



# OPEN Modulation instability, bifurcation analysis, and ion-acoustic wave solutions of generalized perturbed KdV equation with M-fractional derivative

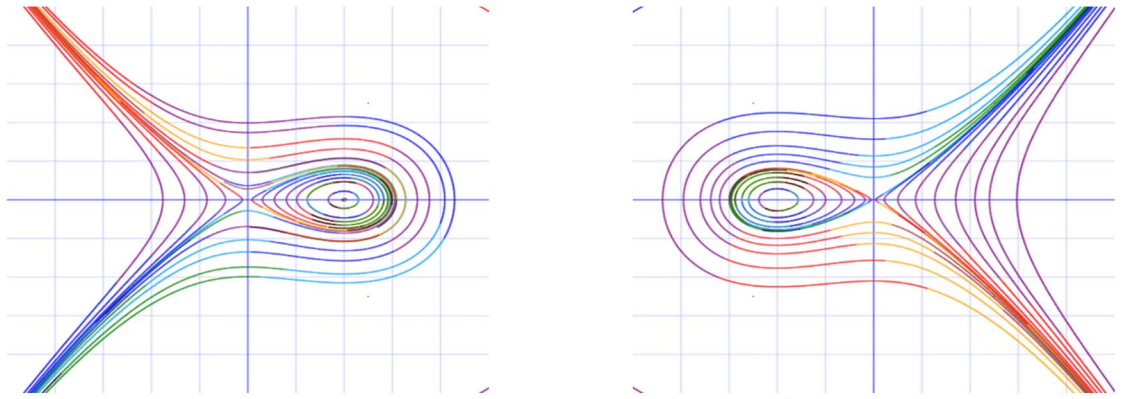
Mohammed Kbiri Alaoui<sup>1</sup>, Md. Mamunur Roshid<sup>2</sup>✉, Mahtab Uddin<sup>3</sup>, Wen-Xiu Ma<sup>4,5,6,7</sup>✉, Harun-Or-Roshid<sup>8</sup> & Mohammad Jahirul Haque Munshi<sup>2</sup>

The perturbed Korteweg-de Vries (PKdV) equation is essential for describing ion-acoustic waves in plasma physics, accounting for higher-order effects such as electron temperature variations and magnetic field influences, which impact their propagation and stability. This work looks at the generalized PKdV (gPKdV) equation with an M-fractional operator. It uses bifurcation theory to look at critical points and phase portraits, showing system changes such as shifts in stability and the start of chaos. Figures 1, 2 and 3 provide detailed analyses of static soliton formation through saddle-node bifurcation. We also use the modified simple equation (MSE) method to look for ion-acoustic wave solutions directly, without having to first define them. This lets us find shapes like hyperbolic, exponential, and trigonometric waves. These solutions reveal complex phenomena, including double periodic waves, periodic lump waves, bright bell-shaped waves, and singular soliton waves. Additionally, we analyze modulation instability in the gPKdV equation, which signifies chaotic transitions and is crucial for understanding nonlinear wave dynamics. Those methods demonstrate their value in generating precise soliton solutions relevant to nonlinear science and mathematical physics. This research illustrates how theoretical mathematics and physics can support solutions to practical world issues, especially in energy and technological advancement.

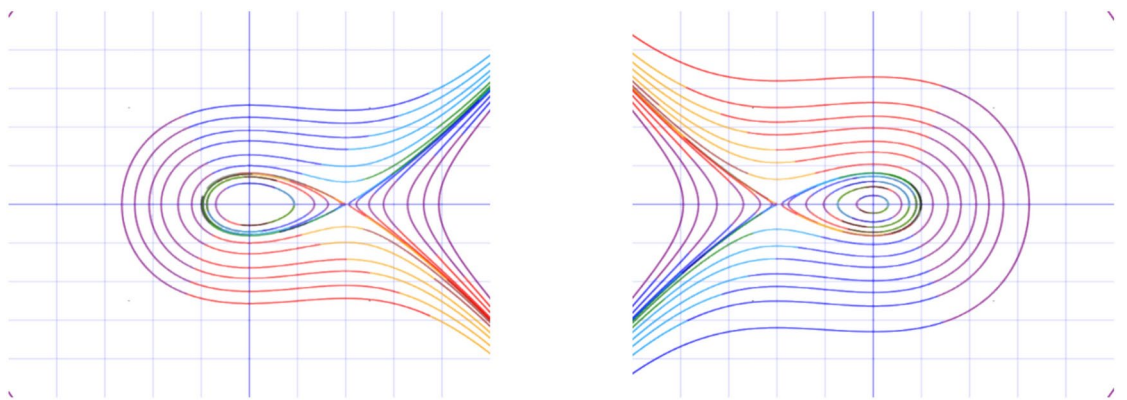
**Keywords** Modified simple equation technique, Generalized perturbed Korteweg-de Vries equation, Bifurcation theory, Phase portrait, Soliton solution, Clean energy technologies

Nonlinear evolution equations are pivotal in capturing the dynamics of complex systems across various scientific and engineering fields. Unlike linear equations, nonlinear ones can model phenomena where interactions lead to unpredictable and rich behaviors, such as turbulence in fluid dynamics<sup>1</sup>, chemical reaction pattern formation<sup>2</sup>, soliton of telecommunication systems<sup>3,4</sup>, quantum physics<sup>5,6</sup>, etc. These equations represent critical thresholds, bifurcations, and chaotic regimes, making them indispensable for understanding and predicting real-world phenomena. Their applications extend to physics, where they describe wave propagation and quantum mechanics, finance for modeling market fluctuations<sup>7</sup>, and ecology for ecosystem dynamics<sup>8</sup>. The study and solutions of nonlinear evolution equations deepen our theoretical understanding and drive advancements in technology and practical problem-solving across disciplines<sup>1–10</sup>. Numerous methods are continually being discovered to explore exact solutions of NLEEs such as multi exp-function technique<sup>11</sup>, Hirota bilinear process<sup>12,13</sup>, JEFÉ scheme<sup>14</sup>,

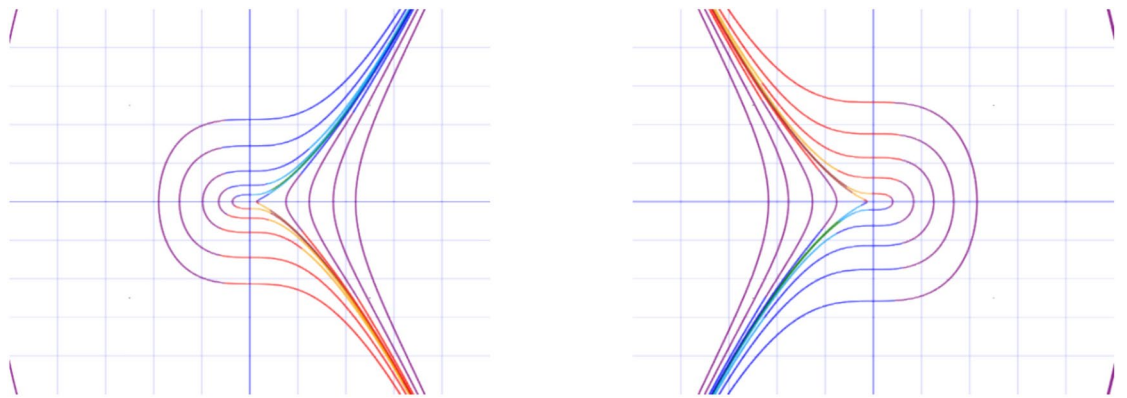
<sup>1</sup>Department of Mathematics, College of Science, King Khalid University, P.O. Box 9004, 61413 Abha, Saudi Arabia. <sup>2</sup>Department of Mathematics, Hamdard University Bangladesh, Dhaka, Bangladesh. <sup>3</sup>Institute of Natural Sciences, United International University, Dhaka, Bangladesh. <sup>4</sup>Department of Mathematics, Zhejiang Normal University, Jinhua 321004, Zhejiang, China. <sup>5</sup>Department of Mathematics, King Abdulaziz University, 21589 Jeddah, Saudi Arabia. <sup>6</sup>Department of Mathematics and Statistics, University of South Florida, Tampa, FL 33620-5700, USA. <sup>7</sup>Material Science Innovation and Modelling, Department of Mathematical Sciences, North-West University, Mafikeng Campus, Mmabatho 2735, South Africa. <sup>8</sup>Department of Mathematics, Sunamganj Science and Technology University, Shantiganj, Sunamganj, Bangladesh. ✉email: mamunmath307@gmail.com; wma3@usf.edu



**Fig. 1.** The phase portraits of the system (5) for  $2Ka_1(\epsilon - a_3K) > 0$ .



**Fig. 2.** The two-dimensional phase portraits of the system (5) for  $2Ka_1(\epsilon - a_3K) < 0$ .



**Fig. 3.** The two-dimensional phase portraits of the system (9) for  $\epsilon = a_3K$ .

variational technique<sup>15</sup>, enhanced MSE technique<sup>16</sup>, NK and improved F-expansion technique<sup>17</sup>, extended direct algebraic scheme<sup>18</sup>, MSSE technique<sup>19</sup>, new mapping scheme<sup>20</sup>, modified extended tanh and NMK techniques<sup>21</sup>, unified method<sup>22,23</sup>, exp-expansion and NMK schemes<sup>24</sup>, extended tanh expansion technique<sup>25</sup>, and so on<sup>26–29</sup>.

The generalized perturbed Korteweg-de Vries (gPKdV) equation is a mathematical model that describes the evolution of waves in a nonlinear and dispersive medium. The gPKdV equation<sup>30</sup> can be written as follows:

$$P_t + a_1PP_x + a_2P_{xxx} + a_3P_x = 0. \quad (1)$$

where  $P$  is the wave function,  $t$  is the time variable,  $x$  is the spatial variable, and  $a_1$ ,  $a_2$ , and  $a_3$  are constants that characterize the nonlinear, dispersive, and perturbative effects, respectively. The term  $a_1PP_x$  represents the

nonlinear interaction in the wave. Nonlinearity often leads to wave steepening, which can cause the formation of shock waves or solitons. The constant  $a_1$  determines the strength of this nonlinear effect. The term  $a_2 P_{xxx}$  accounts for the dispersion in the medium. Dispersion leads to the spreading of the wave packet over time and space. The constant  $a_2$  controls the degree of dispersion, balancing the nonlinear steepening effect. The term  $a_3 P_x$  represents a perturbative effect, which could arise from various physical mechanisms such as damping, external forces, or higher-order corrections. The constant  $a_3$  quantifies the magnitude of this perturbation. The coefficient of perturbation  $a_3$  plays an acute role in articulating the influence of the Coriolis parameter in the horizontal module. The model in Eq. (1) is employed to illustrate particular events in theoretical physics associated with quantum mechanics. This model is applicable for representing phenomena such as the formation of shock waves, solitons, turbulence, boundary layer dynamics, and mass transport, particularly within the areas of fluid dynamics, aerodynamics, and continuum mechanics<sup>33–35</sup>.

In Eq. (1), the parameter  $a_3$  represents the Coriolis effect, which is the deflection of moving objects like air and water caused by Earth's rotation. Equation (1) provides a generalized version of both the geophysical and classical KdV equations. The geophysical KdV equation<sup>36–38</sup> can be reconstructed by setting  $a_2 = \frac{3}{2}$ ,  $a_3 = \frac{1}{6}$  and  $a_1 = w_0$ . The perturbed-KdV model is commonly used in fields such as acoustics, aerodynamics, and medical engineering to explain sound propagation in fluids. Various significant characteristics and applications of the gPKdV equation have been thoroughly discussed in the literature<sup>30–32</sup>.

The primary objective of this work is to apply bifurcation theory to examine the critical points and phase portraits of the gPKdV model, where systems transition to new behaviors, such as changes in stability or the onset of chaos. To explore ion-acoustic wave solutions and the influence of fractional derivatives, we utilize a direct approach known as the modified simple equation technique on the M-fractional gPKdV model. Additionally, we illustrate some complex behaviors of the obtained solutions through three-dimensional, two-dimensional, and density diagrams, highlighting the impact of the fractional parameter in a two-dimensional diagram. Finally, we assess the modulation instability of the M-fractional gPKdV model.

The article is arranged as follows: Section two discusses the significance and features of the M-fractional derivative; Section three explores the working principles of the modified simple equation method; Section four applies bifurcation theory to Eq. (1) and presents the effect of parameters on equilibrium points and phase portraits. We execute orbits such as homoclinic, periodic, and heteroclinic, and obtain their corresponding phenomena, including periodic waves and kink waves. Section five implements the MSE technique to integrate the M-fractional gPKdV model. Section six discusses the numerical form of the obtained solutions with three-dimensional diagrams, density plots, and two-dimensional plots. Section seven presents the modulation instability of the M-fractional gPKdV model. Section eight provides the comparison, advantages, and limitations of this work and methodology. Finally, Section nine provides a summary of this work.

## Fractional derivative

Fractional derivatives are crucial in the study of nonlinear evolution equations (NLEEs) due to their ability to model complex phenomena with greater accuracy than integer-order derivatives. They provide a more flexible framework for capturing the memory and hereditary properties inherent in many physical, biological, and engineering systems. By using fractional derivatives, NLEEs can describe unusual diffusion and wave propagation in various types of media. This lets us make more accurate predictions and find better solutions. This improved modeling capability is essential in fields such as viscoelasticity, fluid dynamics, and signal processing. Using fractional derivatives in NLEEs also makes it easier to come up with new numerical and analytical methods that help us understand and solve difficult nonlinear problems. The growing interest in fractional calculus underscores its significance in advancing theoretical and applied research across various scientific disciplines such as<sup>39–42</sup>.

## Definition and some features of $M$ -fractional derivative

**Definition:** The mapping  $\varphi : [0, \infty) \rightarrow \mathfrak{R}$  and an order  $\chi$ ,  $M$ -fractional operator is described as:

$$D_{M,t}^{\chi,\psi} \varphi = \lim_{\epsilon \rightarrow 0} \frac{\varphi(t\phi_\psi(\epsilon t^{-\chi})) - \varphi(t)}{\epsilon}, t > 0, \psi > 0.$$

Here,  $\phi_\psi(x)$  is Mittag-Leffler function in one parameter clear as<sup>41</sup>, and taking belong to  $(0,1)$  :

$$\phi_\psi(x) = \sum_{n=0}^k \frac{x^n}{\Gamma(\phi n + 1)}.$$

**Features.** Let  $l, m \rightarrow \mathfrak{R}$ ,  $1 \geq \chi > 0$ . Let  $R, H$  be functions. Then.

$$\begin{aligned} D_{M,t}^{\chi,\psi}(lR + mH) &= lD_{M,t}^{\chi,\psi}(R) + mD_{M,t}^{\chi,\psi}(H). \\ D_{M,t}^{\chi,\psi}(RH) &= RD_{M,t}^{\chi,\psi}(H) + HD_{M,t}^{\chi,\psi}(R). \\ D_{M,t}^{\chi,\psi}\left(\frac{R}{H}\right) &= \frac{HD_{M,t}^{\chi,\psi}(R) - RD_{M,t}^{\chi,\psi}(H)}{H^2} \\ D_{M,t}^{\chi,\psi}(t^\varpi) &= \varpi t^{\varpi-\chi}, \varpi \in \mathfrak{R}. \\ D_{M,t}^{\chi,\psi}(P) &= 0, P \in \mathfrak{R}. \\ D_{M,t}^{\chi,\psi}(R \circ H) &= R'(H)D_{M,t}^{\chi,\psi}H(t), \text{ for } F \text{ is differentiable at } G. \end{aligned}$$

$$D_{M,t}^{\chi,\psi} R(t) = \frac{t^{1-\chi}}{\Gamma(\psi+1)} \frac{dR}{dt}, \text{ for } R \text{ is differentiable.}$$

### Methodology of MSE technique

In this point, the working rule of the MSE technique<sup>43,44</sup> is useful to solve any nonlinear PDEs problem. The main advantage of the MSE method is its ability to directly investigate traveling wave solutions for nonlinear partial differential equations (NLPDEs). This method does not use any auxiliary equation to find the solution such as an extended direct algebraic scheme<sup>18</sup>, MSSE technique<sup>19</sup>, new mapping scheme<sup>20</sup>, modified extended tanh and NMK techniques<sup>21</sup>, unified method<sup>22,23</sup> and so on. For this purpose, we study nonlinear PDEs in the succeeding form:

$$L(G_t, G_x, G_{xxx}, G_{tx}) = 0. \quad (2)$$

Using the relation  $\zeta = \Xi - \Theta t$ ;  $G(x, t) = G(\zeta)$  in Eq. (2) then we attain,

$$L(-\Theta G_\zeta, \Xi^2 G_\zeta, \Xi^3 G_{\zeta\zeta\zeta}, -\Theta \Xi G_{\zeta\zeta}) = 0. \quad (3)$$

The proposed trial solution is:

$$G(\zeta) = \sum_{q=0}^s \left( p_q \left( \frac{H'(\zeta)}{H(\zeta)} \right)^q \right). \quad (4)$$

Here,  $q$  is a balanced number that can be calculated using the following formula,

$$\frac{d^p G}{d\zeta^p} = q + p \text{ and } G^N \frac{d^p G}{d\zeta^p} = Nq + (q + p).$$

Now we use the proposed trial solution Eq. (4) in Eq. (3), then we get  $P(H(\zeta)) = C_0 + C_1 H(\zeta)^{-1} + C_2 H(\zeta)^{-2} + C_3 H(\zeta)^{-3} + \dots + C_n H(\zeta)^{-n}$ . If we set the coefficient  $C_k = 0$ ;  $k = 0, 1, 2, 3, \dots, n$ , then we get a system of equations. To get the values of  $p_q, \Xi, H(\zeta), H'(\zeta), \Theta$ , we solve the obtained system of equations. Substituting these parameters into Eq. (4) yields the required solutions.

### Bifurcation analysis and phase portrait

In this section, we apply the bifurcation theory for analysis of the phase portrait of the time M-fractional generalized perturbed KdV (tMf-gPKdV) equation. To fully grasp the dynamics of nonlinear wave propagation, especially in the context of solitons and other wave phenomena, bifurcation analysis of the gPKdV equation is essential. The gPKdV equation models the evolution of shallow water waves and other physical systems with dispersive nonlinearity. Bifurcation analysis identifies critical parameter values at which the system undergoes qualitative changes, such as the creation or destruction of solitons. This analysis reveals how system behavior transitions between stable and unstable states, offering insights into wave stability and the formation of complex patterns. Phase portraits, which represent the state space of the system, provide a visual method for studying the stability and dynamics of solutions. By analyzing fixed points, limit cycles, and the system's behavior near bifurcation points, one can predict the system's long-term behavior. Fluid dynamics, optical fibers, and nonlinear media widely apply this approach for wave stability analysis and pattern formation.

The time M-fractional generalized perturbed KdV (tMf-gPKdV) equation is considered in the subsequent arrangement:

$$D_{M,t}^{\sigma,n} P + a_1 P P_x + a_2 P_{xxx} + a_3 P_x = 0. \quad (5)$$

Using the relation  $\varphi = Kx - \epsilon \frac{\Gamma(n+1)t^\sigma}{\sigma}$ ;  $P(x, t) = P(\varphi)$  in Eq. (5) and we attain,

$$-\epsilon P' + K a_1 P P' + K^3 a_2 H''' + a_3 K P' = 0. \quad (6)$$

Integrating one time and setting integrating constant to zero.

$$K a_1 \frac{P^2}{2} + K^3 a_2 P'' + (a_3 K - \epsilon) P = 0. \quad (7)$$

From Eq. (7), we develop as,

$$\frac{d^2 P}{d\varphi^2} = \frac{-(a_3 K - \epsilon)}{K^3 a_2} P + \frac{-K a_1}{2 K^3 a_2} P^2. \quad (8)$$

According to the bifurcation theory, the Eq. (8) is as,

$$\begin{cases} \frac{dP}{d\varphi} = H, \\ \frac{dH}{d\varphi} = \frac{-(a_3 K - \epsilon)P - \frac{K a_1}{2} P^2}{K^3 a_2}, \end{cases} \quad (9)$$

$$\mathcal{H}(H, P) = \frac{H^2}{2} - \frac{1}{K^3 a_2} \left( (\epsilon - a_3 K) \frac{P^2}{2} - \frac{K a_1}{6} P^3 \right) = L. \quad (10)$$

Here,  $L$  is the Hamiltonian constant.

The charge of  $L$ , Eq. (10) generates a phase portrait of Eq. (9). When  $L$  varies, numerous types of orbits emerge in Eq. (10), resulting in diverse dynamical behaviors as explained for Eq. (10). According to bifurcation theory<sup>39</sup>, the periodic heteroclinic, and homoclinic orbit are identified by using phase portrait diagram. From Eq. (9), we attain the point of equilibrium as follows:

$$\sum_a^b \left\{ \begin{array}{l} H = 0, \\ \frac{-(a_3 K - \epsilon)P - \frac{K a_1}{2} P^2}{K^3 a_2} = 0. \end{array} \right. \quad (11)$$

From the Eq. (11), we get the  $E_0 = (0, 0)$  and  $E_1 = \left( \frac{2(\epsilon - a_3 K)}{K a_1}, 0 \right)$ . The stability of equilibrium points is checked now. The Jacobian matrix of the linearized Eq. (9) is now shown to be:

$$J_E = \begin{pmatrix} 0 & 1 \\ \frac{-(a_3 K - \epsilon) - \frac{K a_1}{2} P_E}{K^3 a_2} & 0 \end{pmatrix},$$

$$D(E) = -\frac{-(a_3 K - \epsilon) - \frac{K a_1}{2} P_E}{K^3 a_2}.$$

According to the theory of planar, the observations are: if  $D(E) < 0$ , then  $E$  is the saddle point. If  $D(E) > 0$  and  $T(E) = 0$ , then the point  $E$  is the center point. If  $D(E) = 0$ , then  $E$  is the cusp point.

Therefore, we get.

- (a) The point  $E_0$  is saddle and  $E_1$  is the center and when  $2K a_1 (\epsilon - a_3 K) > 0$ , and  $D(E_1) > 0$ ,  $D(E_0) < 0$ ,  $T(E_1) = 0$ .
- (b) The point  $E_1$  is saddle and  $E_0$  is the center and when  $2K a_1 (\epsilon - a_3 K) < 0$ , and  $D(E_1) > 0$ ,  $D(E_0) < 0$ ,  $T(E_1) = 0$ .

According to the equilibrium point, the values of  $L$  are:

$$L_0 = \mathcal{H}(E_0) = 0,$$

$$L_1 = \mathcal{H}(E_1) = -\frac{2(a_3 K - \epsilon)^3}{3(K a_1)^2 K^3 a_2}.$$

The behavior of Figure 1, 2, and 3, as follows:

When  $2K a_1 (\epsilon - a_3 K) > 0$ , we obtain.

- (i) Eq. (1) has the bell-type solitary wave if  $L = L_0$ , and the orbit of the homoclinic is obtained.
- (ii) Eq. (1) has the periodic solitary wave if  $L < L_0$ , and the orbit of the heteroclinic is obtained.

When  $2K a_1 (\epsilon - a_3 K) < 0$ , we obtain.

- (i) Eq. (1) has the bell-type solitary wave if  $L = L_0$ , and the orbit of the homoclinic is obtained.
- (ii) Eq. (1) has the periodic solitary wave if  $L < L_0$ , and the orbit of the heteroclinic is obtained.

When  $\epsilon = a_3 K$ .

If  $L = 0$ , then the Eq. (10) developed as:

$$H = 2\sqrt{\frac{1}{K^3 a_2} \left( (\epsilon - a_3 K) \frac{P^2}{2} - \frac{K a_1}{6} P^3 \right)}.$$

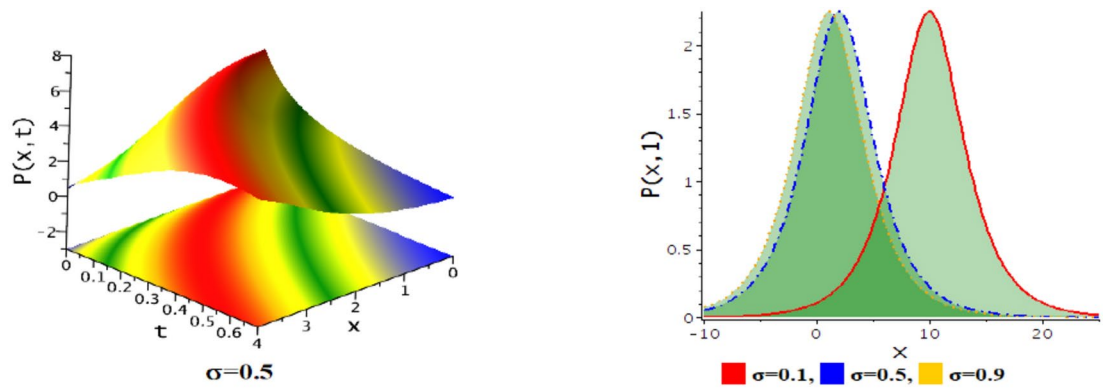
When  $2K a_1 (\epsilon - a_3 K) > 0$ , the heteroclinic orbit is obtained in  $(P, H)$ -plane from Eq. (10). If we set the value of  $H$  into  $dP/d\xi = H$ , then we attain,

$$P(x, t) = \frac{3(\epsilon - a_3 K)}{K a_1} \operatorname{sech}^2 \left( \frac{1}{2} \sqrt{\frac{(\epsilon - a_3 K)}{K^3 a_2}} \varphi \right) \quad (12)$$

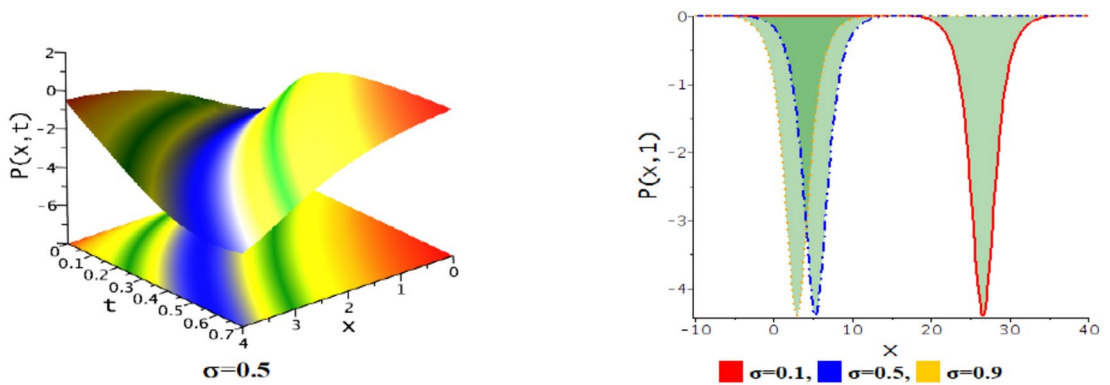
Figure 4 for the values  $\epsilon = 2, K = 1, a_1 = 1.5, a_2 = 1, a_3 = -0.2, n = 1.5$ , and Figure 5 for the values  $\epsilon = 2, K = 1, a_1 = 1.5, a_2 = -1, a_3 = -0.2, n = 1.5$ .

When  $2K a_1 (\epsilon - a_3 K) < 0$ , the heteroclinic orbit is obtained in  $(P, H)$ -plane from Eq. (10). Then similar to Eq. (12) we attain, The solitary periodic wave solution is shown in Fig. 6. for parametric values  $\epsilon = 1, K = 1, a_1 = 1.5, a_2 = 1, a_3 = 2, n = 1.5$

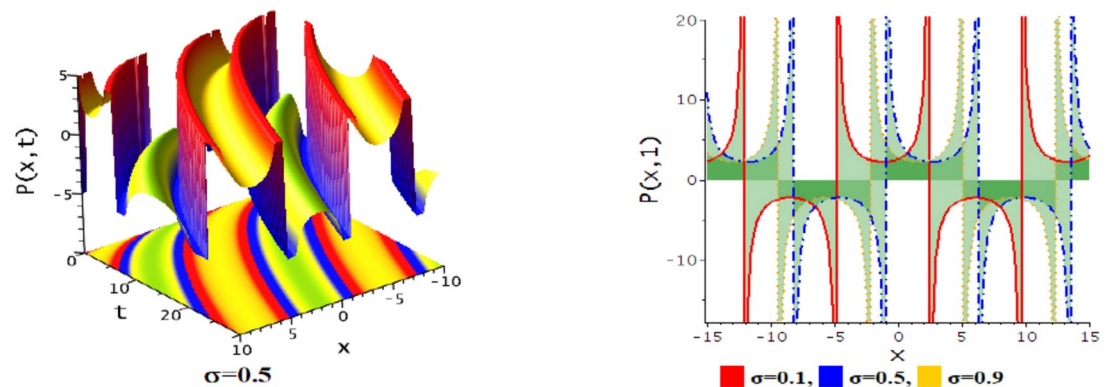




**Fig. 4.** Visualization of the bright bell wave of the solution Eq. (12) for specific parametric values  $\epsilon = 2, K = 1, a_1 = 1.5, a_2 = 1, a_3 = -0.2, n = 1.5$ .



**Fig. 5.** Visualization of dark bell shape wave of the solution Eq. (12) for specific parametric values  $\epsilon = 2, K = 1, a_1 = 1.5, a_2 = -1, a_3 = -0.2, n = 1.5$ .



**Fig. 6.** Visualization of the periodic wave of the solution Eq. (13) for specific parametric values  $\epsilon = 1, K = 1, a_1 = 1.5, a_2 = 1, a_3 = 2, n = 1.5$ .

$$P(x, t) = \frac{3(\epsilon - a_3 K)}{K a_1} \sec^2 \left( \frac{1}{2} \sqrt{\frac{(\epsilon - a_3 K)}{K^3 a_2}} \varphi \right). \quad (13)$$

### Application of modified simple equation technique

In this section, we solve the gPKdV model by using the MSE technique analytically. Under specific conditions on the free parameters, we expressed the obtained solution in terms of trigonometric and hyperbolic function

forms. The main advantage of the proposed method is to solve all NLEEs directly. To solve the NLEEs, the MSE method doesn't use any auxiliary equation or any predefined solutions like as tanh-coth method<sup>1</sup>, ( $G'/G$ )-expansion technique<sup>2</sup>, Sardar sub-equation method<sup>3</sup> and others<sup>4–10</sup>, it derives the exact solitary wave solution directly.

According to the MSE technique the balance number  $s = 2$ . So, the solution of the Eq. (7) is:

$$P = \beta_0 + \beta_1 \left( \frac{d}{d\varphi} \psi(\varphi) \right) + \beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2. \quad (14)$$

If we use the trial solution in Eq. (14) into Eq. (7), then the algebraic equations are obtained below:

$$\begin{aligned} -\epsilon\beta_0 + \frac{1}{2}Ka_1\beta_0^2 + a_3K\beta_0 &= 0, \\ K^3a_2\beta_1 \left( \frac{d^3}{d\varphi^3} \psi(\varphi) \right) + Ka_1\beta_1 \left( \frac{d}{d\varphi} \psi(\varphi) \right) \beta_0 + a_3K\beta_1 \left( \frac{d}{d\varphi} \psi(\varphi) \right) - \epsilon\beta_1 \left( \frac{d}{d\varphi} \psi(\varphi) \right) &= 0, \\ -\epsilon\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2 + \frac{1}{2}Ka_1 \left( 2\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2 \beta_0 + \beta_1^2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2 \right) + K^3a_2 \left( 2\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right) \left( \frac{d^3}{d\varphi^3} \psi(\varphi) \right) - 3\beta_1 \left( \frac{d^2}{d\varphi^2} \psi(\varphi) \right) \left( \frac{d}{d\varphi} \psi(\varphi) \right) + 2\beta_2 \left( \frac{d^2}{d\varphi^2} \psi(\varphi) \right)^2 \right) + a_3K\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2 &= 0, \\ Ka_1\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^3 \beta_1 + K^3a_2 \left( 2\beta_1 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^3 - 10\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^2 \left( \frac{d^2}{d\varphi^2} \psi(\varphi) \right) \right) &= 0, \\ \frac{1}{2}Ka_1\beta_2^2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^4 + 6K^3a_2\beta_2 \left( \frac{d}{d\varphi} \psi(\varphi) \right)^4 &= 0. \end{aligned}$$

Using the Maple software, we get the following solution sets.

**Set 01:**  $K = \sqrt{\frac{-a_1\beta_2}{12a_2}}, \beta_0 = -\frac{2(a_3K - \epsilon)}{Ka_1}, \epsilon = -\frac{1}{72\beta_2} \sqrt{\frac{-3a_1\beta_2}{a_2}} (a_1\beta_1^2 + 2a_3\beta_2),$   
 $\psi(\varphi) = h_1 + h_2 e^{-\frac{\beta_1\varphi}{\beta_2}}.$

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2 h_2 e^{-\frac{\beta_1\varphi}{\beta_2}}}{\beta_2 \left( h_1 + h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)} + \frac{\beta_1^2 \left( h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)^2}{\beta_2 \left( h_1 + h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)^2}. \quad (15)$$

Here,  $\theta = -\frac{\beta_1\varphi}{\beta_2}, \varphi = \sqrt{\frac{-a_1\beta_2}{12a_2}}x + \frac{1}{72\beta_2} \sqrt{\frac{-3a_1\beta_2}{a_2}} (a_1\beta_1^2 + 2a_3\beta_2) \frac{\Gamma(n+1)t^\sigma}{\sigma}.$

If  $a_1\beta_2a_2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{h_1 h_2}{\cosh\theta (h_1^2 + h_2^2) + 2h_1 h_2 + \sinh\theta (h_1^2 - h_2^2)} \right]. \quad (16)$$

If  $h_1 = h_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{1}{2} \operatorname{sech} \left( \frac{\theta}{2} \right) \right]. \quad (17)$$

If  $h_1 = \pm i h_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{1}{1 \mp i \sinh\theta} \right]. \quad (18)$$

If  $a_1^2 a_2 a_3 \beta_1^2 + 36\epsilon^2 a_2^2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{h_1 h_2}{\cos\theta (h_1^2 + h_2^2) + 2h_1 h_2 + i \sin\theta (h_1^2 - h_2^2)} \right]. \quad (19)$$

If  $h_1 = h_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{1}{2} \sec \left( \frac{\theta}{2} \right) \right]. \quad (20)$$

If  $h_1 = \pm ih_2$ , then Eq. (15) becomes,

$$P(x, t) = -\frac{2(a_3K - \epsilon)}{Ka_1} - \frac{\beta_1^2}{\beta_2} \left[ \frac{1}{1 \pm \sinh\theta} \right]. \quad (21)$$

Here,  $\theta = -\frac{\beta_1\varphi}{\beta_2}$ ,  $\varphi = \sqrt{\frac{-a_1\beta_2}{12a_2}}x + \frac{1}{72\beta_2}\sqrt{\frac{-3a_1\beta_2}{a_2}}(a_1\beta_1^2 + 2a_3\beta_2)\frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

**Set 02:**  $K = \sqrt{\frac{-a_1\beta_2}{12a_2}}$ ,  $\beta_0 = 0$ ,  $\epsilon = -\frac{1}{72\beta_2}\sqrt{\frac{-3a_1\beta_2}{a_2}}(a_1\beta_1^2 + 2a_3\beta_2)$ ,  
 $\psi(\varphi) = h_1 + h_2e^{-\frac{\beta_1\varphi}{\beta_2}}$ .

$$P(x, t) = -\frac{\beta_1^2 h_2 e^{-\frac{\beta_1\varphi}{\beta_2}}}{\beta_2 \left( h_1 + h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)} + \frac{\beta_1^2 \left( h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)^2}{\beta_2 \left( h_1 + h_2 e^{-\frac{\beta_1\varphi}{\beta_2}} \right)^2}. \quad (22)$$

Here,  $\varphi = \sqrt{\frac{-a_1\beta_2}{12a_2}}x - \frac{1}{2}\sqrt{\frac{-a_1\beta_2}{12a_2}}(a_1\beta_0 + 2a_3)\frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

If  $\beta_0\beta_2 > 0$  and  $a_2a_1\beta_2 < 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{h_1 h_2}{\cosh\theta(h_1^2 + h_2^2) + 2h_1 h_2 + \sinh\theta(h_1^2 - h_2^2)} \right]. \quad (23)$$

If  $h_1 = h_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{1}{2} \operatorname{sech}\left(\frac{\theta}{2}\right) \right]. \quad (24)$$

If  $h_1 = \pm ih_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{1}{1 \mp i \sinh\theta} \right]. \quad (25)$$

If  $\beta_0\beta_2 < 0$  or  $a_2a_1\beta_2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{h_1 h_2}{\cos\theta(h_1^2 + h_2^2) + 2h_1 h_2 + i \sin\theta(h_1^2 - h_2^2)} \right]. \quad (26)$$

If  $h_1 = h_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{1}{2} \sec\left(\frac{\theta}{2}\right) \right]. \quad (27)$$

If  $h_1 = \pm ih_2$ , then Eq. (22) becomes,

$$P(x, t) = -\frac{\beta_1^2}{\beta_2} \left[ \frac{1}{1 \pm \sinh\theta} \right]. \quad (28)$$

Here  $\theta = \frac{\beta_1\varphi}{\beta_2}$ ,  $\varphi = \sqrt{\frac{-a_1\beta_2}{12a_2}}x + \frac{1}{72\beta_2}\sqrt{\frac{-3a_1\beta_2}{a_2}}(a_1\beta_1^2 + 2a_3\beta_2)\frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

**Set 03:**  $\beta_0 = 0$ ,  $\beta_2 = -\frac{12K^2a_2}{a_1}$ ,  $\beta_1 = \frac{12}{a_1}\sqrt{K\epsilon a_2 - K^2a_3a_2}$ ,  
 $\psi(\varphi) = h_1 + h_2e^{\frac{\sqrt{K\epsilon a_2 - K^2a_3a_2}}{K^2a_2}\varphi}$ .

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2a_3a_2)}{K^2a_2a_1} \left( \frac{h_2e^\theta}{(h_1 + h_2e^\theta)} - \frac{(h_2e^\theta)^2}{(h_1 + h_2e^\theta)^2} \right). \quad (29)$$

Here,  $\theta = \frac{\sqrt{K\epsilon a_2 - K^2a_3a_2}}{K^2a_2}\varphi$ ,  $\varphi = Kx - \epsilon\frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

If  $K\epsilon a_2 - K^2a_3a_2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , Then Eq. (29) becomes,



$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{h_1 h_2}{\cosh \theta (h_1^2 + h_2^2) + 2h_1 h_2 + \sinh \theta (h_1^2 - h_2^2)} \right]. \quad (30)$$

If  $h_1 = h_2$ , then Eq. (29) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{2} \operatorname{sech} \left( \frac{\theta}{2} \right) \right]. \quad (31)$$

If  $h_1 = \pm i h_2$ , then Eq. (29) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{1 \mp i \sinh \theta} \right]. \quad (32)$$

If  $K\epsilon a_2 - K^2 a_3 a_2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (29) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{h_1 h_2}{\cos \theta (h_1^2 + h_2^2) + 2h_1 h_2 + i \sin \theta (h_1^2 - h_2^2)} \right]. \quad (33)$$

If  $h_1 = h_2$ , then Eq. (29) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{2} \sec \left( \frac{\theta}{2} \right) \right]. \quad (34)$$

If  $h_1 = \pm i h_2$ , then Eq. (29) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{1 \pm i \sin \theta} \right]. \quad (35)$$

Here,  $\theta = \frac{\sqrt{K\epsilon a_2 - K^2 a_3 a_2}}{K^2 a_2} \varphi$ ,  $\varphi = Kx - \epsilon \frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

**Set 04:**  $\beta_0 = -\frac{2(Ka_3 - \epsilon)}{Ka_1}$ ,  $\beta_2 = -\frac{12K^2 a_2}{a_1}$ ,  $K = \frac{144a_2\epsilon + 24\sqrt{a_1^2 a_2 a_3 \beta_1^2 + 36\epsilon^2 a_2^2}}{288a_2 a_3}$ ,  
 $\psi(\varphi) = h_1 + h_2 e^\theta$ .

$$P(x, t) = -\frac{576 \left( \frac{L}{288a_2} - \epsilon \right) a_2 a_3}{La_1} + \frac{6912a_1 a_2 a_3^2 \beta_1^2}{L^2} \left( \frac{h_2 e^\theta}{(h_1 + h_2 e^\theta)} - \frac{(h_2 e^\theta)^2}{(h_1 + h_2 e^\theta)^2} \right). \quad (36)$$

Here,  $\theta = \frac{6912a_1 \beta_1 a_2 a_3^2}{L^2} \varphi$ ,  $\varphi = \frac{144a_2\epsilon + 24\sqrt{a_1^2 a_2 a_3 \beta_1^2 + 36\epsilon^2 a_2^2}}{288a_2 a_3} x - \epsilon \frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

If  $a_1^2 a_2 a_3 \beta_1^2 + 36\epsilon^2 a_2^2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{h_1 h_2}{\cosh \theta (h_1^2 + h_2^2) + 2h_1 h_2 + \sinh \theta (h_1^2 - h_2^2)} \right]. \quad (37)$$

If  $h_1 = h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{2} \operatorname{sech} \left( \frac{\theta}{2} \right) \right]. \quad (38)$$

If  $h_1 = \pm i h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{1 \mp i \sinh \theta} \right]. \quad (39)$$

If  $a_1^2 a_2 a_3 \beta_1^2 + 36\epsilon^2 a_2^2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{h_1 h_2}{\cos \theta (h_1^2 + h_2^2) + 2h_1 h_2 + i \sin \theta (h_1^2 - h_2^2)} \right]. \quad (40)$$

If  $h_1 = h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{2} \sec\left(\frac{\theta}{2}\right) \right]. \quad (41)$$

If  $h_1 = \pm i h_2$ , then Eq. (36) becomes,

$$P(x, t) = \frac{12(K\epsilon a_2 - K^2 a_3 a_2)}{K^2 a_2 a_1} \left[ \frac{1}{1 \pm \sin\theta} \right]. \quad (42)$$

Here,  $\theta = \frac{\sqrt{K\epsilon a_2 - K^2 a_3 a_2}}{K^2 a_2} \varphi$ ,  $\varphi = Kx - \epsilon \frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

**Set 05:**

$$K = \sqrt{\frac{-a_1 \beta_2}{12a_2}}, \epsilon = \frac{1}{2} \sqrt{\frac{-a_1 \beta_2}{12a_2}} (a_1 \beta_0 + 2a_3), \beta_1 = \sqrt{6\beta_0 \beta_2}, \psi(\varphi) = h_1 + h_2 e^{-\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}}.$$

$$P(x, t) = \beta_0 - \frac{6\beta_0 h_2 e^{-\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}}}{h_1 + h_2 e^{-\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}}} + \frac{6\beta_0 \left( h_2 e^{-\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}} \right)^2}{\left( h_1 + h_2 e^{-\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}} \right)^2}. \quad (43)$$

Here,  $\varphi = \sqrt{\frac{-a_1 \beta_2}{12a_2}} x - \frac{1}{2} \sqrt{\frac{-a_1 \beta_2}{12a_2}} (a_1 \beta_0 + 2a_3) \frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

If  $\beta_0 \beta_2 > 0$  and  $a_2 a_1 \beta_2 < 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{h_1 h_2}{\cosh\theta (h_1^2 + h_2^2) + 2h_1 h_2 + \sinh\theta (h_1^2 - h_2^2)} \right]. \quad (44)$$

If  $h_1 = h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{1}{2} \operatorname{sech}\left(\frac{\theta}{2}\right) \right]. \quad (45)$$

If  $h_1 = \pm i h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{1}{1 \mp i \sinh\theta} \right]. \quad (46)$$

If  $\beta_0 \beta_2 < 0$  or  $a_2 a_1 \beta_2 > 0$  then the following forms are derived.

If  $h_1 \neq h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{h_1 h_2}{\cos\theta (h_1^2 + h_2^2) + 2h_1 h_2 + i \sin\theta (h_1^2 - h_2^2)} \right]. \quad (47)$$

If  $h_1 = h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{1}{2} \sec\left(\frac{\theta}{2}\right) \right]. \quad (48)$$

If  $h_1 = \pm i h_2$ , then Eq. (43) becomes,

$$P(x, t) = \beta_0 - 6\beta_0 \left[ \frac{1}{1 \pm i \sin\theta} \right]. \quad (49)$$

Here,  $\theta = -\frac{\sqrt{6\beta_0 \beta_2} \varphi}{\beta_2}$ ,  $\varphi = \sqrt{\frac{-a_1 \beta_2}{12a_2}} x - \frac{1}{2} \sqrt{\frac{-a_1 \beta_2}{12a_2}} (a_1 \beta_0 + 2a_3) \frac{\Gamma(n+1)t^\sigma}{\sigma}$ .

## Figure analysis

This section provides an in-depth examination of the graphical representation of solutions derived from the M-fractional generalized perturbed Korteweg–de Vries (gPKdV) equations, focusing on their waveforms and unique characteristics. The analysis highlights the importance of ion acoustic waves within the gPKdV framework due to their stability and persistence, making them essential for modeling nonlinear wave phenomena in environments like shallow water and plasma systems. The gPKdV equation integrates perturbation terms that account for real-world complexities, including higher-order dispersion and nonlinear effects, allowing it to model a wide range of wave behaviors. Solitary wave solutions arise from the balance between nonlinearity, which steepens the wave, and dispersion, which spreads it. This balance leads to stable, localized waveforms that travel without changing shape. Numerical simulations illustrate diverse wave patterns using 3D density plots

and 2D representations. These include double periodic waves, interactions between kink and periodic lump waves, periodic rogue waves, bell-shaped bright and dark waves, singular soliton waves, and V-shaped periodic rogue waves. Different waveforms serve distinct purposes in modeling physical phenomena. Kink waves, which describe transitions between states, are crucial for systems with topological structures and phase interfaces. Bell waves, balancing nonlinearity and dispersion, are useful for studying localized energy transport in optical fibers and water waves. Periodic waves capture oscillatory behavior and pattern formation, aiding in the study of fluids and plasmas. Rogue waves, characterized by extreme amplitudes and sudden energy localization, are critical for understanding rare events in fields like oceanography and nonlinear optics.

The gPKdV equation effectively models nonlinear wave phenomena across diverse systems, enhancing insights into wave stability, energy transport, and the emergence of complex patterns. This comprehensive approach reveals the equation's versatility in capturing the intricate dynamics of nonlinear interactions. The effect of  $M$ -fractional parameters is shown in two-dimensional plots, here the red color is used for  $\sigma = 0.1$ ; the red color is used for  $\sigma = 0.5$ ; the red color is used for  $\sigma = 0.9$ . Figure 7 represents the double periodic wave of Eq. (19) for specific parametric values  $a_1 = 0.2, a_2 = 0.5, a_3 = 0.2, h_1 = 2, h_2 = -0.2, \beta_1 = 3, \beta_2 = 1, n = 1.5$ . Figure 8 represents the periodic lump-type wave of the solution Eq. (20) for specific parametric values  $a_1 = 0.2, a_3 = 0.2, h_1 = 2, a_2 = 0.5, h_2 = -2, \beta_1 = 3, \beta_2 = 1, n = 1.5$ . The real portion of this solution provides an interaction of periodic lump wave and kink wave and the imaginary portion characterizes the episodic lump wave. Figure 9 represents the bell shape wave of Eq. (25) for specific parametric values  $h_1 = 1, a_3 = 0.1, \beta_2 = 0.5, a_1 = 1, h_2 = -i, \beta_1 = 0.5, a_2 = -0.1, n = 1.5$ . The real portion characterizes a bright bell shape wave and the imaginary portion provides a singular soliton wave. Figure 10 visualizes the wave of the interaction of periodic wave and kink of Eq. (32) for specific parametric values  $\epsilon = 1, a_1 = 1, a_2 = -0.1, a_3 = 0.1, h_1 = 1, h_2 = -i, K = 0.5, \beta_2 = 0.5, n = 1.5$ . Figure 11, Visualizes the wave of double periodic wave periodic lump type wave interaction of periodic lump wave and kink wave episodic lump wave bright bell shape wave singular soliton wave interaction of periodic wave and kink wave of Eq. (33) for specific parametric values  $\epsilon = -1, a_1 = 1, a_2 = 3, a_3 = 0.5, h_1 = 0.5, h_2 = 1, K = 0.5, \beta_2 = 0.5, n = 1.5$ . Figure 12, visualizes the diagram of the V-shape periodic rogue wave of Eq. (41) for specific parametric values  $\epsilon = 3, a_1 = 3, a_2 = 0.5, a_3 = -3, h_1 = 1, h_2 = 1, \beta_1 = 3, n = 1.5$ . Figure 13, visualizes the diagram of Eq. (47) for specific parametric values  $a_1 = 1, h_1 = 1, a_2 = 1, \beta_0 = -0.5, a_3 = 2, \beta_2 = 0.5, n = 1.5, h_2 = -2$ . The real portion provides a linked rogue wave with a dark bell shape and the imaginary part provides an episodic rogue wave. Figure 14, visualizes the diagram of Eq. (48) for the values  $a_1 = 1, a_2 = 1, a_3 = 2, h_1 = 1, h_2 = 1, \beta_0 = 0.5, \beta_2 = 0.5, n = 1.5$ . The real portion provides the interaction of periodic wave and kink and the imaginary portion characterizes periodic lump wave.

### Modulation instability

Modulation instability (MI)<sup>45–48</sup> refers to the exponential growth of small perturbations in a continuous wave or a uniform background, leading to the formation of localized structures or patterns. This phenomenon arises in nonlinear and dispersive media and plays a crucial role in various physical systems, including optical fibers, fluid dynamics, and plasma physics. The underlying mechanism of MI involves a balance between nonlinearity and dispersion (or diffraction), where a small initial disturbance can draw energy from the continuous background. This process amplifies specific frequencies, causing the system to evolve into localized wave packets or soliton-like structures. In this section, we investigate the modulation instability of traveling waves for  $M$ -fractional perturbed Korteweg-de Vries. Modulation instability in the gPKdV equations is significant because it describes the exponential growth of small perturbations in a wave train, leading to the formation of localized, high-amplitude structures. Nonlinear and dispersive effects drive this instability, which is crucial in generating rogue waves and other extreme events in fluids, plasmas, and optical systems. By analyzing modulation instability, researchers can better understand wave-breaking mechanisms, pattern formation, and the transition from regular wave patterns to chaotic states, enhancing the predictive power and real-world applicability of the gPKdV model in diverse nonlinear wave phenomena.

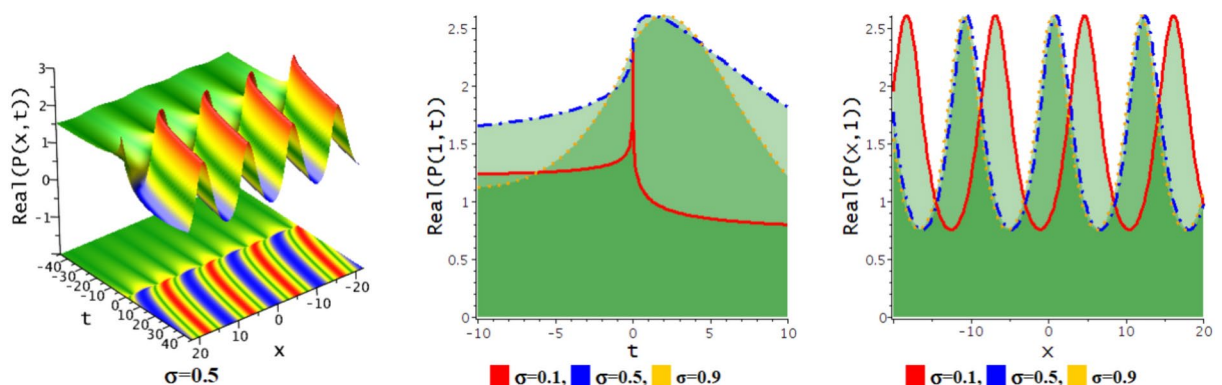


Fig. 7. Periodic wave of the solution Eq. (19).

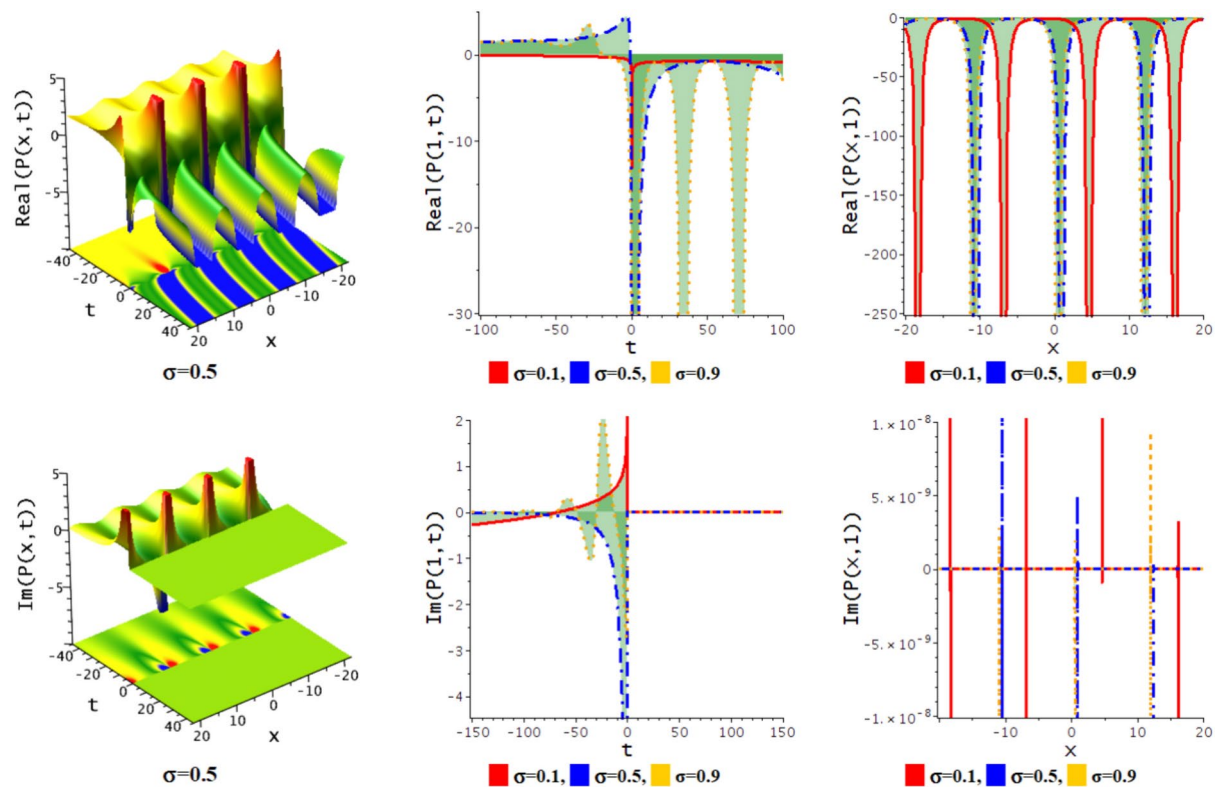


Fig. 8. Periodic wave of the solution Eq. (20).

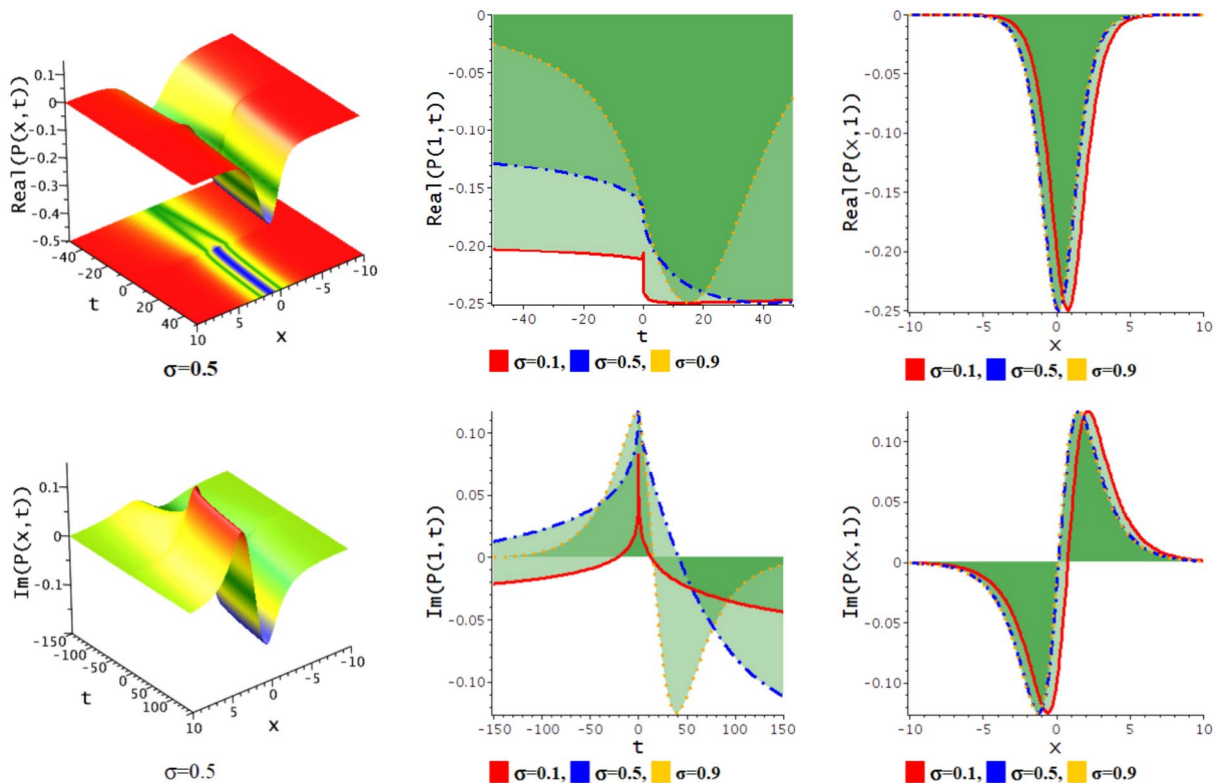


Fig. 9. Double bell waves the solution Eq. (25).

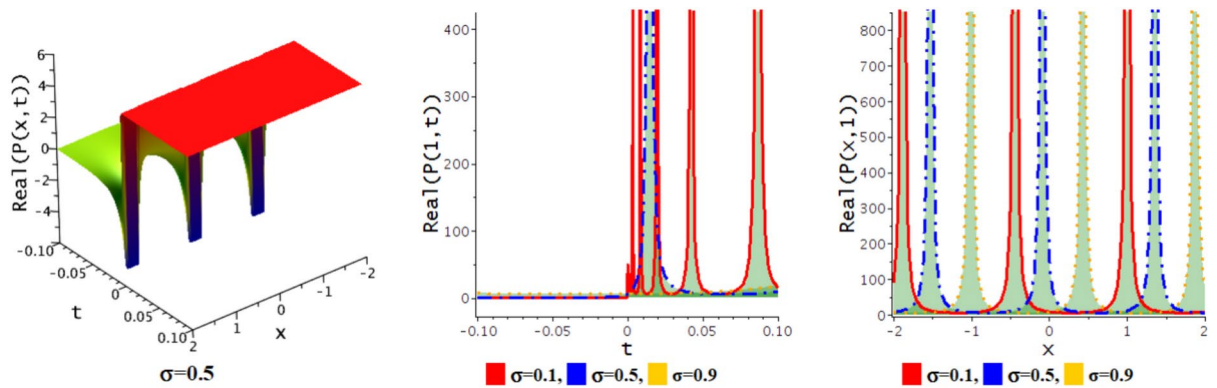


Fig. 10. Kinky periodic of the solution Eq. (32).

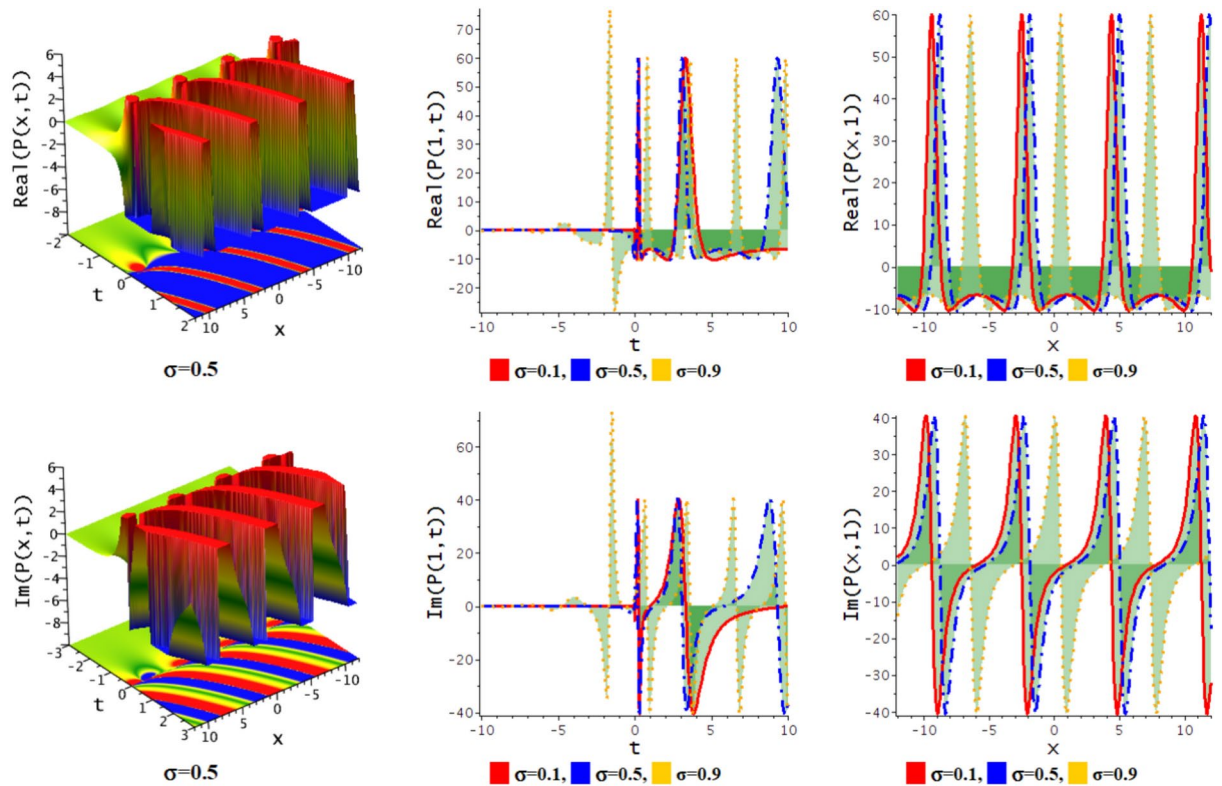


Fig. 11. Kinky periodic wave of the solution Eq. (33).

$$D_{M,t}^{\sigma,n} P + a_1 P P_x + a_2 P_{xxx} + a_3 P_x = 0. \quad (50)$$

An MI analysis is performed by looking for perturbed solutions of the following form:

$$P(x, t) = g + EH(x, t). \quad (51)$$

Inserting Eq. (51) into Eq. (50), then we get.

$$ED_{M,t}^{\sigma,n} H + Ea_1 (g + EH) H_x + Ea_2 H_{xxx} + Ea_3 H_x = 0.$$

And linearizing in  $E$ ,

$$ED_{M,t}^{\sigma,n} H + Ega_1 H_x + Ea_2 H_{xxx} + Ea_3 H_x = 0. \quad (52)$$

Let us consider the solution of Eq. (52) as:



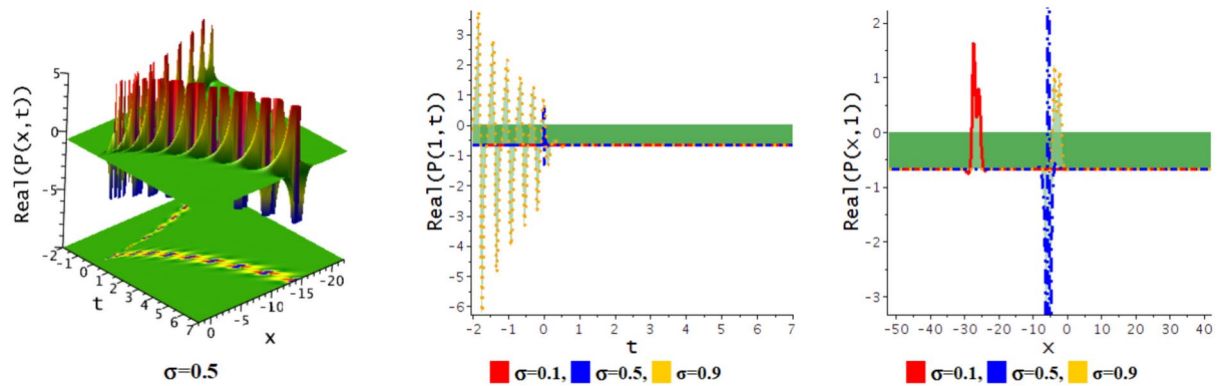


Fig. 12. Linked periodic rogue wave of the solution Eq. (41).

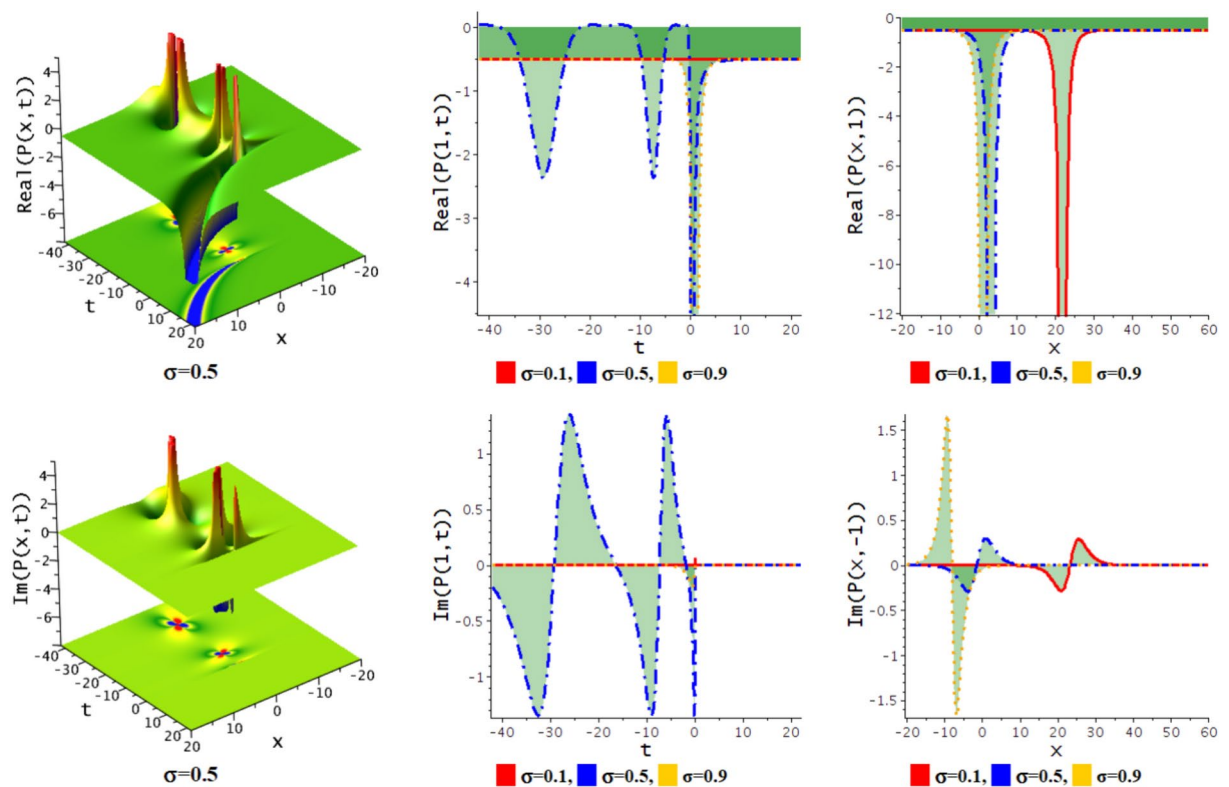


Fig. 13. Rogue wave of the solution Eq. (47).

$$P(x, t) = e^{i(\theta x + \epsilon \frac{\Gamma(n+1)}{\sigma} t^\sigma)}. \quad (53)$$

Inserting Eq. (53) into Eq. (52) and divide the entire equation by  $e^{i(\theta x + \epsilon \frac{\Gamma(n+1)}{\lambda} t^\lambda)}$ , then we get,

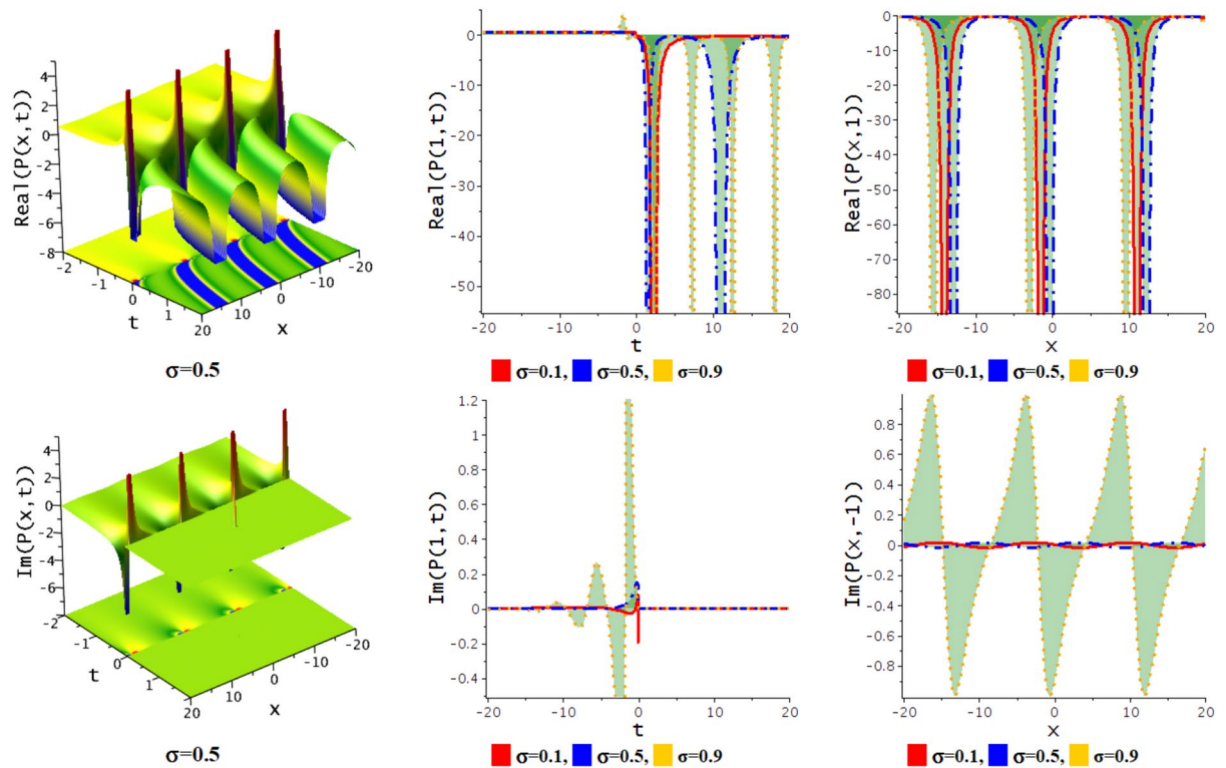
$$iE\epsilon + iEga_1\theta - iEa_2\theta^3 + iEa_3\theta = 0.$$

This equation defines the dispersion relation:

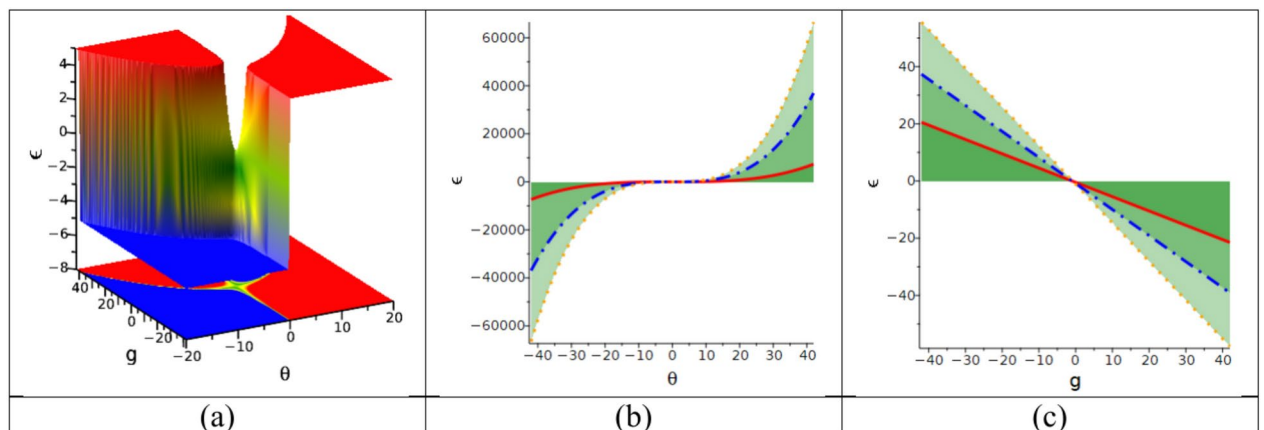
$$\epsilon = -(ga_1\theta - a_2\theta^3 + a_3\theta). \quad (54)$$

It is evident from Eq. (54) that the dispersion is stable and that, for negative values of  $\epsilon$ , any superposition of the solutions will seem to decay. Figure 15 depicts the 3-D and 2-D diagram of the Eq. (54).





**Fig. 14.** Interaction of periodic wave and kink wave of the solution Eq. (48).



**Fig. 15.** The 3-D and 2-D diagram of the Eq. (54) (a) 3-D diagram for  $a_1 = 1, a_2 = 0.1, a_3 = 1$ . (b) red  $[a_1 = 1, a_2 = 0.1, a_3 = 1, g = 0.5]$ , blue  $[a_1 = 1.3, a_2 = 0.5, a_3 = 1.5, g = 0.7]$ , orange  $[a_1 = 1.5, a_2 = 0.9, a_3 = 2, g = 0.9]$ , (c) red  $[a_1 = 1, a_2 = 0.1, a_3 = 1, \theta = 0.5]$ , blue  $[a_1 = 1.3, a_2 = 0.5, a_3 = 1.5, \theta = 0.7]$ , orange  $[a_1 = 1.5, a_2 = 0.9, a_3 = 2, \theta = 0.9]$ .

## Comparison, advantages and limitations

In this section, we perform both analytical and graphical comparisons between our work and the published results<sup>42</sup>. Additionally, we highlight some advantages and limitations of our applied method.

### Comparison with extended Tanh-method

In this subsection, we compare the solutions obtained in our study with those presented in<sup>32</sup>, which were derived using the extended tanh method and the generalized Kudryashov (GK) method. Using these methods, Sayed Saifullah et al.<sup>36</sup> investigated the fractional gPKdV equation and identified eight analytical soliton solutions. Their work revealed known phenomena, including bright and dark solitons, singular solutions, hyperbolic traveling wave solutions, and singular periodic solutions under specific parameter conditions. In contrast, our research employed the modified simple equation method to address the fractional gPKdV equation, yielding

thirty-four analytical soliton solutions. For selected parameter values, we uncovered phenomena such as double periodic waves, periodic lump-type waves, interactions between periodic lump and kink waves, bright bell-shaped waves, singular soliton waves, and V-shaped periodic rogue waves. Our approach proves effective and straightforward for identifying unique solitary wave solutions, introducing new phenomena for the M-fractional gPKdV equation.

### Advantages and limitation of the MSE technique

The primary advantage of the MSE (Modified Simple Equation) technique lies in its ability to derive soliton solutions for nonlinear partial differential equations (NLPDEs) in various forms, including hyperbolic, trigonometric, or exponential functions. Unlike other methods, this technique does not require auxiliary differential equations or predefined solutions. In contrast, methods such as the NK and improved F-expansion technique<sup>17</sup>, the extended direct algebraic approach<sup>18</sup>, the MSSE method<sup>19</sup>, the new mapping approach<sup>20</sup>, the modified extended tanh and NMK methods<sup>21</sup>, the unified approach<sup>22,23</sup>, and the exp-expansion and NMK techniques<sup>24</sup> rely on predefined auxiliary differential equations and predetermined solutions. Consequently, the MSE method uniquely allows for the direct resolution of NLPDEs. However, the scope of the MSE approach is limited; it cannot handle all types of NLPDEs, particularly when the balance number of an NLPDE exceeds two, making the solution process significantly more complex.

### Conclusion

This work applies the modified simple equation (MSE) method to the gPKdV equation with an M-fractional derivative, elucidating key properties of the M-fractional operator. By using the MSE technique, we explored ion-acoustic wave solutions in hyperbolic, exponential, and trigonometric forms. These solutions manifested as double periodic waves, periodic lump-type waves, interactions between periodic lump and kink waves, bright bell-shaped waves, singular soliton waves, V-shaped periodic rogue waves, and more for specific constraint values. We further analyzed the system's dynamic behavior through bifurcation and phase portrait studies. Additionally, we conducted a detailed modulation instability analysis, confirming the stability of the derived solutions in Fig. 15. To our knowledge, this approach to fractional nonlinear PDEs is novel. Consequently, our methods provide valuable tools for generating unique and precise soliton solutions, relevant for applications in nonlinear science and mathematical physics. In future work, we plan to investigate the chaotic dynamics and sensitivity analysis of the gPKdV equation and explore the efficacy of various fractional derivatives using a novel generalized approach.

### Data availability

All data generated or analyzed during this study are included in this article.

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## Author contributions

Mohammed Kbiri Alaoui: Configured the methodology and analyzed the figures. Md. Mamunur Roshid: Conceptualized this work, wrote and validated the codes, generated the figures, and wrote the main manuscript. Mahtab Uddin: Updated the result discussions and finalized the manuscript. Wen-Xiu Ma: Analyzed the findings and modified the manuscript. Harun-Or- Roshid: Reviewed the present works and wrote the initial manuscript. Mohammad Jahirul Haque Munshi: Investigated the work and wrote the initial manuscript. All authors have read and agreed to the published version of the manuscript.

## Competing interests

The authors declare no competing interests.

### Additional information

**Correspondence** and requests for materials should be addressed to M.M.R. or W.-X.M.

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