

High plating currents without dendrites at the interface between a lithium anode and solid electrolyte

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Avoiding lithium dendrites at the lithium/ceramic electrolyte interface and, as a result, avoiding cell short circuit when plating at practical current densities remains a significant challenge for all-solid-state batteries. Typically, values are limited to around 1 mA cm^{-2} , even, for example, for garnets with a relative density of $>99\%$. It is not obvious that simply densifying ceramic electrolytes will deliver high plating currents. Here we show that plating currents of 9 mA cm^{-2} can be achieved without dendrite formation, by densifying argyrodite, $\text{Li}_6\text{PS}_5\text{Cl}$, to 99% . Changes in the microstructure of $\text{Li}_6\text{PS}_5\text{Cl}$ on densification from 83 to 99% were determined by focused ion beam-scanning electron microscopy tomography and used to calculate their effect on the critical current density (CCD). Modelling shows that not all changes in microstructure with densification act to increase CCD. Whereas smaller pores and shorter cracks increase CCD, lower pore population and narrower cracks act to decrease CCD. Calculations show that the former changes dominate over the latter, predicating an overall increase in CCD, as observed experimentally.

All-solid-state batteries based on a lithium metal anode and ceramic electrolyte have the potential to achieve high energy densities^{1–3}. Avoiding lithium dendrites (filaments of lithium) that penetrate the solid electrolyte (SE) on charging (lithium plating) leading ultimately to a short circuit at practical current densities remains a substantial challenge^{4–22}. The need for fast charging in, for example, electric vehicles means that plating current densities in excess of 5 mA cm^{-2} are desirable^{23,24}. The major classes of solid electrolyte under consideration, garnets, sulfides and halides, struggle to achieve lithium plating currents in this range without dendrite penetration, short circuit and cell failure^{25–29}.

For example, extensive work on garnets, including lithium lanthanum zirconium oxide (LLZO), have shown that plating currents are typically limited to around 1 mA cm^{-2} even when the electrolyte is 99% dense or above, calling into question whether densifying polycrystalline ceramic electrolytes can achieve high plating currents^{30–32}.

Work to address the challenge of avoiding dendrites at high plating currents, that is achieving a high critical current density (CCD), has explored altering the Li/SE interface or modifying the solid electrolyte. At the Li/SE interface, notable strategies have included the addition of a thin interlayer and patterning of the solid electrolyte

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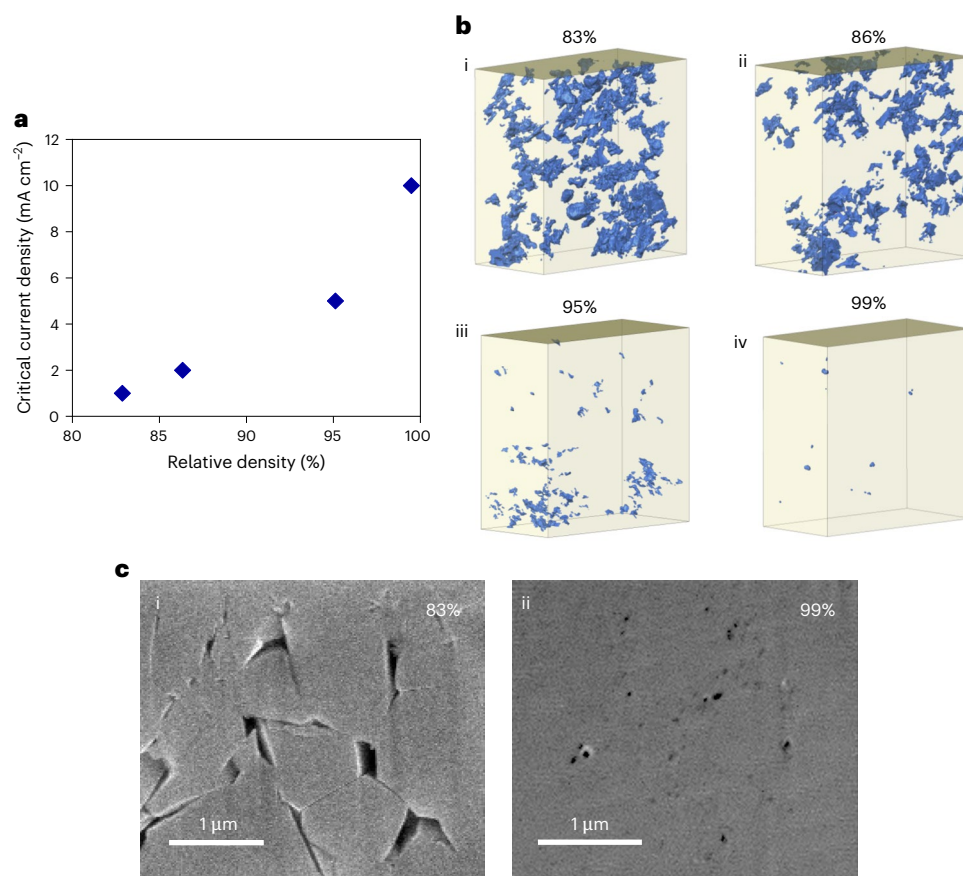


Fig. 1 | Critical currents for dendrite formation, with 3D reconstructions and FIB-SEM cross sections showing how CCDs increase with densification.

a, CCD for dendrite formation versus relative density of Li₆PS₅Cl solid electrolytes densified under different conditions: 83% (cold pressed), 86%, 95% and 99% spark plasma sintered at 300, 350 and 400 °C for 5 mins.

b, 3D reconstructed microstructures of Li₆PS₅Cl solid electrolytes obtained using FIB-SEM serial sectioning and imaging: (i) 83% dense, (ii) 86% dense, (iii) 95% dense, (iv) 99% dense. Regions of porosity are coloured blue. **c**, (i) and (ii) are magnified FIB-SEM cross sections of the 83% and 99% dense solid electrolytes in (i) and (iv) from panel **b**.

surface^{33–47}. Regarding the solid electrolyte, the addition of composite phases, bilayered electrolytes and the use of single crystals have all been explored^{30,48–54}. An important study by Krauskopf et al.⁵⁵ has demonstrated that high plating currents can be obtained at micro-electrodes placed within grains, that is, in defect-free regions, yet, as the authors point out, such currents remain elusive when using bulk Li anodes. In an attempt to avoid grain boundaries, glasses have been investigated^{52,56–59}. However, generally, glasses are found to be brittle, and despite avoiding grain boundaries, limiting imperfections may still remain^{56,60}. It has been reported that thin-film batteries based on LiPON can avoid dendrites when charged at fast rates, but their low conductivity means that they are not suitable for larger cells^{61,62}.

Sulfide-based solid electrolytes such as argyrodite, Li₆PS₅Cl, have high conductivities and have therefore received considerable attention as possible solid electrolytes for solid-state batteries⁶³. Here we investigate the effect of densification on the Li₆PS₅Cl microstructure and CCD. We show that plating currents as high as 9 mA cm⁻² can be achieved without lithium dendrites (and as a consequence without short circuit) by densifying the solid electrolyte to 99%. Extracting the changes in the microstructure on densification from focused ion beam-scanning electron microscopy (FIB-SEM) tomography and then using the changes in the microstructural dimensions to calculate CCD shows that the smaller pores and shorter cracks that occur on densification act to increase CCD, whereas narrower cracks and greater distance between pores (lower population density) act to reduce CCD. Therefore, not all the changes in the microstructure of Li₆PS₅Cl upon densification act to increase CCD. The calculations show that the reduction in pore size

and the shorter crack lengths outweigh the narrower cracks and lower population density, predicting higher CCDs overall with densification, which is confirmed by experiment. These high plating currents avoiding dendrites are achieved without the need to form alloys with lithium, without the need to modify the solid electrolyte surface or to introduce interlayers.

CCD and microstructure on densification

Li₆PS₅Cl polycrystalline solid electrolytes with a range of densities from 83% to 99% were prepared by sintering under different conditions. Powder X-ray diffraction patterns confirming the single-phase nature of the argyrodite after sintering are shown in Supplementary Fig. 1. The critical currents at which dendrites form (CCD) for the solid electrolytes of different density are shown in Fig. 1, along with 3D reconstructed microstructures obtained by FIB-SEM serial sectioning and imaging and FIB-SEM cross sections. The CCD measurements and FIB-SEM are described in Methods. The CCD increases from 1 mA cm⁻² for the 83% dense solid electrolyte to 10 mA cm⁻² for the material with 99% relative density. The CCD results are supported by the micro X-ray computed tomography (XCT) data in Fig. 2 discussed below, which show no dendritic cracks in the 99% dense solid electrolyte at 9 mA cm⁻² but a dendrite at 10 mA cm⁻². The lowest density solid electrolyte was formed by cold pressing the powder, which is the most widely used method for forming sulfide solid electrolytes^{5,64,65}. Higher densities were achieved by spark plasma sintering. The process involves passing a current through the die and solid electrolyte disc, resulting in more homogeneous and rapid heating than conventional methods. Densities

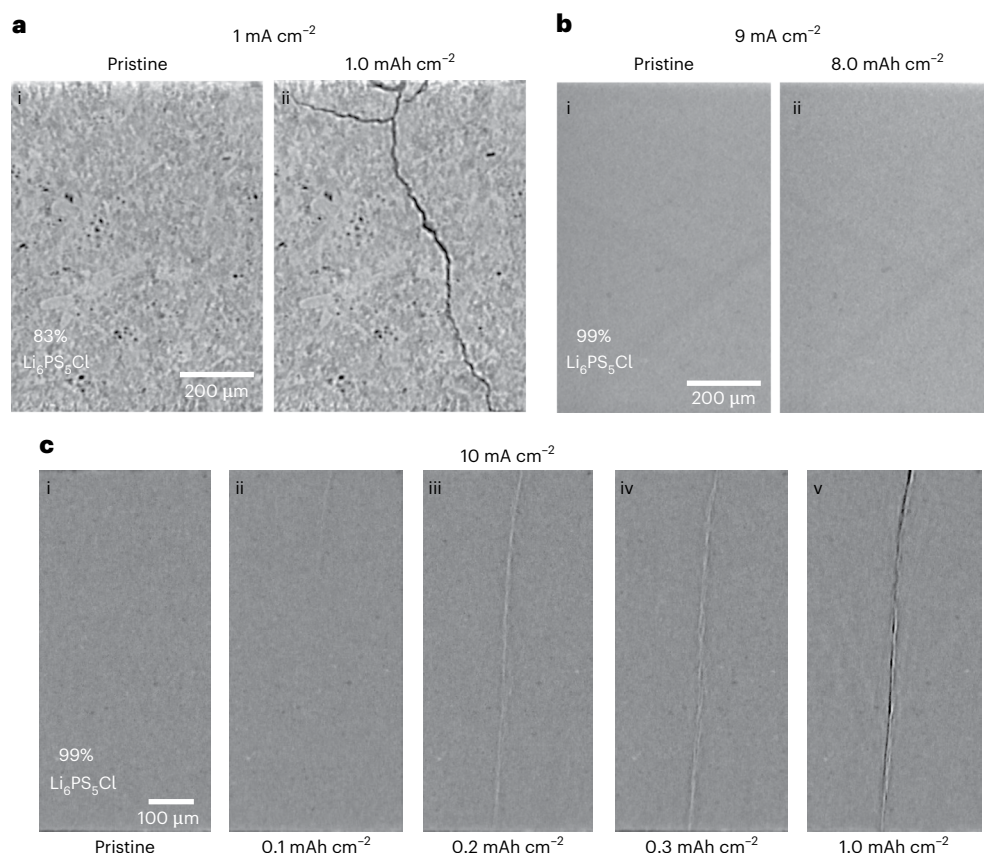


Fig. 2 | In situ XCT cross sections during plating in Li/Li₆PS₅Cl/Li cells with high- and low-density electrolytes show dendrite cracks at the CCD but not below. **a–c**, Virtual XCT cross sections of Li/Li₆PS₅Cl/Li cells with plating at the top electrode for a cold pressed Li₆PS₅Cl electrolyte (83% dense): pristine (i) and after plating 1 mAh cm⁻² at 1 mA cm⁻² (ii), (a); for a 400 °C sintered Li₆PS₅Cl electrolyte

(99% dense), pristine (i) and after plating 8 mAh cm⁻² at 9 mA cm⁻² (ii), (b), and the same density electrolyte with plating at 10 mA cm⁻², pristine (i), after plating 0.1 mAh cm⁻² (ii), 0.2 mAh cm⁻² (iii), 0.3 mAh cm⁻² (iv) and 1 mAh cm⁻² (v) (c). The Li electrodes have been removed from the images for clarity.

of 86, 95 and 99% were obtained by heating respectively at 300, 350 and 400 °C for 5 mins. CCDs were determined as described in Methods (Supplementary Figs. 2 and 3), which also contains more details of the sintering. Densities were determined by weighing the discs and establishing the volume of the electrolyte from XCT measurements.

The microstructure changes significantly between 83% and 99% dense solid electrolytes. The 3D reconstructions (Fig. 1b) show that the porosity decreases markedly with increasing densification. The SEM images in Fig. 1c show the presence of grains in the 83% dense material, whereas the 99% dense material exhibits greater homogeneity, making it harder to identify individual grains. Sulfides are mechanically soft. Their Young's moduli are typically in the range of 10–30 GPa (refs. 12,60,66,67) compared with oxides of 150–200 GPa (refs. 68–70). We have previously measured the Young's modulus of Li₆PS₅Cl to be as low as 28 GPa (ref. 12). Conventional thinking based on oxide ceramics may not be the best approach to understand these compliant materials, which densify to some extent even with cold pressing and to a very significant extent when heated at relatively low temperatures (several hundred degrees) for short durations.

Whereas the CCD measurements indicate that densifying the electrolyte suppresses dendrites, and as a consequence short circuits, it is instructive to examine with XCT what happens when lithium is plated. XCT virtual cross sections are presented in Fig. 2. The data were collected from symmetric two-electrode cells with lithium being plated at the electrode at the top of the images. The lithium electrodes have been removed from the XCT for clarity. Figure 2a shows the virtual cross section of the cell with the 83% dense solid electrolyte before and after plating a capacity of 1 mAh cm⁻² at a rate of 1 mA cm⁻².

The image after plating is very similar to those reported previously for lower-density argyrodite^{12,13}, with a spalled-out piece of the ceramic and the transverse crack reaching the other electrode. The current density at which this occurs is also consistent with previous work^{12,13,25}. In contrast, the cell with the 99% solid electrolyte plated at a current density of 9 mA cm⁻² shows no evidence of a dendritic crack (Fig. 2b). This is also demonstrated in the 3D reconstructions of the XCT data in Supplementary Fig. 4. Repeating the experiment on the 99% dense solid electrolyte but increasing the current density to 10 mA cm⁻² does result in a dendritic crack (Fig. 2c and Supplementary Fig. 4). The cracks are however different between the low- and high-density ceramics. In the case of the dense solid electrolyte there is no evidence of spallation, and the transverse crack follows a very straight path to the other electrode. The XCT virtual cross-sectional images in Fig. 2b,c for 99% dense solid electrolytes confirm that the CCD is between 9 and 10 mA cm⁻², in accord with the electrochemical measurements in Fig. 1 and Supplementary Figs. 2 and 3.

Microstructure determines critical current density

The experimental results in Figs. 1 and 2 show that sintering changes the Li₆PS₅Cl solid electrolyte microstructure and increases CCD. However, are the changes in microstructure with densification responsible for the high CCD? To explore how the changes in microstructure influence CCD, we need to extract the dimensions of the microstructure (pores and cracks) and how they change on densification, then use these to calculate the CCDs. For the latter, we employ the recently published model of dendrite penetration into polycrystalline ceramic electrolytes¹².

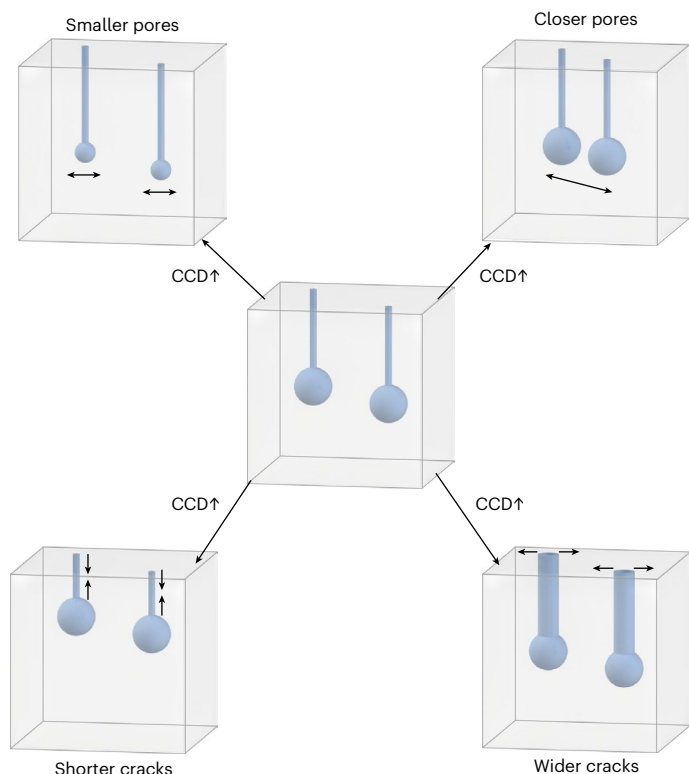


Fig. 3 | Schematic demonstrating key parameters affecting dendrite initiation. Reduction in pore size and pore separation, shorter and wider cracks lead to an increase in the CCD.

The dendrite model is based on the viscoplastic flow of lithium. On plating, Li deposits and grows along cracks in the ceramic from the anode to fill sub-surface pores. If the current density is sufficiently high that the Li being plated into the pores exceeds the Li flow out to the extent that pressure in the pores exceeds the local fracture strength at the grain boundaries, then a dendritic crack initiates. As illustrated in Fig. 3, the CCD for crack initiation depends on the pore size and population density (pore proximity) and the length and width of cracks connecting the pores to the interface¹². More details of the model can be found in the Supplementary Methods and Supplementary Fig. 5.

The dimensions of the microstructure pores and cracks for each of the densified $\text{Li}_6\text{PS}_5\text{Cl}$ solid electrolytes were determined by analysing the FIB-SEM tomographic images from the near surface region of each electrolyte, providing a resolution (~40 nm) which is below the limit typically available from XCT measurements. The method involves imaging a series of FIB-SEM sections allowing reconstruction of a 3D volume, which can be analysed. How the dimensions were obtained from the imaging for each microstructure is described further in the Methods. The microstructure of a polycrystalline ceramic contains a distribution of pore and crack sizes. Our previous work showed that it is the large pores within the distribution of pore sizes that are primarily responsible for dendrite formation¹². They are the weakest points as their larger surface area results in greater influx of lithium, hence greater pressure in the pore for a given overall current density. Considering the differently densified solid electrolytes, the largest pores range from 0.5 (high density) to 2.3 μm (low density) in radii. Our model does not consider an isolated pore but rather an array of pores distributed in three dimensions. The current at a given pore is affected by the size and proximity of the neighbouring pores¹². The neighbouring pores spread the current reducing it at each pore. When evaluating the proximity of pores, we consider only the large pores, that is, within a pore radius of 1/3rd. This is because we have shown that it is the large pores that take by far the most current. Pores that are less than 1/3rd of

the radius of the larger pores attract less than 10% of the current. This is shown in Supplementary Fig. 6, where we calculate the current distribution between pores and show that the central smaller pore, which is 1/3rd smaller in radius, experiences only 10% of the current. Because it is these larger pores that are experiencing and influencing the vast majority of the current, they are those critical to predicting dendrite crack initiation. When considering the proximity of the large pores to one another across different density electrolytes, the average pore proximity ranges from 10 (in low-density microstructures) to 22 μm (in high-density microstructures), as shown in Fig. 4a. The pore sizes and proximities from the microstructural analysis were used to define the ranges shown in Fig. 4a. Pores very close to the surface cannot build pressures to crack the ceramic as extrusion is too fast (easy), also the modelling shows deeper pores attract more current and so will build greater pressure (Supplementary Fig. 7). Therefore, we consider the longer cracks (deeper pores) for each density and their associated widths. The longest cracks measured within the differently densified solid electrolytes were found to range from 2.5 (low density) to 1.5 μm (high density). The average widths of the longest cracks (within 33% of the longest cracks) within each electrolyte was found to vary from 158 (low density) to 125 nm (high density). The crack lengths and widths were used to define the ranges in Fig. 4b.

Figure 5a–d illustrates how CCD varies with pore size and separation and crack length and width. The ranges of the microstructure dimensions used in Fig. 5 are obtained from Fig. 4, and thus represent typical values expected when densifying $\text{Li}_6\text{PS}_5\text{Cl}$. Looking at Fig. 5a,b, CCD decreases with increasing pore size and separation, both due to their effect on the magnitude of current experienced by a given pore. As shown in Fig. 5c,d, CCD increases with increasing crack width and decreases with increasing crack length, both due to the effect of width and length on the rate of Li extrusion. The dependence of CCD on a wider range of microstructural dimensions is shown in Supplementary Figs. 8 and 9, but the trends are the same as in Fig. 5. Figure 4 shows that densifying $\text{Li}_6\text{PS}_5\text{Cl}$ results in shorter and narrower cracks and smaller pores with greater separation. Therefore modelling shows that not all changes in microstructure during the densification of $\text{Li}_6\text{PS}_5\text{Cl}$ act to increase CCD. Whereas decreasing crack width and increasing pore separation acts to decrease CCD, at the same time, the decreasing crack length and pore size act to increase it. The calculations show (Fig. 5e) that the former outweigh the latter predicting an overall increase in CCD as observed experimentally (Fig. 1). Whereas the trends in the predicted and experimental CCDs are the same (Figs. 1a and 5e), we do not expect a close match between the absolute values; for example, even the larger pores for a given density are not of identical size. Also, the microstructure model assumes a simple geometry for pores (spheres) and cracks (cylinders) whereas actual pores and cracks will have more complex shapes. We used a combination of cold pressing and spark plasma sintering because no single method alone could provide a density range as wide as 83–99%, the alignment between experimental and modelled CCDs shows that our model is more general than just one densification method, and it shows that the CCD is indeed determined by the microstructure, however formed.

Stable lithium plating at high current density

The high CCD of 10 mA cm^{-2} for the 99% densified solid electrolyte (Fig. 1) enables extended cycling without short circuit at high, practical, plating current densities. This is illustrated in Fig. 6, where three electrode cells with a lithium working electrode are plated at a current density of 9 mA cm^{-2} . In the case of the 83% dense (cold pressed) $\text{Li}_6\text{PS}_5\text{Cl}$ solid electrolyte the cell short circuits after six cycles. Whereas the 99% dense $\text{Li}_6\text{PS}_5\text{Cl}$ solid electrolyte cycled continuously until the cycling was stopped after 100 cycles. In both cases stripping was carried out at the lower current density of 0.05 mA cm^{-2} to avoid voiding at the interface. It is worth noting that even the 99% dense $\text{Li}_6\text{PS}_5\text{Cl}$ discs retain surface roughness, as shown by the atomic force microscopy

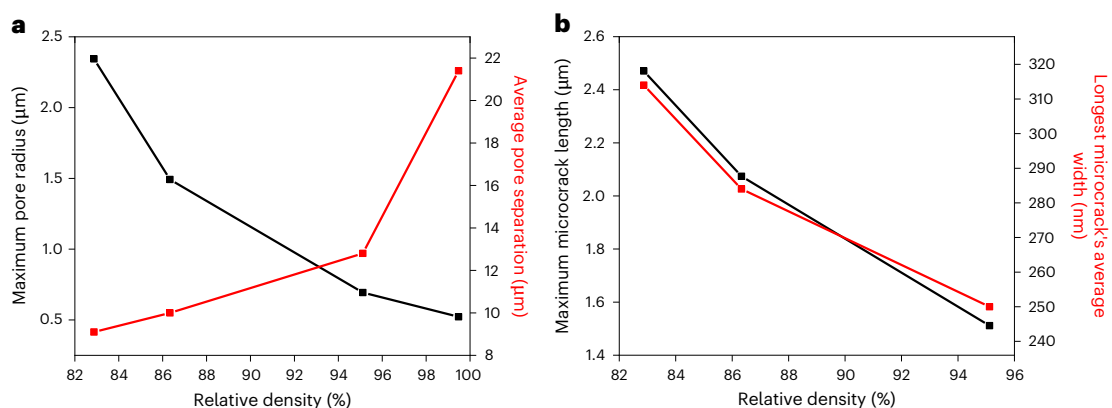


Fig. 4 | Plots of how pore and crack dimensions vary with density of $\text{Li}_6\text{PS}_4\text{Cl}$ constructed using values from FIB-SEM tomography. a, The largest pore size and their average separation. b, The maximum crack length and their associated width. Black data correspond to the left y axis and red data correspond to the right y axis.

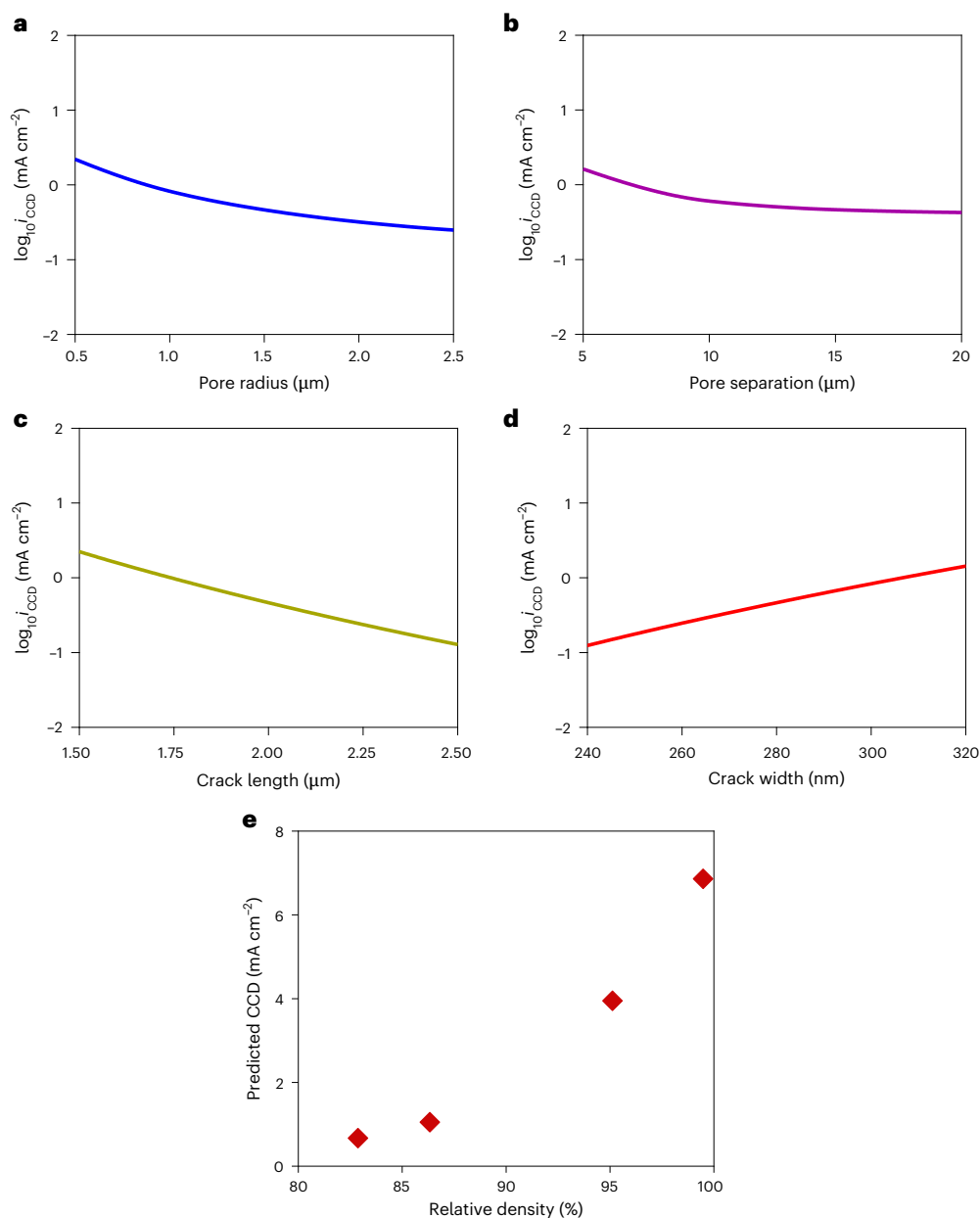


Fig. 5 | Effect of pore size and separation as well as crack length and width on the CCD calculated using the dendrite model. a–d, The effect on the CCD for dendrite initiation of pore radius (a), pore separation (b), crack length (c) and crack width (d). e, Plot of predicted overall CCDs for each density of $\text{Li}_6\text{PS}_4\text{Cl}$.

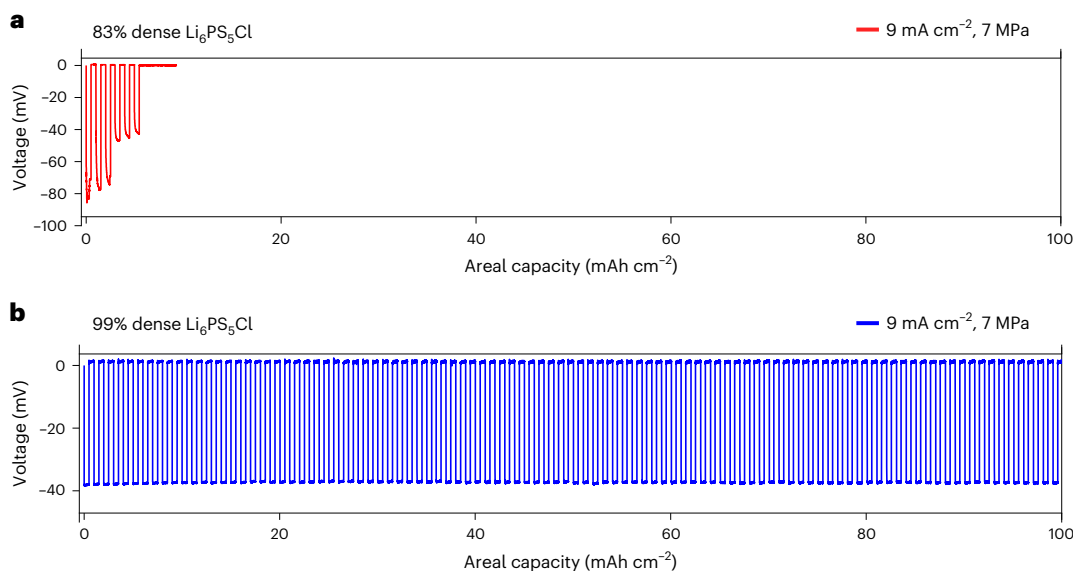


Fig. 6 | Repeated fast plating at 9 mA cm⁻² enabled by densifying the Li₆PS₅Cl microstructure. a, b, Cycling of Li/Li₆PS₅Cl 3-electrode cells with current densities of 9 mA cm⁻² for plating, 0.05 mA cm⁻² for stripping (to avoid void

formation), both for a capacity of 0.5 mAh cm⁻², under 7 MPa stack pressure, with a cold pressed Li₆PS₅Cl disk (83% dense) (a) and a 400 °C sintered Li₆PS₅Cl disk (99% dense) (b).

(AFM) scans in Supplementary Fig. 10, therefore the differences in CCD are not due to difference in surface roughness. A previous interesting study by Diallo et al.⁷¹ discusses how dendrites may form in low-density electrolytes where there is continuous and connected porosity. The study is complementary to our work which further explores the dendrite mechanism of the high-density regime where the porosity of the electrolyte is isolated, and cracks are initiated for the lithium to grow. In this case it is necessary to reach a critical current, CCD, to initiate dendrites. Together the studies offer important insights to the dendrite problem. As the focus of this Article is the effect of densification on dendrite suppression on plating and not on voiding, when stripping, we used a 7 MPa stack pressure along with the 0.05 mA cm⁻² stripping current mentioned above to prevent voiding formation. The sustained cyclability of the 99% dense solid electrolyte is not related to any suppression of dendrite propagation after initiation. As shown by the XCT images in Fig. 2, there are no dendrites at this plating current density.

It has been shown previously that Li₆PS₅Cl reacts with Li metal to form a solid electrolyte interphase. However, the growth slows with time and the layer stabilizes regardless of electrolyte density (Supplementary Fig. 11) and is invariant with cycling (Supplementary Fig. 12).

Conclusions

Sintering argyrodite, Li₆PS₅Cl, solid electrolytes to a density of 99% results in a critical current for lithium dendrites of 10 mA cm⁻². Plating lithium at the solid electrolyte interface at a current density of 9 mA cm⁻² shows no evidence of lithium dendrites electrochemically or in XCT imaging. The early motivation to develop ceramic electrolyte batteries was based on the belief that dense ceramics would not permit Li dendrites to form and short circuit cells. Yet, previous studies have suggested that dense ceramics, even with 99% density, do not guarantee Li plating at practical currents^{30,31}. Critical currents are often limited to a few mA cm⁻². The results here show that not all the changes in microstructure on densification necessarily favour increased CCDs. Densification decreases crack width and increasing pore separation, which act to decrease CCD, whereas the decreasing crack length and smaller pores act to increase CCD. For Li₆PS₅Cl, the latter outweighs the former. The high degree of densification is achieved here by spark plasma sintering. A future challenge is to achieve high degrees of densification by more scalable routes. It is known that spark plasma sintering

typically induces densification of ceramics with limited grain growth as compared to other techniques, and in general, different densification protocols are known to result in different microstructures of polycrystalline ceramics^{72–77}. The results reported here inspire investigation of other densification methods directed to optimizing the microstructural parameters that increase CCD, such as small pores, while suppressing those that act to decrease CCD, such as narrow cracks. Our results also highlight the need to study microstructure and CCD relations across a range of material systems. Future work focussed on scalable sintering of sulfide electrolytes may enable practical charging rates for solid-state cells employing lithium anodes without the need to form composite electrolytes or use interlayers.

Methods

Spark plasma sintering and electrolyte preparation

All procedures were performed in an Ar-filled glovebox (O₂ and H₂O levels < 1 ppm). Li₆PS₅Cl powder was purchased from Ampcera (scanning electron microscopy (SEM) of as-received powder shown in Supplementary Fig. 13), and measurements were performed to verify the absence of any carbon impurities (Supplementary Fig. 14). The powder was loaded into a graphite die set, which was then inserted into a spark plasma sintering system (FCT system GmbH) within an Ar-filled glovebox and subject to 50 MPa uniaxial pressure at either 300 °C, 350 °C or 400 °C for 5 min while under vacuum by passing a current through the die set. Five-mm-diameter sintered electrolyte discs of identical thickness (1 mm) were then used for cell assembly. Electrochemical impedance spectroscopy (EIS) measurements confirm an increase in conductivity with densification (Supplementary Figs. 15 and 16), consisted with lower grain boundary resistances on densification. Cold pressed electrolyte discs were formed by loading the same pristine electrolyte powder into a 5-mm-diameter stainless-steel die set and pressing under a uniaxial pressure of 400 MPa. To determine the relative density of the electrolytes, discs were weighed, and their total volume measured using micro-XCT. The densities calculated from the FIB-SEM tomography deviate within 2% of the bulk micro-XCT measurement. When using micro-XCT, the volume of the entire electrolyte disc is determined with a spatial resolution of ~1.6 μm and the weight of the disc measured to ±1 μg accuracy, and so its value is considered the most accurate.

Three-electrode cells

For a three-electrode cell, two disc-shaped lithium electrodes of 1-mm diameter were pressed onto the same side of the electrolyte disc as the working and reference electrodes, respectively. The remaining areas of the electrolyte surface were then covered by insulating polymer tape. A 5-mm lithium disc was pressed onto the opposite side as the counter electrode. The assembly was placed into a pouch cell and sealed under vacuum. A stack pressure of 7 MPa was applied to all cells in this study. The CCD is strictly defined as the current density at which dendrites form and below which they do not. Commonly implemented ‘rate-test’ protocols using symmetric cells with two Li electrodes can suffer from a number of problems including void formation at the interface^{31,78}. Here we determine CCD using three-electrode cells to avoid interference from the second electrode. By using asymmetrical three-electrode cells, we can further mitigate any void formation at the counter, which may influence the CCD measurement in symmetrical cells. To enable a more statistical conclusion, three tests at each current were repeated at the same relative density; currents below the CCD showed no indications of dendrite failure whereas those at the CCD always failed. Galvanostatic single plating for a capacity of 8 mAh cm⁻² was used. For cycling studies, cells were cycled galvanostatically with a capacity on each half cycle of 0.5 mAh cm⁻². Although the plating current on the working electrode was 9.0 mA cm⁻², a current density of 0.05 mA cm⁻² was used on stripping to avoid void formation at the working electrode–solid electrolyte interface (Supplementary Fig. 17). The 1-mm diameter working electrode compared with the 5-mm diameter counter can lead to some current focussing at the working electrode edge. This would mean that the CCD could be even higher than observed here. The 1-mm diameter working electrode still represents a sampling of thousands of pores and cracks, even for high-density electrolytes. All galvanostatic measurements were performed using a Gamry Interface 1010 potentiostat.

Potentiostatic EIS

Potentiostatic electrochemical impedance spectroscopy (PEIS) was taken using a 5-mV voltage perturbation, collecting ten points per decade across a frequency range of 1 MHz–1 Hz. PEIS data were analysed and fitted by equivalent circuit models using the ZView software package.

In situ X-ray tomography

X-ray tomograms were recorded at the I13-2 beamline at the Diamond Light Source and the TOMCAT beamline at the Swiss Light Source (Paul Scherrer Institute). For both experiments the projections were collected with a PCO pco.edge 5.5 sCMOS camera combined with an optical microscope with a 4× magnification, resulting in a pixel size of 1.63 µm. For each tomogram, 1,801 equiangularly distributed projections were taken over 180°. Additionally, X-ray tomograms were collected using a Zeiss Xradia 610 Versa X-ray microscope and 3,201 equiangularly distributed projections were taken over 360° for each tomogram, with a resultant pixel size of 1.63 µm. For the in situ measurements, metallic lithium foil was punched into disc-shaped lithium electrodes and pressed onto both sides of the electrolyte disc, 1-mm diameter electrodes were used for the working electrode and 3 mm for the counter. X-ray tomograms were obtained at the pristine state and after charging in successive steps of 0.1 mAh cm⁻². During the experiment, 7 MPa of pressure was applied using a constantly applied spring force, which was measured and calibrated using a transducer.

Plasma FIB-SEM and microstructure analysis

Solid electrolyte discs were transferred into the chamber of a plasma FIB scanning electron microscope (Helios G4 PFIB DualBeam, ThermoFisher Scientific) for FIB milling with an iLoad vacuum sample transfer system without exposure to air. Automated serial sectioning

was used to obtain a three-dimensional depiction of the porosity. A 30 kV and 15 nA Xe⁺ beam was used to slice the volume of interest and capture a stack of secondary electron cross-sectional images. Each slice had a thickness of 100 nm. The freshly milled cross section was then imaged with the SEM to visualize the pores and cracks in proximity of the pellet surface. To construct a three-dimensional view, the sequence of 2D secondary electron images were aligned using Avizo 3D (version 2021.2, ThermoFisher Scientific). Example 2D secondary electron images used to generate the 3D microstructures in Fig. 1b are shown in Supplementary Fig. 18. Using the Avizo 3D software the images were subjected to a semi-automated segmentation process to differentiate between pores, cracks and bulk electrolyte. In our microstructural analysis, we only considered features with a 3:1 length:width aspect ratio or higher as cracks. The dimensions of the largest pores and longest cracks were obtained from the segmentation analysis and used as inputs for the modelling simulations.

Powder XRD

The electrolyte discs were ground into powders using a pestle and mortar. X-ray diffraction was performed using a Rigaku MiniFlex 600 instrument located within a nitrogen glovebox.

Data availability

The data supporting the findings of this study are available within the article and its Supplementary Information files.

Code availability

The code used for the modelling in this study is available via Zenodo at <https://doi.org/10.5281/zenodo.10523607> (ref. 79).

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D.L.R.M. contributed to all aspects of the research. D.L.R.M., D.S.-J., B.H. S.Z., Z.N., A.B., J.I. and S.M. carried out the in situ synchrotron XCT. D.L.R.M. performed the preparation of electrolyte discs and cell assembly. D.L.R.M., Z.N. and J.B. optimized the sintering conditions for the electrolyte. D.L.R.M. and Z.N. performed the powder X-ray diffraction. G.J.R. performed the Raman spectroscopy. M.S. and D.L.R.M. performed the plasma FIB imaging. D.L.R.M., G.L., Z.N. and C.W.M. developed the model. G.L. performed and implemented the modelling. D.L.R.M., G.L., D.S.J., M.S., P.S.G., T.J.M. and P.G.B. discussed the data. All authors contributed to the interpretation of data. D.L.R.M., G.L., D.S.J. and P.G.B. wrote the paper with contributions and revisions from all authors. The project was supervised by D.S.J. and P.G.B.

Competing interests

The authors declare no competing interests.

Additional information

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