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## Pyrocurrent anomalies and intrinsic magnetodielectric behavior near room temperature in $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , a compound with distorted honeycomb and spin-chains

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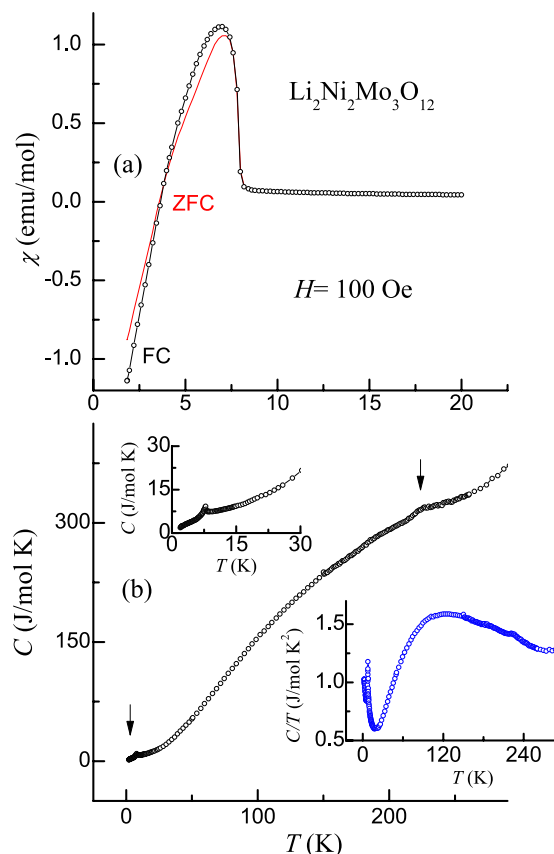
Keeping current interests to identify materials with intrinsic magnetodielectric behaviour near room temperature and with novel pyroelectric current anomalies, we report temperature and magnetic-field dependent behavior of complex dielectric permittivity and pyroelectric current for an oxide,  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , containing magnetic ions with (distorted) honey-comb and chain arrangement and ordering magnetically below 8 K. The dielectric data reveal the existence of relaxor ferroelectricity behaviour in the range 160–240 K and there are corresponding Raman mode anomalies as well in this temperature range. Pyrocurrent behavior is also consistent with this interpretation, with the pyrocurrent peak-temperature interestingly correlating with the poling temperature.  $^7\text{Li}$  NMR offer an evidence for crystallographic disorder intrinsic to this compound and we therefore conclude that such a disorder is apparently responsible for the randomness of local electric field leading to relaxor ferroelectric property. Another observation of emphasis is that there is a notable decrease in the dielectric constant with the application of magnetic field to the tune of about  $-2.4\%$  at 300 K, with the magnitude varying marginally with temperature. Small loss factor values validate the intrinsic behaviour of the magnetodielectric effect at room temperature.

The search for materials exhibiting large magnetodielectric (MDE) coupling in a broad temperature range, particularly around room temperature, continues unabated in the literature. It is of great interest to identify such materials with this property at room temperature to facilitate applications [see, for instance, refs 1–5]. It is now well-known that, around room temperature, many insulating materials become sufficiently conducting to result in detectable magnetoresistance. Under such circumstances, the debate of intrinsic versus extrinsic behavior of magneto-capacitance often arises<sup>6</sup>. In this article, keeping such trends of research activities in mind, we present the results of our investigation in this direction on a Li-based compound,  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , which was not paid any attention in the literature after initial magnetic studies by Hase *et al.*<sup>7</sup>.

In the compound,  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , Ni has been shown to possess magnetic moment (spin = 1), undergoing a complex magnetic order below  $T_C = 8$  K [ref. 7]. This system consists of linear chains as well as distorted honey-comb of the divalent Ni ions. Thus, there are two independent sites for Ni ions. Through neutron diffraction measurements, it has been found that the Ni chains exhibit ferromagnetic order, while distorted honeycomb of Ni shows the antiferromagnetic structure. The nominal composition of the title compound considering disorder<sup>8</sup> can be written as  $\text{Li}(\text{Li}_{0.5}\text{Ni}_{1.5})(\text{Li}_{0.5}\text{Ni}_{0.5})\text{Mo}_3\text{O}_{12}$ . The crystal structure [orthorhombic, space group *Pmcn* (*No.* 62)] in which the compound forms is related to that of  $\text{NaCo}_{2.31}\text{Mo}_3\text{O}_{12}$  [ref. 9]. The detailed crystal structure and magnetic structure of the title compound can be found in the references 8 and 7 respectively.

It is important to note that the magnetic properties of compounds with honey-comb arrangement [see, for instance, refs 10–14] of atoms and/or spin-chain systems [see, for instance, references 15–19] have been attracting a lot of attention experimentally as well as theoretically. However, there are only a few reports in the literature,

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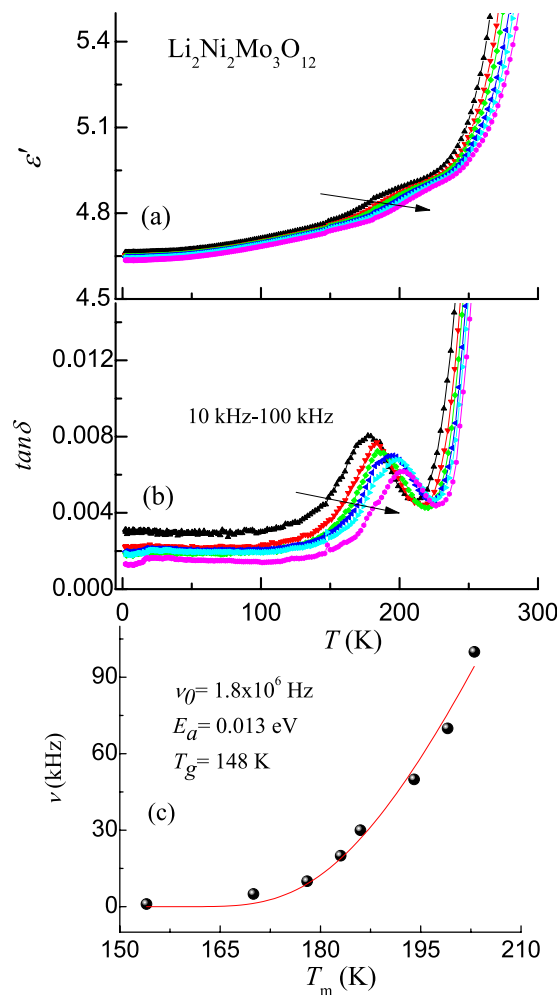
**Figure 1.** (a) Magnetic susceptibility obtained in 100 Oe for zero-field-cooled and field-cooled conditions of the specimen; (b) heat-capacity as a function of temperature for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  in the absence of magnetic field; vertical arrows mark the features of interest. Top and bottom insets in (b) show the  $C$  and  $C/T$  plots to highlight the features.

addressing the dielectric and magneto-dielectric behavior of honey-comb and spin-chain-based oxides [see, for instance, refs 15–19]. Since the  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  was reported to be an insulator, interestingly containing (distorted) honeycomb lattice and spin-chain features<sup>7</sup>, we considered it worthwhile to explore its detailed dielectric and pyroelectric properties. With this motivation, we have carried out dielectric, magnetodielectric and pyrocurrent ( $I$ ) studies as a function of temperature ( $T$ ) on this compound. In order to supplement the arguments, we have also carried out magnetization ( $M$ ), heat-capacity ( $C$ ), Raman spectroscopy and  $^7\text{Li}$  (nuclear spin = 3/2) NMR studies.

## Results

**X-ray diffraction.** The x-ray diffraction (XRD) pattern is shown in Supplementary Information (see Supplementary Fig. S1). Within the detection limit (<2%) of the XRD measurement, the polycrystalline sample was found to be single phase ( $Pmcn$  space group). The experimental data was subjected to Rietveld refinement by using Fullprof programme<sup>20</sup> and the Wyckoff positions have been taken from ref. 8. The lattice parameters ( $a = 10.4118(4)$  Å,  $b = 17.4958(4)$  Å,  $c = 5.0778(4)$  Å) are in good agreement with the literature<sup>8</sup>.

**Magnetization.** Since dc magnetization ( $M$ ) behavior has been reported in ref. 7, we present here essential features from our data for the sake of completeness. We obtained  $T$ -dependence of dc magnetic susceptibility ( $\chi$ ) in the presence of various magnetic fields ( $H = 100$  Oe and 5 kOe). Our results are in good agreement with those reported [ref. 7]. The readers may see Supplementary Information for the  $\chi(T)$  behavior obtained in 5 kOe (see Supplementary Fig. S2(a)) for the zero-field-cooled (ZFC) condition (from 300 K to 2 K) of the specimen. The curve was fitted with the Curie-Weiss law in the interval of 50–300 K. The value of the effective magnetic moment obtained from this fitting is about  $3.23 \mu_B$  per Ni atom, which is higher than that expected ( $2.83 \mu_B$ ) for divalent Ni ( $S = 1$ ). We believe that such a marginally higher value could be the result of short-range magnetic correlations extending to such a high temperature range and intrinsic crystallographic disorder. The paramagnetic Curie-temperature, if inferred from this linear region, is about  $-40$  K, and the negative sign indicates the dominance of antiferromagnetic correlations. These values agree with those reported in the literature<sup>7</sup>. Fig. 1a shows the ZFC and field-cooled (FC)  $\chi$  curves obtained in 100 Oe. ZFC and FC curves exhibit a sudden upturn at 8 K with a lowering of temperature, due to the onset of magnetic ordering, followed by a peak around 7.2 K and finally a decrease. On further lowering of temperature, the sign of magnetization gets reversed (i.e. negative magnetization) around 3.6 K – the compensation temperature. ZFC and FC curves deviate only marginally from



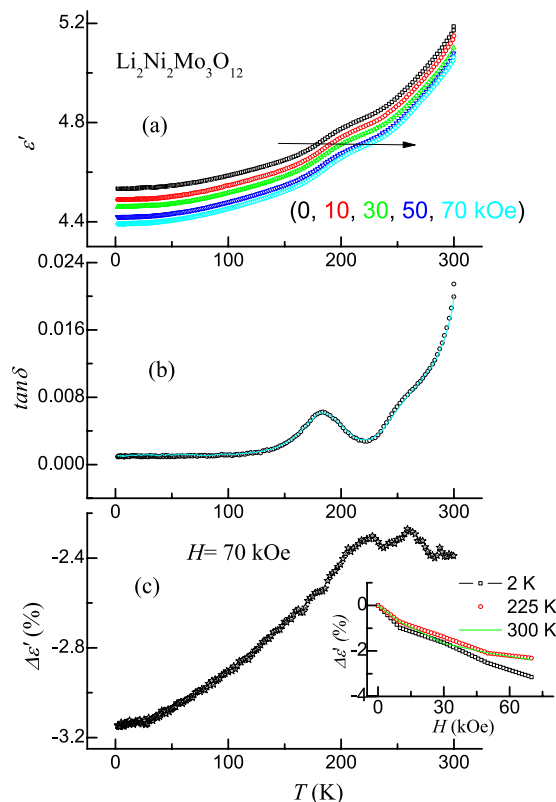
**Figure 2.** (a) Real part of the permittivity and (b) loss as a function of temperature with varying frequency for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ . Arrows are drawn to show that the curves shift to higher temperature range as the frequency is increased (shown for 10, 20, 30, 50, 70 and 100 kHz). In (c), Vogel-Fulcher fitting for the  $\tan \delta$  peak values is shown.

each other below 8 K, with a sign-crossover around the compensation temperature (3.6 K). From the absence of bifurcation of ZFC-FC curves, we conclude that this compound is not a spin-glass, despite the existence of crystallographic disorder. We carried out ac  $\chi$  studies and there is a frequency ( $\nu$ ) independent peak in the real part (not shown here) at the onset of magnetic order, without any feature in the imaginary part; this establishes the absence of any kind of magnetic frustration.

We have also performed isothermal magnetization  $M(H)$  studies up to 140 kOe (see Supplementary Fig. S2) at selected temperatures (1.8, 6, 10 and 20 K). At 1.8 K, the virgin curve lies outside the envelope curve and starts with the negative value (see top inset of Fig. S2(b)). This finding is agreeing with those reported for 1.7 K by Hase *et al.*<sup>7</sup> who performed the studies up to 50 kOe. We would like to stress that an extension of isothermal  $M$  studies to higher fields to 140 kOe in this investigation reveals another interesting feature - that is, a weak hysteresis, in the range 25 to 50 kOe in the curve for 1.8 K (see Fig. S2(c)). This finding implies the existence of disorder-broadened metamagnetic transition, which could be responsible for the virgin curve lying outside the envelope curve. Additionally, in the  $M(H)$  curve at 6 K, a small hysteresis is observed in the low-field region (<2 kOe) (see Fig. S2(d)). Isothermal  $M(H)$  plots continue to increase without any evidence for saturation at high fields in the magnetically ordered state. These findings, indicating the existence of both ferromagnetic and anti-ferromagnetic components, are consistent with the complex magnetic structure proposed in ref. 7.

**Heat capacity.** Fig. 1b shows the plot of  $C(T)$  below 290 K and two features are worth noting - one broad feature around 225 K, and a sharper one near 8 K (see also top inset of Fig. 1b). These are seen in  $C/T$  plots also (see bottom inset of Fig. 1b). The 8K-peak arises from long range magnetic ordering<sup>7</sup>. The high temperature peak, not reported earlier<sup>7</sup>, may be correlated to relaxor ferroelectric behavior, reported in this article.

**Complex permittivity.** Complex permittivity studies have been performed from room temperature to 2 K. Dielectric constant ( $\epsilon'$ ) and the loss factor ( $\tan \delta$ ) as a function of temperature, obtained with various frequencies (10, 20, 30, 50, 70 and 100 kHz), are shown in Fig. 2a and b. The observation we would like to stress is that  $\epsilon'$

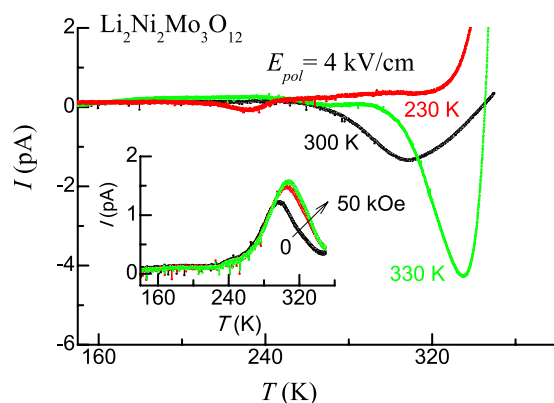


**Figure 3.** (a) Dielectric constant and (b)  $\tan\delta$  in zero field and in various magnetic fields (10, 30, 50 and 70 kOe), measured with 70 kHz, for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ . In (b), data points correspond to zero field, while in-field behavior (70 kOe) is represented by a continuous line. Both these curves overlap. (c) Magnetodielectric effect ( $\Delta\epsilon'$ ) for a change of magnetic field from zero to 70 kOe as extracted from the data shown in (a). Inset of (c) shows  $\Delta\epsilon'$  versus  $H$  for some selected temperatures.

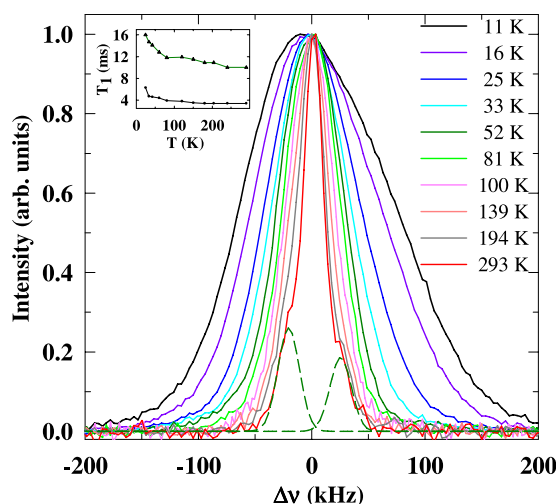
undergo a gradual increase with  $T$  from 2 K, with a broad shoulder, say, for example at 215 K for 100 kHz, and the shoulder temperature is frequency dependent. The feature is spread over a range of  $\sim 160$ –240 K. Corresponding feature, appearing in the form of a distinct peak, is clearer in the  $\tan\delta$  data; one can see a systematic variation of this peak temperature with frequency, and this dispersion indicates relaxor-type ferroelectricity in this compound. We have fitted the value of the peak temperature of  $\tan\delta$  with Vogel-Fulcher (V-F) equation,  $\nu = \nu_0 \exp[-E_a/k_B(T_m - T_g)]$ , where  $E_a$  is the activation energy,  $k_B$  is the Boltzmann constant,  $\nu_0$  is the pre-exponential factor,  $T_m$  is the temperature at which  $\tan\delta$  peaks and  $T_g$  is the freezing temperature. This fitting (see Fig. 2c) resulted in the value of  $E_a = 0.013$  eV,  $\nu_0 = 1.8 \times 10^6$  Hz and  $T_g = 148$  K, which agrees well with the literature<sup>21,22</sup>. We would like to mention here that, in the temperature range of above interest, the  $\tan\delta$  values are in the range of 0.001–0.01, which establishes highly insulating behaviour and negligible influence of extrinsic contributions.

Dielectric constant and the loss factor ( $\tan\delta$ ) measured with 70 kHz are shown below 300 K in Fig. 3a and b respectively in the presence of various magnetic fields (0, 10, 30, 50 and 70 kOe). It is clear from this figure that there is a systematic decrease in the dielectric constant as the magnetic field is increased. However, interestingly, there is no observable change in the  $\tan\delta$  values (see Fig. 3b) with the application of  $H$ ; low values even in high fields establish intrinsic nature of magneto-dielectric behaviour in this compound. Fig. 3c shows the MDE (derived from the curves obtained with 70 kHz) in the form of  $\Delta\epsilon'$ , where  $\Delta\epsilon' = [(\epsilon'(H) - \epsilon'(0))/\epsilon'(0)]$ . The values of  $\epsilon'(H)$  at 2 K and 300 K are  $-3.15$  and  $-2.4\%$  respectively. For some fixed temperatures (2 K, 225 K and 300 K),  $\Delta\epsilon'$  measured as a function of  $H$  and the curves thus obtained are presented in the inset of Fig. 3c. It is interesting that intrinsic and significant MDE is observed in a broad temperature range, even at room temperature.

**Pyrocurrent behaviour.** We did detailed pyroelectric current measurements for various protocols (temperature, magnetic field, rate dependence etc.). We first poled from 300 K to 2 K with the applied electric field of  $+4$  kV/cm. A peak is observed around 300 K for the rate of warming ( $dT/dt = 2$  K/min) (see Fig. 4). It is found that the direction of polarization gets reversed, if one changes the direction of the poling electric field (Fig. 4, inset). We did additional measurements as discussed in refs 23–25, mainly, varying the poling temperature for the same electric field and then recording the corresponding pyrocurrent behavior as a function of temperature. We obtained  $I$  vs  $T$  curves by poling from 330 K, 300 K and 230 K (to 100 K) with the electric field of 4 kV/cm. We noted (see Fig. 4) that the peak appears essentially at the respective poling temperatures; the magnitude at the peak also increases with the poling temperature. Such a behavior, rarely known in the literature, has been



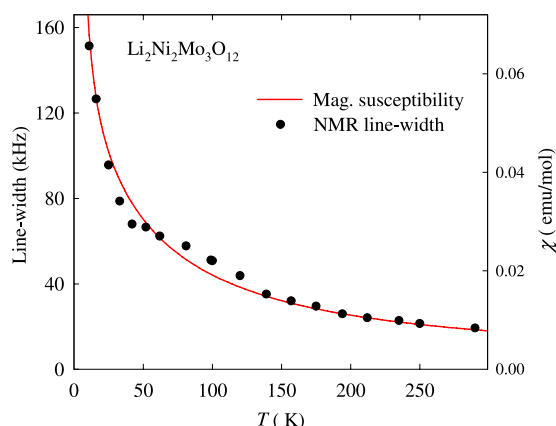
**Figure 4.** Pyrocurrent behaviour obtained with a poling field of 4 kV/cm for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  for poling temperatures of 230, 300 and 330 K. Inset shows pyrocurrent behavior in the presence of magnetic fields at 300 K for the poling electric field of  $-4\text{ kV/cm}$  and an arrow is drawn to show that the curves move towards higher temperature range with increasing magnetic field with the curves corresponding to 0, 30 and 50 kOe.



**Figure 5.** Temperature dependence of the  $^7\text{Li}$  NMR spectra for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  obtained at  $(\nu =) 27\text{ MHz}$ . The origin of the x-axis  $\Delta\nu = 0$  corresponds to the centre of the spectra at 293 K. The dotted green lines represent the asymmetric satellites at 293 K. The inset shows the  $T_1$  variation with temperature. The solid line is guide to eyes.

demonstrated for relaxor ferroelectric state in a delafossite<sup>23</sup>. We therefore attribute such a feature to the role played by polar nanoregions<sup>22</sup>. We have also obtained the pyropeak for different rates of heating, 2, 3, 4 and 5 K/min at 300 K and we found that there is an observable rate dependence (see Supplementary Fig. S3). This implies that, at room temperature, thermal stimulation effect plays a role. Interestingly enough, in the presence of various magnetic fields (10, 30 and 50 kOe), the pyro peak at this temperature arising from such an extrinsic effect shifts to higher temperature side with the increase in the magnitude as shown in inset of Fig. 4.

**$^7\text{Li}$  NMR.** NMR measurements were carried out on powdered sample at 27 MHz and the spectrum was obtained by Fourier transformation. A single Gaussian line was obtained at 293 K along with asymmetric quadrupole satellites (see dotted curves in Fig. 5). The quadrupolar coupling constant turns out to be about 92 kHz at 290 K. The full width at half maximum ( $\Delta\nu$ ) of the spectral line is found to be 22 kHz at 290 K. We found that the Knight shift is negligible down to 20 K. Many lithium based systems show a similar behaviour due to weak coupling with the neighbouring nuclei. However, the line-width  $\Delta\nu$  is found to vary drastically with temperature (Fig. 6) and it has a close resemblance to the variation of magnetic susceptibility. This is a strong indication that the  $^7\text{Li}$  NMR line-width and macroscopic magnetization have a common origin. It is notable that, below 20 K, there is a marked asymmetry in the spectra with respect to Gaussian line attributable to resonances with different widths from different Li sites (experiencing different hyperfine interactions depending on local magnetic environment, as the magnetically ordered regime is approached). Strictly speaking, in the absence of crystallographic disorder, there is only one site for Li and therefore this asymmetry supports the conclusion on the presence of anti-site Li ions in this compound<sup>7</sup>. The magnetic Ni ion has two sites forming distorted honeycomb lattice and linear chains with different types of magnetic interactions as revealed by neutron diffraction and some fraction of



**Figure 6.** The variation of magnetic susceptibility and NMR line-width as a function of temperature for  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ .

Li occupies both these sites as briefed in the introduction. The spectra well below 8 K is broadened considerably due to the transferred hyperfine field, signifying the onset of long range magnetic order. Spin-lattice relaxation time ( $T_1$ ) behavior also is consistent with multiple Li sites. At least two distinct components, one with a  $T_1$  of 10 ms and another with a  $T_1$  of 3.5 ms, could be resolved at 290 K. As the temperature is decreased both the components remain distinct, and the magnitude of  $T_1$  increases monotonously to 16.3 ms and 6.3 ms at 24 K. (Fig. 6, inset). The ratio of these two components remains at about 75:25 throughout the temperature range of measurement. In short, NMR thus provides evidence for crystallographic disorder, which must be the source of nanopolar behavior.

**Raman spectroscopy.** The Raman spectrum at room temperature is shown in the top panel of Fig. 7. The spectral features at  $332\text{ cm}^{-1}$  and  $365\text{ cm}^{-1}$  correspond to the bending modes of the  $\text{MO}_4$  tetrahedra while those at  $795\text{ cm}^{-1}$ ,  $822\text{ cm}^{-1}$ ,  $890\text{ cm}^{-1}$ ,  $933\text{ cm}^{-1}$  and  $962\text{ cm}^{-1}$  correspond to the stretching modes of the  $\text{MO}_4$  tetrahedra. These modes are concomitant to the bending modes ( $326\text{ cm}^{-1}$  and  $370\text{ cm}^{-1}$ ) and stretching bands ( $823\text{ cm}^{-1}$ ,  $853\text{ cm}^{-1}$ ,  $910\text{ cm}^{-1}$ ,  $974\text{ cm}^{-1}$ ,  $995\text{ cm}^{-1}$ ) of the isomorphous structure  $\text{Li}_2\text{Mg}_2\text{Mo}_3\text{O}_{12}$  [ref. 26]. The entire spectrum was fitted with this assignment of the peaks. The temperature dependencies of the peak position of strongest spectral features  $\nu_1$  ( $365\text{ cm}^{-1}$ ),  $\nu_2$  ( $890\text{ cm}^{-1}$ ),  $\nu_3$  ( $933\text{ cm}^{-1}$ ),  $\nu_4$  ( $962\text{ cm}^{-1}$ ) are also shown in Fig. 7. All of these spectral features vary non-monotonically in the neighborhood of 180 K and 10 K as demonstrated in the right and left panels respectively.

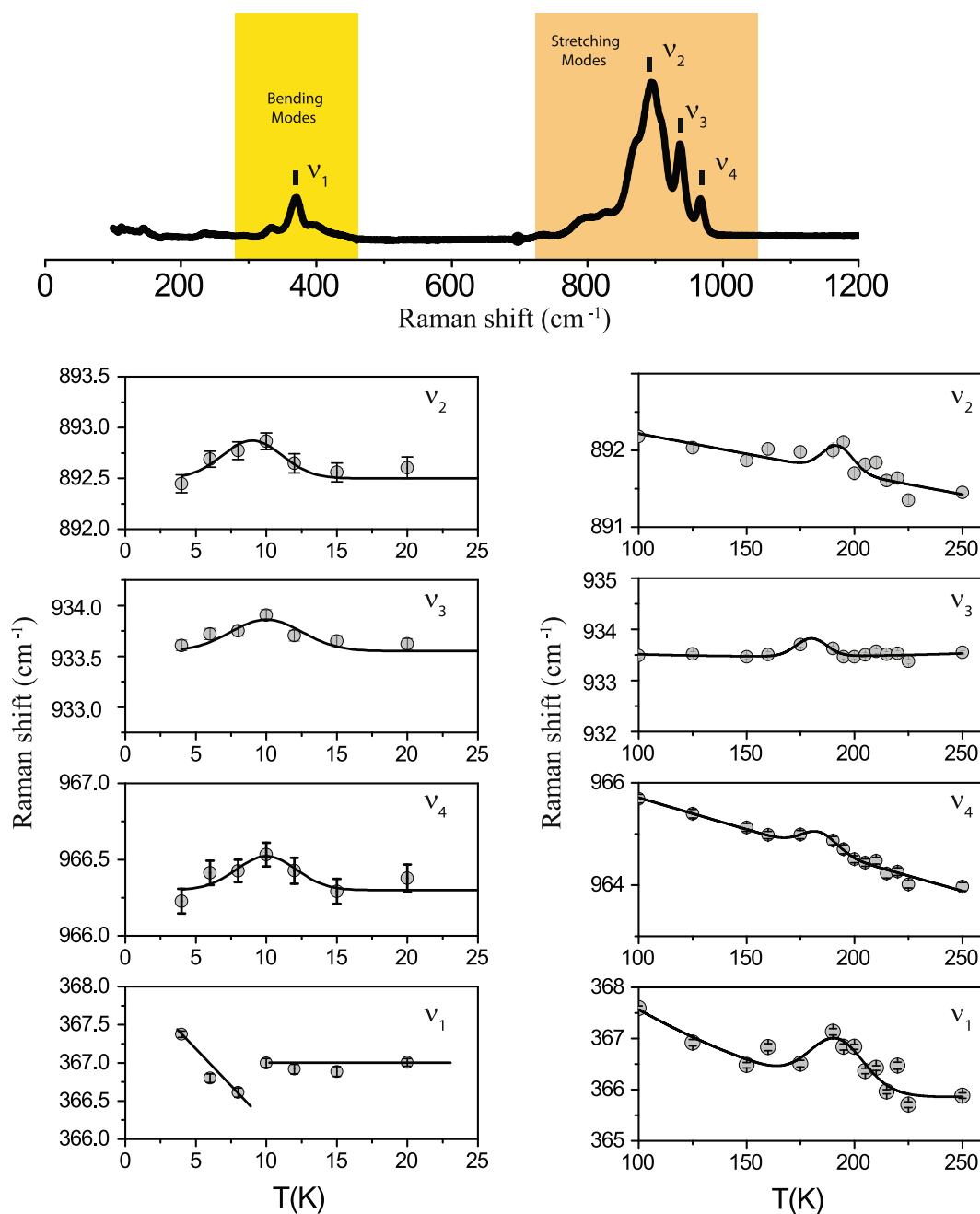
**Discussions.** The magnetization results establish that the title compound undergoes a complex long-range magnetic order below 8 K with a magnetic compensation point at a lower temperature. Heat-capacity data reveals another feature at a higher temperature, that is, around 225 K. The dielectric constant exhibits considerable frequency dispersion in this high temperature range. This is attributable to relaxor ferroelectric behaviour. This is a new finding as far as the property of this compound is concerned. A fitting of the peak in  $\tan\delta$  to Vogel-Fulcher relation yields parameters which are consistent with those for relaxor ferroelectrics.

The source of this phenomenon can be attributed to nanopolar regions created by crystallographic disorder intrinsic to this structure. It is known<sup>7</sup> that about 25% of Li goes to Ni1 site (honeycomb site), while 50% of Li goes to Ni2 site (chain site). The present  $^7\text{Li}$  NMR results provide a microscopic evidence for such a disorder, as evidenced by asymmetric spectra at low temperatures with different Li ions experiencing different internal hyperfine fields (with the onset of magnetic order of Ni) as well as by spin-lattice relaxation data. It is also important to note that the charge states of Li, Ni and Mo are different (+1, +2 and +6 respectively). Therefore, the presence of antisite Li results in random local electric field in the compound. Needless to stress that relaxor ferroelectricity due the charge imbalance at a particular site in the material due to disorder has been known to be one of the main ingredients of the relaxor ferroelectric<sup>21,22</sup>. Therefore we conclude that the generation of the polar nano region due to charge imbalance at various sites is responsible for the observed relaxor ferroelectricity.

The spectral features in the Raman spectroscopy vary non-monotonically in the neighborhood of 180 K and 10 K and it is highly likely that these variations in the peak positions near 180 K and 10 K are triggered by electrostriction due to relaxor ferroelectric behaviour and magnetostriction due to onset of magnetic order respectively.

We have made an intriguing finding in the magnetocapacitance data. The values of  $\tan\delta$  in the entire temperature range of investigation is extremely small, even in the presence of external magnetic fields. This finding emphasizes that the observed magnetodielectric effect at room temperature is an intrinsic effect and not due to any extrinsic factors.

It may be remarked that the sensitivity of the peak temperature in the pyrocurrent plot to heating rate is usually attributed to 'thermally stimulated depolarization current'<sup>27–31</sup> due to random trapping of mobile carriers. Since there appears to be the formation of nanopolar regions, the observed rate dependence may originate from



**Figure 7.** The room temperature Raman spectrum of  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  obtained by exciting the sample with an incident laser beam of wavelength 488 nm is shown in the top panel. The bottom left and right panels show the variation of the spectral features corresponding to  $\nu_1, \nu_2, \nu_3, \nu_4$  in the temperature range 4.2 to 25 K and 100 to 250 K. The solid lines are guides to the eyes.

‘dynamic disorder’, possibly due to the diffusion of small Li ion at various sites, which can be controlled by thermal effects. Such a dynamic disorder may contribute features mimicking ‘thermally stimulated depolarization current’.

In conclusion, the compound,  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , characterized by unique crystallographic features (namely, coexisting honeycomb and linear chains of magnetic ions) and ordering magnetically below 7.2 K, is investigated by magnetization, heat-capacity,  $^7\text{Li}$  NMR, Raman spectroscopy, dielectric, magnetodielectric and pyroelectric current measurements. The results reveal a new feature in the range 160–240 K - that is, relaxor ferroelectric behaviour attributable to crystallographic disorder intrinsic to this compound), with the pyrocurrent peak temperature tracking the poling temperature. In addition, intrinsic magnetodielectric effect is seen over a broad temperature range including room temperature.



## Methods

The polycrystalline  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$  was prepared by a conventional solid state route. The stoichiometric amounts of high purity (>99.9%)  $\text{Li}_2\text{CO}_3$ ,  $\text{NiO}$  and  $\text{MoO}_3$  were mixed together in an agate mortar, following by calcination at  $650^\circ\text{C}$  for 12 hours. The calcinated powder was then sintered at  $700^\circ\text{C}$  for 144 hours with intermediate grindings. XRD pattern was obtained with  $\text{Cu K}_\alpha$  ( $\lambda = 1.54 \text{ \AA}$ ).

Temperature dependent dc magnetization measurements were carried out with the help of a commercial SQUID magnetometer (Quantum Design, USA). Heat-capacity studies were carried out with a commercial Physical Properties Measurements System (PPMS) (Quantum Design, USA). Unlike otherwise stated, all these measurements were performed for the zero-field-cooled state (from 300 K) of the specimen. The same PPMS system was used to measure complex dielectric permittivity using an Agilent E4980A LCR meter with a home-made sample holder with several frequencies ( $\nu = 1$  to 100 kHz) and with a bias voltage of 1 V; this sample holder was also used for pyroelectric studies with Keithley 6517B electrometer by poling at 300 K with electric fields of 4 kV/cm.

$^7\text{Li}$  NMR experiments were carried out at 27 MHz by sweeping the magnetic field.  $T_1$  was measured by saturation recovery method using  $\pi/2 - \tau - \pi/2$  pulse sequence. Raman spectra at different intervals of temperature below 300 K were recorded using an optical cryostat with continuous helium flow (MicrostatHe: Oxford instruments) and single stage adaptation 1800grooves/mm spectrometer (T64000: Horiba Jobin Yvon). The sample was excited with 488 nm excitation wavelength from a mixed gas laser (Stabilite 2018: Spectra Physics).

## References

1. Tang, R., Zhou, H., You, W. & Yang, Y. Room-temperature multiferroic and magnetocapacitance effects in M-type hexaferrite  $\text{BaFe}_{0.2}\text{Sc}_{1.8}\text{O}_{19}$ . *App. Phys. Lett.* **109**, 082903 (2016).
2. Balédent, V. *et al.* Evidence for oom temperature electric polarization in  $\text{RMn}_2\text{O}_5$ . *Multiferroics. Phys. Rev. Lett.* **114**, 117601 (2015).
3. Ikeda, N., Mori, S. & Kohn, K. Charge ordering and dielectric dispersion in mixed valence oxides  $\text{RFe}_2\text{O}_4$ . *Ferroelectrics* **314**, 41–56 (2005).
4. Ikeda, N. *et al.* Ferroelectricity from iron valence ordering in the charge-frustrated system  $\text{LuFe}_2\text{O}_4$ . *Nature* **436**, 1136–1138 (2005).
5. Subramanian, M. A. *et al.* Giant room-temperature magnetodielectric response in the electronic ferroelectric  $\text{LuFe}_2\text{O}_4$ . *Adv. Mater.* **18**, 1737–1739 (2006).
6. Catalan, G. Magnetocapacitance without magnetoelectric coupling. *Appl. Phys. Lett.* **88**, 102902 (2006).
7. Hase, M. *et al.* Negative magnetization of  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ : A spin system composed of distorted honeycomb lattices and linear chains. *Phys. Rev. B* **84**, 104402 (2011).
8. Ozima, M., Sato, S. & Zoltai, T. The crystal structure of a lithium-nickel molybdate,  $\text{Li}_2\text{Ni}_2\text{Mo}_3\text{O}_{12}$ , and the systematics of the structure type. *Acta Cryst. B* **33**, 2175–2181 (1977).
9. Ibers, J. A. & Smith, G. W. Crystal structure of sodium cobalt molybdate. *Acta Cryst.* **17**, 190–197 (1964).
10. Pujari, S., Damle, K. & Alet, F. Neel-state to valence-bond-solid transition on the honeycomb lattice: Evidence for deconfined criticality. *Phys. Rev. Lett.* **111**, 087203 (2013).
11. Ye, F. *et al.* Direct evidence of a zigzag spin-chain structure in the honeycomb lattice: A neutron and x-ray diffraction investigation of single-crystal  $\text{Na}_2\text{IrO}_3$ . *Phys. Rev. B* **85**, 180403(R) (2012).
12. Mallik, R. *et al.* Complex magnetism in a new alloy,  $\text{Eu}_2\text{PdSi}_3$ , with two crystallographically inequivalent sites. *J. Magn. Magn. Mater.* **185**, L135–L143 (1998).
13. Patil, S., Iyer, K. K., Maiti, K. & Sampathkumaran, E. V. Behavior of magnetic ordering and the Kondo effect in the alloys  $\text{Ce}_2\text{Rh}_{1-x}\text{Co}_x\text{Si}_3$ . *Phys. Rev. B* **77**, 094443 (2008).
14. Mallik, R., Sampathkumaran, E. V., Strecker, M. & Wortmann, G. Observation of a minimum in the temperature-dependent electrical resistance above the magnetic-ordering temperature in  $\text{Gd}_2\text{PdSi}_3$ . *Europhys. Lett.* **41**, 315 (1998).
15. Ranjith, K. M. *et al.* Commensurate and incommensurate magnetic order in spin-1 chains stacked on the triangular lattice in  $\text{Li}_2\text{NiW}_2\text{O}_8$ . *Phys. Rev. B* **94**, 014415 (2016).
16. Paulose, P. L., Mohapatra, N. & Sampathkumaran, E. V. Spin-chain magnetism in Eu-doped  $\text{Ca}_3\text{Co}_2\text{O}_6$  and  $\text{Ca}_3\text{CoRhO}_6$  investigated by Mössbauer spectroscopy. *Phys. Rev. B* **77**, 172403 (2008).
17. Basu, T. *et al.* Displacive-type ferroelectricity from magnetic correlations within spin-chain. *Sci. Rep.* **4**, 5636 (2014).
18. Basu, T. *et al.* A reentrant phenomenon in magnetic and dielectric properties of  $\text{Dy}_2\text{BaNiO}_5$  and an intriguing influence of external magnetic field. *J. Phys.: Condens. Matter* **26**, 172202 (2014).
19. Basu, T., Iyer, K. K., Singh, K. & Sampathkumaran, E. V. Novel dielectric anomalies due to spin-chains above and below Néel temperature in  $\text{Ca}_3\text{Co}_2\text{O}_6$ . *Sci. Rep.* **3**, 3104 (2013).
20. Carvajal-Rodríguez, J. Recent advances in magnetic structure determination by neutron powder diffraction. *Physica B* **192**, 55–69 (1993).
21. Bokov, A. A. & Ye, Z.-G. Recent progress in relaxor ferroelectrics with perovskite structure. *Journal of Materials Science* **41**, 31–52 (2006).
22. Cross, L. E. Relaxor ferroelectrics. *Ferroelectrics* **76**, 241–267 (1987).
23. Kumar, S. *et al.* From spin induced ferroelectricity to spin and dipolar glass in a triangular lattice: The  $\text{CuCr}_{1-x}\text{V}_x\text{O}_2$  ( $0 \leq x \leq 0.5$ ) delafossite. *J. Sol. Stat. Chem.* **203**, 37–43 (2013).
24. Guo, J. *et al.* Normal-relaxor ferroelectric modulation of a-site complex perovskite ferroelectric  $(\text{K}_{1/2}\text{Bi}_{1/2})\text{TiO}_3$  by post-annealing. *J. App. Phys.* **121**, 014101 (2017).
25. Hemberger, J. *et al.* Relaxor ferroelectricity and colossal magnetocapacitance coupling in ferromagnetic  $\text{CdCr}_2\text{S}_4$ . *Nature* **434**, 364–367 (2005).
26. Smit, J. P., Kim, H.-S., Pless, J. D., Stair, P. C. & Poeppelmeier, K. R. Probing the vanadyl and molybdenyl bonds in complex vanadomolybdate structures. *Inorg. Chem.* **45**, 521–528 (2006).
27. De, C., Ghara, S. & Sundaresan, A. Effect of internal electric field on ferroelectric polarization in multiferroic  $\text{TbMnO}_3$ . *Solid State Commun.* **205**, 61–65 (2015).
28. Zou, T. *et al.* Excess-hole induced high temperature polarized state and its correlation with the multiferroicity in single crystalline  $\text{DyMnO}_3$ . *App. Phys. Lett.* **105**, 052906 (2014).
29. Kohara, Y., Yamasaki, Y., Onose, Y. & Tokura, Y. Excess-electron induced polarization and magnetoelectric effect in yttrium iron garnet. *Phys. Rev. B* **82**, 104419 (2010).
30. Upadhyay, S. K., Paulose, P. L., Iyer, K. K. & Sampathkumaran, E. V. Spin-glass behavior and pyroelectric anomalies in a new lithium-based oxide,  $\text{Li}_3\text{FeRuO}_5$ . *Phys. Chem. Chem. Phys.* **18**, 23348–23353 (2016).
31. Upadhyay, S. K., Iyer, K. K., Rayaprol, S., Paulose, P. L. & Sampathkumaran, E. V. A rock-salt-type Li-based oxide,  $\text{Li}_3\text{Ni}_2\text{RuO}_6$ , exhibiting a chaotic ferrimagnetism with cluster spin-glass dynamics and thermally frozen charge carriers. *Sci. Rep.* **6**, 31883 (2016).



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## Author Contributions

S.K.U. prepared the sample and characterized the same. He carried out magnetization, heat-capacity dielectric and pyrocurrent studies jointly with K.K.I. and analysed the results. P.L.P. performed NMR studies and analysed the results. Smita Gohil and Shankar Ghosh performed Raman studies and analysed the results. E.V.S. formulated the paper and finalized the manuscript in consultation with others for respective studies.

## Additional Information

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**Competing Interests:** The authors declare that they have no competing interests.

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