electron transfer from the excited photosensitizer, the subsequent transport of electrons to surface active sites is governed by the electronic structure of the catalyst material. We therefore investigate the electronic structure of CoAs₃ by DFT calculations and compare it with that of other photocatalysts investigated in the current work.

The concept of topological catalysis has recently attracted considerable interest^{38,63–67}. In contrast to TIs, the bulk states of

topological semimetals are conductive. This conductivity, along with elevated electronic mobility, is thought to be beneficial for catalytic processes^{37,38}. In addition, topological semimetals are found to carry a large number of topological emergent characters, including different categories of nodal points, nodal lines, and nodal surfaces^{65,68,69}. These topological emergent characters also show diverse surface states. Consequently, both theoretical and experimental studies have identified topological semimetals as promising catalysts. Their proposed catalytic advantages are attributed to a combination of properties, including bulk-state conductivity, topologically protected surface states, and favorable molecular adsorption energies^{6,41,70–79}. However, the mechanism responsible for the photocatalytic activity of CoAs₃ remains unclear, and the specific material properties that determine its performance in the HER have yet to be identified.

DFT calculations were performed to investigate the band structure of CoAs₃ (Fig. 4a), which provides insight into its photocatalytic behavior. The relaxed crystalline structure of CoAs₃ is provided as Supplementary Dataset 1. Without considering the spin–orbit coupling (SOC) effect, a threefold band crossing node appears at the high-symmetric Γ point, characterized by a topological charge $c=0$ (Fig. 4b). Upon incorporating SOC, this threefold degeneracy evolves into a sixfold band crossing node. This node further splits into a

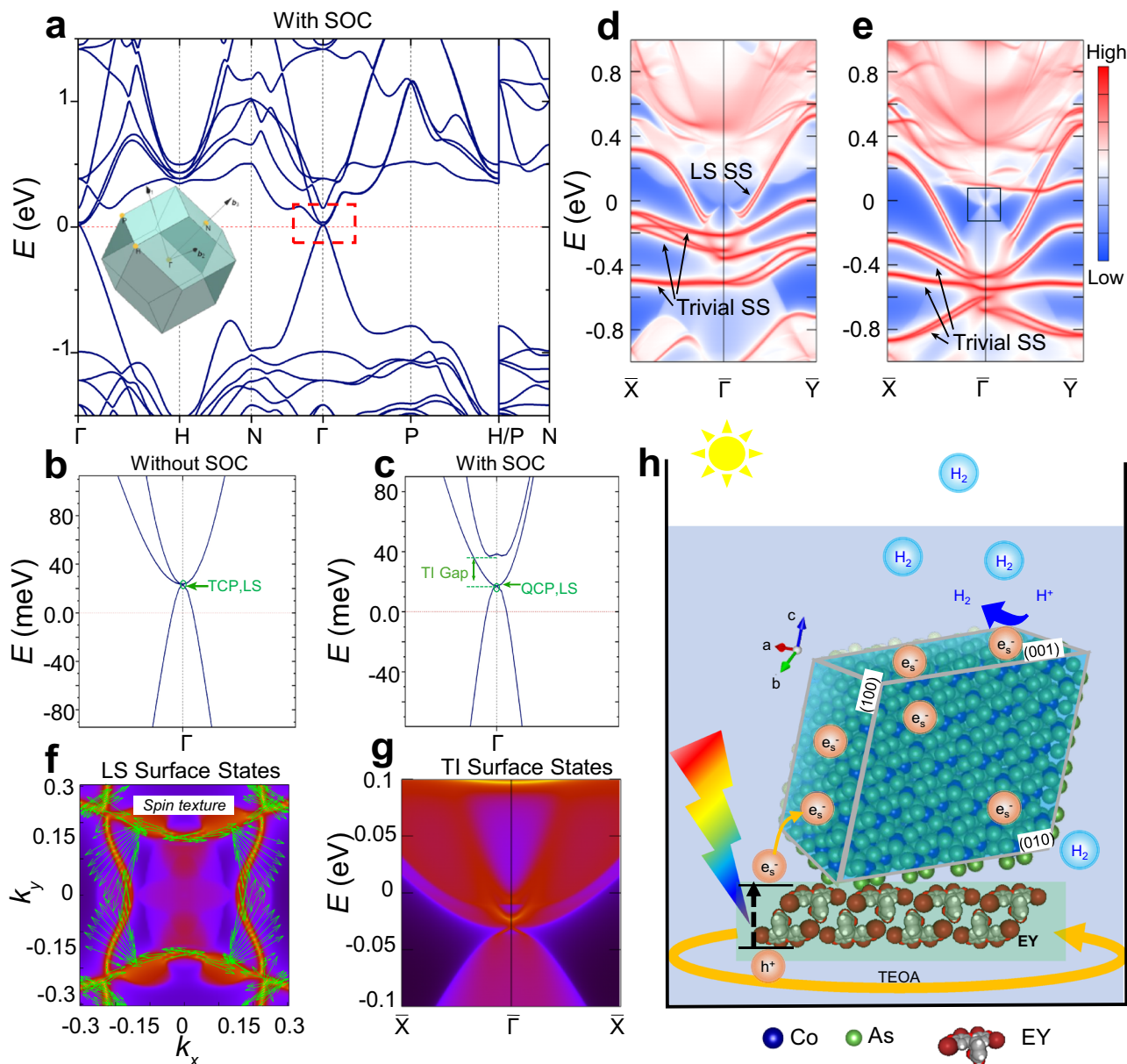


Fig. 4 | The calculated band structure and surface states of CoAs₃, and the proposed mechanism of enhanced photocatalytic HER performance.

a Electronic band structure of CoAs₃ considering the SOC. Insert: bulk Brillouin zones of CoAs₃. **b, c** The enlarged view of the dashed red area for the enlarged band crossing points with and without SOC along the high symmetry path N-Γ-P in (a). The quadratic contact point at Γ above E_F . (TCP triple contact-point, QCP quadratic contact-point, LS Luttinger semimetal). **d, e** Surface band structures of the top and

down surfaces. Note, the trivial surface states are induced by unsaturated dangling bonds on the surface, where the surface states of Luttinger semimetal are clearly shown. **f** Color mapping of the calculated spin texture of Luttinger semimetal (LS) surface states. **g** Color plot of the calculated topological insulator (TI) surface states. **h** The proposed mechanism of enhanced visible light photocatalytic HER performance of CoAs₃. Source data for Fig. 4 are provided as a Source data file.

twofold band crossing (triple contact-point (TCP)) and a fourfold degenerate quadratic contact-point (QCP), a phenomenon attributed to the coexistence of center-reversal symmetry and time-reversal symmetry. The presence of these degenerate points is indicative of the electronic structure of CoAs₃, potentially offering high carrier mobility and efficient electron pathways as well as efficient charge separation and transport in photocatalytic hydrogen evolution. In between the TCP and QCP, there is a TI gap and a Dirac-like topological surface state (Fig. 4c). The QCP along with specific surface states constitute a Luttinger semimetal (LSM)^{80–82}, where this degeneracy appears as a quadratic band touching of $j = 3/2$ fermions. LSMs have been studied due to their proximity to Z2 TIs⁸³, Dirac⁸⁴, and Weyl semimetals⁸⁵. This nodal point is formed by the crossing of the valence band maximum

and two conduction bands near the conduction band minimum and is protected by time-reversal symmetry and spatial-inversion symmetry. Detailed band structure analysis (Fig. 4b) indicates that the fourfold band degeneracy at the Γ point, near the Fermi level, primarily originates from the Co *d*-orbitals. Concurrently, with the inclusion of SOC, the sixfold nodal point transforms into a fourfold degenerate quadratic Dirac point (Fig. 4c)⁸⁴. To further investigate the nontrivial band topology of CoAs₃, its surface band structure was calculated (Fig. 4d, e). These calculations show both Luttinger surface states (LSS) crossing the E_F and the concomitant appearance of trivial surface states. The presence of these conductive surface states, particularly those crossing E_F , can serve as efficient channels for transferring photogenerated electrons to the surface active sites, thereby reducing charge

recombination and enhancing hydrogen evolution efficiency. Further calculations reveal the spin texture of the LSS (Fig. 4f), exhibiting a characteristic spin-momentum locked pattern. This spin-momentum locked pattern can suppress backscattering of electrons, thereby reducing recombination, a process beneficial for photocatalysis. A magnified view of the region near the Γ point at the Fermi level (E_F) shows that the TI surface states become apparent (Fig. 4g). Since both the LSS and the TI surface states are protected by chiral symmetry and time-reversal symmetry, respectively, they are insensitive to disorder or impurities that do not break these specific symmetries. This property helps to ensure that efficient and stable charge carrier transport channels persist even in imperfect or complex catalytic environments, thereby contributing to sustained activity and high overall efficiency. The surface states emanated from the bulk nodal points of CoP_3 have been observed by angle-resolved photoemission spectroscopy measurements⁸⁶ and are seen to cross through E_F . Compared to Bi_2Se_3 , CoAs_3 possesses both a fourfold degenerate quadratic Dirac point near E_F and the LSS crossing the E_F . These features contribute to a higher carrier density and high mobility in CoAs_3 , whereas Bi_2Se_3 has only conventional topological surface states crossing E_F . The band structure calculations (Figs. 4 and S8) indicate that the d -orbitals of Co are the primary contributors to the electronic states near the E_F , leading to the formation of a topological Dirac cone. These topological features could enhance the efficiency of catalytic reactions such as the photocatalytic HER.

The intrinsic nature of high dye-sensitized activity of HER using CoAs_3 catalyst can be rationally explained in Fig. 4h. Upon light absorption, the EY photosensitizer generates an electron-hole pair. The electron is injected into the CoAs_3 catalyst, leaving a hole on the EY molecule. The conductive topological surface states of CoAs_3 facilitate the transport of these injected electrons away from the interface, thus suppressing charge recombination. This process contrasts with that in conventional semiconductors, where charge carriers are often annihilated by defects. In CoAs_3 , the topological surface states allow electrons to bypass such defects, which helps to preserve the population of photo-generated electrons. Then, the surviving surface electrons are captured by protons adsorbed on the surface and generate H_2 . Furthermore, the Co d -band crosses the E_F (Supplementary Fig. S7). These states are closely related to hydrogen interaction and are proposed to act as the active sites for HER. Meanwhile, the sacrificial electron donor (TEOA) contributes electrons to the photosensitizer and maintains the surface redox reactions.

In summary, this work introduces the topological LSM CoAs_3 as a catalyst for dye-sensitized photocatalytic hydrogen evolution. A comparison with other topological materials was performed to evaluate its performance. Under visible light ($\lambda \geq 420$ nm), the CoAs_3 -based DSP system achieved a hydrogen production rate of $2688 \mu\text{mol h}^{-1} \text{g}^{-1}$. Ultrafast spectroscopy confirmed rapid electron transfer at the EY/ CoAs_3 interface, and theoretical calculations pointed to the role of the material's topological semimetal nature in its photocatalytic activity. This work identifies a topological semimetal as a stable and active catalyst for photocatalytic HER under visible light. The activity of the system under broad-spectrum illumination suggests a potential for hydrogen production using natural sunlight.

Methods

Reagents

Cobalt (Sigma Aldrich, 99.99%), Arsenic (Sigma Aldrich, 99.99%), Tungsten (Alfa Aesar, 99.95%), Tellurium (Alfa Aesar, 99.999%), molybdenum (Alfa Aesar, 99.95%), carbon (Alfa Aesar, 99.9%), Bismuth selenide (Sigma Aldrich, 99.99%) Bismuth(III) telluride (Sigma Aldrich, 99.99%), Eosin Y (Sigma Aldrich, 99%), TEOA (ACROS organics, 99+%), TiO_2 (Degussa P25).

Sample preparation

CoAs_3 : The sample were prepared via solid-phase reaction. Co and As were sealed in quartz tubes in a stoichiometric ratio, heated to 800°C at a rate of 100°C per hour, held for 48 h, and cooled down with the furnace.

WTe_2 : W and Te were ground stoichiometrically in an argon-filled glove box and sealed in an ampoule under vacuum, heated to 700°C for 7 h and held for 1 day, then heated to 750°C for 1 h and held for a period of time before being cooled in the furnace.

Mo_2C : Carbon and molybdenum were mixed and ground in stoichiometric ratio in an argon-filled glove box and pressed into tablets, flame sealed in ampoules under vacuum and cooled with the furnace after 10 h of heating to 950°C .

HfP_2 : The ingredients were combined and ground in stoichiometric proportions in an argon-filled glove box, then compressed into tablets. These were subsequently sealed under vacuum in ampoules and cooled with the furnace after 8 h of heating to 750°C .

Characterization

The element distribution and the chemical composition of the products were determined through EDS mapping and spot analysis (Hitachi S-4800 scanning electron microscope) and TEM (HITACHI HF 5000), respectively. The chemical composition and crystal lattice structure were determined by XRD (Bruker-D8) with $\text{CuK}\alpha$ radiation over an angular range of 10 – 90° in 1 h at room temperature. The chemical valence states were detected using an XPS (Kratos AXIS ULTRA XPS).

DSP HER

The photocatalytic reactions occurring in hydrogen were carried out in a sealed gas-solid atmospheric pressure cylindrical glass reactor. The light source was a 300 W xenon lamp with a 420 nm cut-off filter. In all, 100 mg of catalyst powder was dispersed in 100 ml of freshly prepared aqueous TEOA (15%v/v), and 0.014 mmol (9.07 mg) of EY was added to the glass reactor, followed by sonication for 30 min to ensure homogeneous particle dispersion ($\text{ph} = 10.15 \pm 0.1$). The vessel was positioned upon a magnetic stirrer, and the vacuum pump was activated to commence the photocatalytic process. This was followed by the illumination and agitation of the sample to initiate the reaction.

Apparent quantum efficiency

The AQE can be described based on the following formula:

$$\text{AQE} (\%) = \frac{A R N_A}{I} \times 100$$

where A , R , N_A , and I represent a coefficient (2 for H_2 evolution), the evolution rate of H_2 , Avogadro's constant, and the number of photons absorbed by the reaction solution, respectively. For evaluating the AQE, the reactions were performed under similar reaction conditions except that the monochromatic light was used. As the light source, the monochromatic light was generated by Xe lamp equipped with the band-pass filters centered at 380, 420, 450, 500, 550, and 600 nm, respectively.

TA spectroscopy

The femtosecond pump-probe TA measurements were performed using a regenerative amplified Ti:sapphire laser system (Coherent; 800 nm, 70 fs, 6 mJ pulse^{-1} , and 1 kHz repetition rate) as the laser source and a femto-TA100 spectrometer (Time-Tech Spectra). All experiments were performed under ambient conditions.

First-principles calculations

First-principles calculations were performed based on DFT by utilizing the Vienna ab initio simulation package^{87,88}. The plane-wave basis was adopted with an energy cutoff of 350 eV. The electron-ion interactions

were modeled with the projector augmented wave potential⁸⁹, and the exchange–correlation functional was approximated by the Perdew–Burke–Ernzerhof-type generalized gradient approximation⁹⁰. For optimized lattice constants and atomic positions, the structural relaxation was performed with an energy (force) criterion of 10^{-8} eV (0.01 eV \AA^{-1}) and by utilizing the DFT-D3⁹¹ method to include van der Waals corrections. Surface state calculations were performed with the Wannier Tools package⁹², which is based on the tight-binding Hamiltonians constructed from maximally localized Wannier functions⁹³.

Data availability

All data generated or analyzed during this study are included in this published article and its Supplementary Information files. Source data are provided with this paper.

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Author contributions

Y. Ma and M. Pan conceived and designed the experiments. J. Gao, C. Wang, Y. Yu, F. Song, and W. Huang helped to locate the research resource related to this work. Y. Cao, X. Tian, Y. Liu, and Z. Han performed HER measurements. Y. Cao, S. Xia, X. Tian, X. Zhang, J. Leng, and C. Zhou performed absorption and photoluminescence (PL) spectra

measurement. Y. Cao, J. He, Y. Zou, and W. Zhu grew the samples. R. Song carried out the DFT calculations. Y. Ma and M. Pan analyzed the data. M. Pan, C. Zhou, and Y. Ma wrote the manuscript with inputs from all other authors.

Competing interests

The authors declare no competing interests.

Additional information

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