



# OPEN Exploring optoelectronic and photovoltaic properties of $\text{Be}_3\text{MF}_3$ ( $\text{M} = \text{P}, \text{As}, \text{Sb}$ ) perovskites via machine learning and numerical simulation

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This study presents a comprehensive Density Functional Theory (DFT) analysis of the structural, mechanical, electrical, and optical characteristics of cubic  $\text{Be}_3\text{MF}_3$  ( $\text{M} = \text{P}, \text{As}, \text{Sb}$ ) perovskites, highlighting its potential for optoelectronic applications. The computed elastic constants adhere to the Born stability criterion, affirming mechanical resilience and ductility. Analyses of band structure and density of states reveal that all  $\text{Be}_3\text{MF}_3$  compounds exhibit indirect band gaps between 2.0 and 3.3 eV, categorizing them as wide-bandgap semiconductors suitable for high-power and UV optoelectronic applications. Optical investigations demonstrate significant absorption, especially in  $\text{Be}_3\text{SbF}_3$ , highlighting its promise as a lead-free absorbing material for solar cell applications. Device-level performance was modeled utilizing SCAPS-1D, incorporating  $\text{Be}_3\text{SbF}_3$  as the absorber across diverse configurations with four electron transport layers ( $\text{In}_2\text{S}_3$ ,  $\text{TiO}_2$ ,  $\text{Mg:ZnO}$ ,  $\text{IGZO}$ ) and six-hole transport layers ( $\text{Cu}_2\text{O}$ ,  $\text{MoO}_3$ ,  $\text{P3HT}$ ,  $\text{NiO}$ ,  $\text{PTAA}$ ,  $\text{V}_2\text{O}_5$ ). The best configuration,  $\text{Al/FTO/In}_2\text{S}_3/\text{Be}_3\text{SbF}_3/\text{Cu}_2\text{O/Ni}$ , attained a maximum power conversion efficiency (PCE) of 18.28%, with open-circuit voltage ( $V_{\text{OC}}$ ) of 1.79 V, short-circuit current ( $J_{\text{SC}}$ ) of 11.04  $\text{mA}/\text{cm}^2$ , and a fill factor (FF) of 89.6%. Subsequent examinations of defect density, interface states, thickness variation, and temperature impacts clarify charge transport and recombination kinetics. A Random Forest machine learning (ML) model was created to forecast device performance, with high accuracy ( $R^2 = 0.987$ ,  $\text{MSE} = 0.00305$ ,  $\text{MAPE} = 0.075$ ) and pinpointing bandgap energy and interface defect density as critical limiting factors. The integrated DFT-SCAPS 1D-ML framework identifies  $\text{Be}_3\text{SbF}_3$  perovskites as viable, lead-free materials for advanced photovoltaic and optoelectronic technologies.

**Keywords**  $\text{Be}_3\text{MF}_3$  ( $\text{M} = \text{P}, \text{As}, \text{Sb}$ ),  $\text{Be}_3\text{SbF}_3$ , Optoelectronic and mechanical properties, SCAPS-1D solar cell simulation, Machine learning optimization

Addressing pollution, climate change, and energy insecurity requires a thorough shift to renewable energy systems, especially solar, wind, and hydropower resources. Developing countries receive sun energy much exceeding the world's total coal reserves and current global consumption, highlighting solar power's unparalleled potential<sup>1,2</sup>. As a result, photovoltaic (PV) technologies have emerged as a central focus of research and industry advancement. Recent advancements in perovskite-based solar cells have intensified this trend owing to their exceptional efficiency, cost-effectiveness, and adjustable material characteristics<sup>3,4</sup>. Perovskite solar cells (PSCs) are chiefly classified into  $\text{A}_3\text{BX}_3$  and  $\text{ABX}_3$  structural kinds, each providing distinct advantages

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regarding efficiency, stability, and scalability<sup>5</sup>. The  $ABX_3$  structure, consisting of metal cations (A, B) and halogen anions (X), is the most extensively researched and forms the basis of most metal halide perovskites<sup>6</sup>. Despite significant advancements improving PSC performance, cubic inorganic perovskites continue to encounter issues like moisture sensitivity, structural instability, and scalability in manufacturing. Organic perovskites have excellent efficiency but suffer from inadequate stability, whereas lead-based compounds encounter toxicity and deterioration when exposed to humidity, heat, and ultraviolet radiation<sup>7,8</sup>. Conversely, materials such as carbon nanotubes (CNTs), graphene, and transition metal dichalcogenides (TMDs) have enhanced stability, flexibility, and conductivity, rendering them optimal for hybrid solar designs<sup>9,10</sup>. Carbon nanotubes improve carrier mobility, moisture resistance, and interface quality when combined with perovskites, hence enhancing long-term stability. However, CNT-based devices frequently necessitate intricate surface modifications to attain good band alignment with the absorber layer. Furthermore, bilayer or multi-junction solar cells that incorporate materials with complementary band gaps enhance light absorption and improve power conversion efficiencies; yet, their production necessitates precise lattice and energy-level alignment, frequently resulting in interfacial instability<sup>11</sup>.

Researchers are investigating lead-free alternatives that utilize non-toxic cations, demonstrating enhanced chemical stability, tunable bandgaps, and extensive optoelectronic applications to surmount these challenges. Parmar et al.<sup>12</sup> investigated  $RbGeX_3$  ( $X = I, Br, Cl$ ) perovskites by DFT, demonstrating significant structural stability, direct bandgaps, and remarkable optical and thermoelectric properties, with anticipated efficiencies of 15.67%, 22.08%, and 28.93%, respectively. Rehman et al.<sup>13</sup> examined  $CsSnX_3$  ( $X = I, Br, Cl$ ) perovskites using DFT-CASTEP (GGA-PBE), validating their narrow bandgaps, elevated absorption coefficients, and malleable characteristics, hence identifying  $CsSnI_3$  as a notably effective contender for solar and LED applications. Zhang et al.<sup>14</sup> examined  $InXY_3$  ( $X = Ca, Sr, Ba$ ;  $Y = Cl, Br$ ) perovskites, revealing their mechanical strength, extensive bandgaps, and significant UV absorption coupled with visible-infrared transparency, making them suitable for transparent and UV-protective optoelectronic devices.

Inorganic compounds are pivotal to advancements in solar technology because of their distinctive physical and electrical characteristics. Recent research emphasizes materials characterized by the generic formula  $A_3BX_3$ , when A represents a large inorganic cation, B denotes a smaller metal cation, and X signifies a halogen anion<sup>15</sup>. Compounds like  $Mg_3SbF_3$ ,  $Ba_3PF_3$ , and  $Sr_3AsBr_3$  exhibit significant thermodynamic stability, ideal bandgaps, elevated absorption coefficients, and anticipated power conversion efficiencies (PCEs) over 25%, making them very interesting candidates for advanced solar cells<sup>16,17,18</sup>. Harun-Or-Rashid et al.<sup>19</sup> investigated Pb-free cubic  $A_3SbI_3$  ( $A = Ba, Sr, Ca$ ) perovskites by first-principles DFT and SCAPS-1D simulations, validating their mechanical and optical stability. Through the optimization of layer thickness, doping, and flaws,  $Ba_3SbI_3$  attained an exceptional power conversion efficiency (PCE) of 30.26%, surpassing previous versions. Hossain et al.<sup>20</sup> investigated  $A_3BF_3$  ( $A = Mg, Ca, Sr, Ba$ ;  $B = N, Bi$ ) perovskites, identifying cubic symmetry, mechanical stability, and significant visible-light absorption. Mg-based compounds functioned as insulators, but Ca-, Sr-, and Ba-based complexes exhibited semiconducting properties appropriate for solar, thermal, and optoelectronic applications. Ghosh et al.<sup>21</sup> conducted an in-depth analysis of  $A_3BX_3$  ( $A = Ca, Sr$ ;  $B = P, As$ ;  $X = I, Br$ ) halide perovskites, documenting direct bandgaps ranging from 1.26 to 2.63 eV, substantial absorption, and elevated efficiency, with  $Sr_3AsI_3$  attaining a PCE of 28.16%. Phosphorus (P), Arsenic (As), and Antimony (Sb) are crucial to improving perovskite efficacy. P serves as a bandgap-tuning dopant, enhancing light absorption and charge mobility; As alters the band structure and prolongs carrier lifetimes; and Sb substitutes hazardous lead, thereby enhancing stability and environmental compatibility<sup>22,23,24</sup>. Collectively, these components provide efficient, lead-free, and resilient perovskite configurations. Nonetheless, synthesis difficulties remain due to the volatility of halogens particularly fluorine (F) and the necessity for meticulous control of stoichiometry, phase purity, and oxidation states, especially in antimony-based systems<sup>25</sup>. Notwithstanding these challenges, ongoing enhancement of synthesis procedures and sophisticated modeling approaches, including GW-BSE simulations to evaluate exciton binding energy<sup>26</sup>, are anticipated to expedite the advancement of these materials for next-generation optoelectronic and solar applications.

This study extensively investigated the structural, mechanical, optical, and electrical properties of  $Be_3MF_3$  ( $M = P, As, Sb$ ) halide perovskites by first-principles DFT calculations to evaluate their suitability for PV applications. Among the candidates,  $Be_3SbF_3$  demonstrated the most advantageous bandgap, stability, and light absorption properties, establishing it as a formidable prospect for solar energy conversion. The SCAPS-1D simulation tool was utilized to assess the PV performance of  $Be_3SbF_3$ -based PSCs. An analysis was conducted on various hole transport layers (HTLs) such as  $Cu_2O$ ,  $MoO_3$ ,  $P3HT$ ,  $NiO$ ,  $PTAA$ , and  $V_2O_5$  and electron transport layers (ETLs) including  $In_2S_3$ ,  $TiO_2$ ,  $Mg:ZnO$ , and  $IGZO$  to determine the best device configuration. The impact of absorber thickness, doping concentration, and defect density on the principal PV parameters was methodically investigated. Furthermore, Nyquist plots, J-V characteristics, and quantum efficiency (QE) investigations were performed to investigate charge transport, recombination, and overall device efficiency. Temperature-dependent performance evaluations validated the material's thermal stability. A Random Forest (RF) machine learning model was created to forecast the photovoltaic performance of  $Be_3SbF_3$ -based solar cells utilizing an extensive dataset of 13,113 data points obtained from device simulations. The model underwent training and evaluation utilizing several interpretative techniques, such as correlation heatmaps, SHAP (Shapley Additive Explanations) importance analysis, parity plots, and confusion matrix assessments, guaranteeing thorough feature evaluation and predictive reliability. The RF model exhibited significant predictive accuracy, successfully identifying bandgap energy and interface defect density as the primary factors influencing device efficiency. The combined use of DFT-derived material descriptors, SCAPS-1D device simulations, and ML-based predictions creates a cohesive framework for elucidating performance-limiting mechanisms and progressing  $Be_3SbF_3$  perovskite solar cells towards scalable, lead-free, and sustainable photovoltaic applications.

## Computational methodology

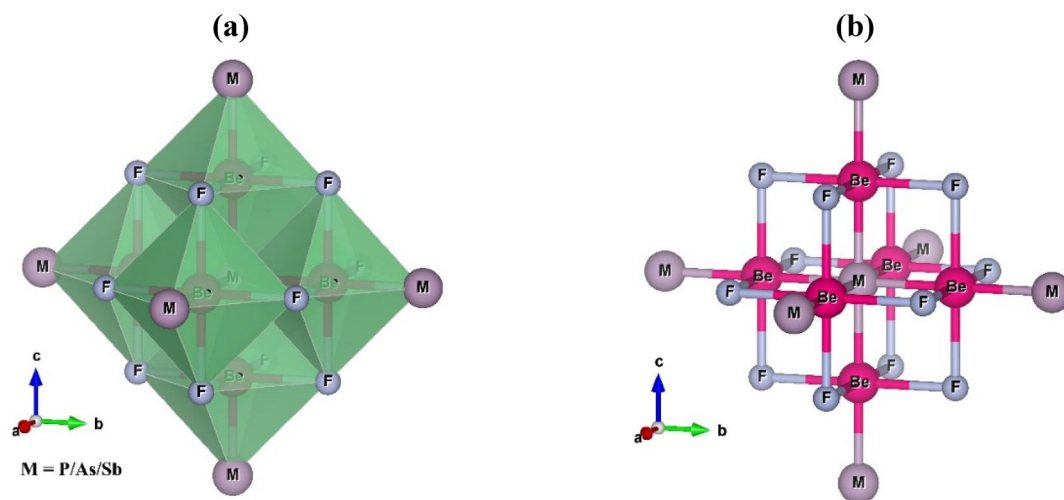
A Density Functional Theory (DFT)-based technique was employed inside the Cambridge Serial Total Energy Package (CASTEP) framework to examine the structural, electrical, and optical properties of  $\text{Be}_3\text{MF}_3$  ( $\text{M}=\text{P}, \text{As}, \text{Sb}$ ) perovskite material. The Generalized Gradient Approximation (GGA) utilizing the Perdew–Burke–Ernzerhof (PBE) functional was utilized to address the exchange–correlation potential and assess the electronic band structure<sup>27</sup>. The Vanderbilt-type ultrasoft pseudopotential On–The–Fly–Generated (OTFG) was employed to precisely simulate the interaction between valence electrons and ionic cores, ensuring computational efficiency and accuracy in total energy calculations<sup>28</sup>. The Broyden–Fletcher–Goldfarb–Shanno (BFGS) minimization technique was employed to optimize the crystal shape, allowing the system to achieve its minimum energy configuration<sup>29</sup>. The Brillouin zone was sampled with a Monkhorst–Pack k-point grid of  $10 \times 10 \times 10$ , and a plane-wave cutoff energy of 450 eV was utilized. This selection was established by rigorous convergence testing, which verified that these parameters provide very accurate assessments of the structural, electrical, and optical properties without incurring excessive computational costs. The selected k-point density effectively encapsulates the fundamental electronic interactions stemming from the inherent symmetry of the crystal lattice, guaranteeing numerical accuracy in the calculation of mechanical properties, including the elastic constants. The alignment of these computational settings with previously reported<sup>21</sup> further validates their reliability and appropriateness for attaining a balanced trade-off between accuracy and computational efficiency. The elastic properties were ascertained utilizing the finite strain–stress methodology, as executed in CASTEP, to compute the elastic stiffness constants ( $C_{ij}$ ) under equilibrium conditions<sup>30</sup>. Geometry optimization adhered to rigorous convergence criteria: a maximum displacement of  $5 \times 10^{-4}$  Å, a maximum force of 0.01 eV/Å, a maximum stress of 0.05 GPa, and a total energy convergence of  $1 \times 10^{-5}$  eV per atom. A strain magnitude of 0.003 was utilized in the elastic calculations, guaranteeing computational stability and alignment with previous theoretical studies. These criteria combined guaranteed the dependability and reproducibility of the results obtained. Additionally, to augment the first-principles findings, the SCAPS-1D software was utilized to model charge carrier transit, electrostatic potential distribution, and overall device performance under steady-state conditions<sup>31,32</sup>. The model quantitatively resolved Poisson's and continuity equations for electrons and holes, yielding a comprehensive insight into recombination kinetics, carrier diffusion, and drift phenomena<sup>33,34</sup>. This hybrid computational method integrating quantum–mechanical DFT insights with numerical device modeling facilitated a thorough evaluation of the  $\text{Be}_3\text{SbF}_3$ -based PSC, providing essential guidance for enhancing its optoelectronic characteristics and operational efficiency in advanced PV applications.

## Result and discussion

### DFT result

#### Structural properties

The beryllium-based single halide perovskite compounds, represented as  $\text{Be}_3\text{MF}_3$  (where  $\text{M}=\text{P}, \text{As}, \text{Sb}$ ), crystallize in a cubic  $\text{A}_3\text{BX}_3$ -type perovskite structure that corresponds to the Pm-3 m crystallographic space group (No. 221)<sup>35</sup>. Figure 1a, b illustrates the optimal crystal structure of the  $\text{Be}_3\text{MF}_3$  systems, showcasing a well-ordered cubic layout including a single formula unit per cell, consisting of seven atoms. The optimized lattice parameters and associated unit cell volumes, derived from GGA-PBE functional, are presented in Table 1. In these structures, the Be atoms predominantly occupy the 3a Wyckoff positions ( $1/2, 1/2, 0$ ), whereas the M-site elements (P, As, Sb) are located at the 1b Wyckoff locations ( $1/2, 1/2, 1/2$ ), and the F atoms are situated at the 3c Wyckoff sites ( $1/2, 0, 0$ ). A comprehensive geometry optimization was conducted, permitting the total relaxation of lattice parameters to achieve the equilibrium configuration. The derived lattice constants and total energy values exhibit remarkable concordance with existing theoretical and experimental findings, hence affirming the



**Fig. 1.** Graphic representation of  $\text{Be}_3\text{NX}_3$  compounds: (a) polyhedral view, (b) ball-and-stick model, for  $\text{Be}_3\text{MF}_3$ .

Material name	Band gap ( $E_g$ )	Lattice constant ( $\text{\AA}$ )	Unit cell volume ( $\text{\AA}^3$ )	$\tau_G$	$E_f$ (eV/atom)	Refs.
$\text{Be}_3\text{PF}_3$	3.327	4.187	73.409	0.94	−3.98	This work
$\text{Be}_3\text{AsF}_3$	3.201	4.281	78.500	0.90	−3.72	This work
$\text{Be}_3\text{SbF}_3$	2.158	4.565	95.166	0.85	−3.30	This work
$\text{Mg}_3\text{PBr}_3$	1.43	5.463	163.003	0.854	−3.35	38
$\text{Mg}_3\text{SbF}_3$	3.46	5.02	126.51	0.794	−3.18	16
$\text{Ba}_3\text{SbF}_3$	1.07	6.407	263.03	0.89	−4.59	39
$\text{Ca}_3\text{AsCl}_3$	1.742	5.756	190.7	0.849	−4.85	40

**Table 1.** Optimized structural models and associated thermodynamic stability parameters for  $\text{Be}_3\text{MF}_3$ .

computational precision and dependability of the current results. To further validate the geometric stability of the cubic structure, the Goldschmidt tolerance factor ( $\tau_G$ ) was computed using Eq. (1)<sup>36</sup>.

$$\tau_G = \frac{2(R_A + R_X)}{\sqrt{2}(R_B + 2R_X)} \quad (1)$$

where  $R_A$ ,  $R_B$ , and  $R_X$  specify the ionic radii of the A-site cation (Be), B-site atom (P, As, or Sb), and X-site anion (F), accordingly. The calculated tolerance factor of  $\tau_G$  between 0.804 to 1 is comfortably situated within the defined stability range for cubic perovskites ( $0.71 < \tau < 1.00$ ), affirming the maintenance of a stable cubic phase devoid of significant lattice aberrations. This geometric stability is a vital prerequisite for sustaining consistent optoelectronic properties, which are crucial for dependable device performance over prolonged operational durations. The thermodynamic stability of the  $\text{Be}_3\text{MF}_3$  compounds was assessed by computing the formation energy ( $E_f$ ) as per Eq. (2)<sup>37</sup>.

$$E_f = E_{\text{Be}_3\text{MF}_3} - (3E_{\text{Be}} + E_{\text{M}} + 3E_{\text{F}}) \quad (2)$$

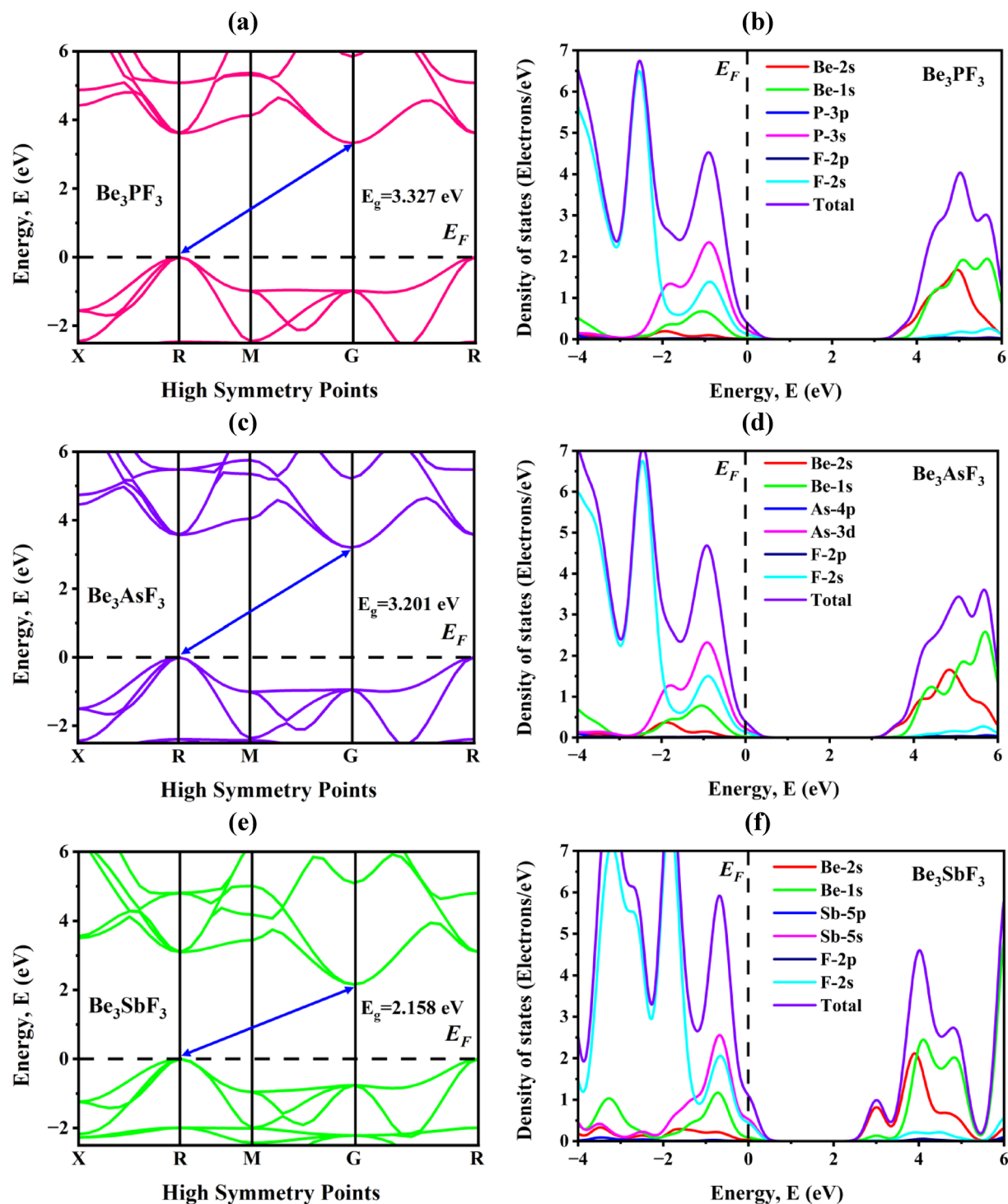
where  $E_{\text{Be}_3\text{MF}_3}$  denotes the cumulative energy of the optimum configuration, and  $E_{\text{Be}}$ ,  $E_{\text{M}}$  and  $E_{\text{F}}$  correspond to the aggregate energies of the constituent Be, M, and F atoms, respectively. The negative formation energy estimates (Table 1) indicate that all  $\text{Be}_3\text{MF}_3$  molecules are energetically favorable and thermodynamically stable, hence affirming their viability for experimental synthesis. The comprehensive evaluation of optimized lattice parameters, tolerance factor, and formation energy confirms the significant structural integrity of the  $\text{Be}_3\text{MF}_3$  perovskites. This intrinsic stability improves their prospective incorporation into solar, optoelectronic, and advanced energy-harvesting applications, where phase resilience and minimum structural variation are essential for attaining superior device performance and operational longevity.

#### Electronic band structure and density of states

The electronic band structure and density of states (DOS) are essential instruments for elucidating the electronic characteristics of crystalline solids, since they disclose the distribution of electrons across different energy levels  $E$  (in eV). Examining the band structure yields essential insights into a material's conductive, semiconductive, or insulating properties, which directly influence its electrical, optical, and thermal responses. This comprehension is essential for customizing materials for sophisticated electrical applications, such as semiconductors, field-effect transistors, photodetectors, and PV cells. Figure 2a, c, e illustrates the computed band structures of  $\text{Be}_3\text{MF}_3$  ( $M = \text{P, As, Sb}$ ) compounds derived from the GGA-PBE functional. The computations were executed along high-symmetry directions ( $X-R-M-\Gamma-R$ ) within the Brillouin zone, with the Fermi energy level ( $E_F$ ) indicated by a red dashed line set at 0 eV. The valence bands are situated beneath the  $E_F$  whilst the conduction bands are positioned above them. The lack of band overlap at  $E_F$  signifies that all  $\text{Be}_3\text{MF}_3$  compounds possess semiconducting properties. The PBE-GGA functional, despite inherent self-interaction inaccuracies, frequently underestimates the band gap owing to its semi-local exchange–correlation characteristics. This constraint is particularly apparent in halide-based perovskites, where robust electron–electron interactions govern the electrical and optical properties<sup>41</sup>. Nonetheless, PBE-GGA continues to be highly proficient in capturing qualitative patterns in electrical structure and band dispersion. Advanced methodologies, including hybrid functionals (HSE06) and techniques such as GW or mBJ, enhance bandgap precision by integrating accurate exchange corrections, however they entail significantly more computational expense<sup>42</sup>. PBE-GGA, while producing lower absolute gap values, provides an exceptional balance between computational efficiency and physical reliability, rendering it ideal for extensive and defect-focused simulations. Its strength in consistently recreating electrical and optical trends guarantees dependable insights into structural and compositional changes in intricate perovskite systems.

The investigation of the electronic band structure indicates that the valence band maximum (VBM) and conduction band minimum (CBM) of  $\text{Be}_3\text{MF}_3$  ( $M = \text{P, As, Sb}$ ) are located at different high-symmetry points (R-G) in the Brillouin zone, hence validating their indirect band gap characteristic. The determined band gaps 3.327 eV for  $\text{Be}_3\text{PF}_3$ , 3.201 eV for  $\text{Be}_3\text{AsF}_3$ , and 2.158 eV for  $\text{Be}_3\text{SbF}_3$ , as shown in Table 1, categorize these compounds as wide-bandgap semiconductors. Indirect band gap materials demonstrate phonon-assisted electronic transitions, leading to less non-radiative recombination and extended carrier lifetimes. These attributes offer improved thermal stability, radiation resistance, and mechanical durability, which are advantageous for high-power, high-temperature, and long-lifespan electronics. Despite their reduced efficacy in light emission, these materials are exceptionally appropriate for nonlinear optical (NLO) and passive photonic applications, where accurate light





**Fig. 2.** The calculated electronic band structure and DOS plots for (a, b)  $\text{Be}_3\text{PF}_3$ , (c, d)  $\text{Be}_3\text{AsF}_3$ , and (e, f)  $\text{Be}_3\text{SbF}_3$  perovskite.

modulation and elevated damage thresholds are critical<sup>43,44</sup>. In optoelectronic device engineering, the indirect characteristics of  $\text{Be}_3\text{MF}_3$  compounds provide benefits for UV and deep-UV photodetectors, electro-optic modulators, waveguides, thermophotovoltaic converters, and optical limiters, where regulated light absorption and minimal recombination losses are essential. The extensive band gaps indicate robust interatomic bonding, elevated breakdown voltage, and substantial optical phonon energies, resulting in enhanced thermal conductivity and prolonged device lifespan. The intrinsic properties of  $\text{Be}_3\text{MF}_3$  compounds render them viable candidates for energy-harvesting systems, radiation-resistant photovoltaic cells, and high-performance optoelectronic devices in sophisticated solar and photonic technologies. Among them,  $\text{Be}_3\text{SbF}_3$  is notable for its optical transitions, substantial mechanical stability, and a moderate band gap, facilitating significant absorption throughout the

visible to near-UV spectrum. Despite its band gap being marginally greater than that of optimal single-junction absorbers,  $\text{Be}_3\text{SbF}_3$  is still highly efficient for tandem solar cell configurations, serving as the wide-bandgap top absorber in conjunction with a low-bandgap bottom cell (such as silicon or tin-based perovskites) to expand the spectral response and improve overall efficiency<sup>45</sup>.

The Total and Partial density of states (TDOS and PDOS) investigations offer significant insights into the electronic configuration and bonding properties of crystalline materials. TDOS provides a comprehensive overview of the total accessible electronic states across various energy levels, whereas PDOS disaggregates these contributions by specific atomic orbitals or elements, thereby clarifying the local electronic environment and the characteristics of orbital hybridization<sup>46</sup>. Collectively, these investigations are essential for elucidating the origins of band structure characteristics and comprehending the chemical bonding mechanisms that dictate material stability and electrical performance. Figure 2b, d, f displays the calculated TDOS and PDOS spectra for  $\text{Be}_3\text{MF}_3$  ( $M = \text{P, As, Sb}$ ) compounds, plotted against energy with  $E_F$  set at 0 eV (denoted by a black dashed line). The valence band, ranging from approximately  $-4$  to  $0$  eV, is predominantly influenced by the F-2p orbitals, signifying that fluorine atoms significantly contribute to the characteristics of the upper valence band and the nature of chemical bonding.

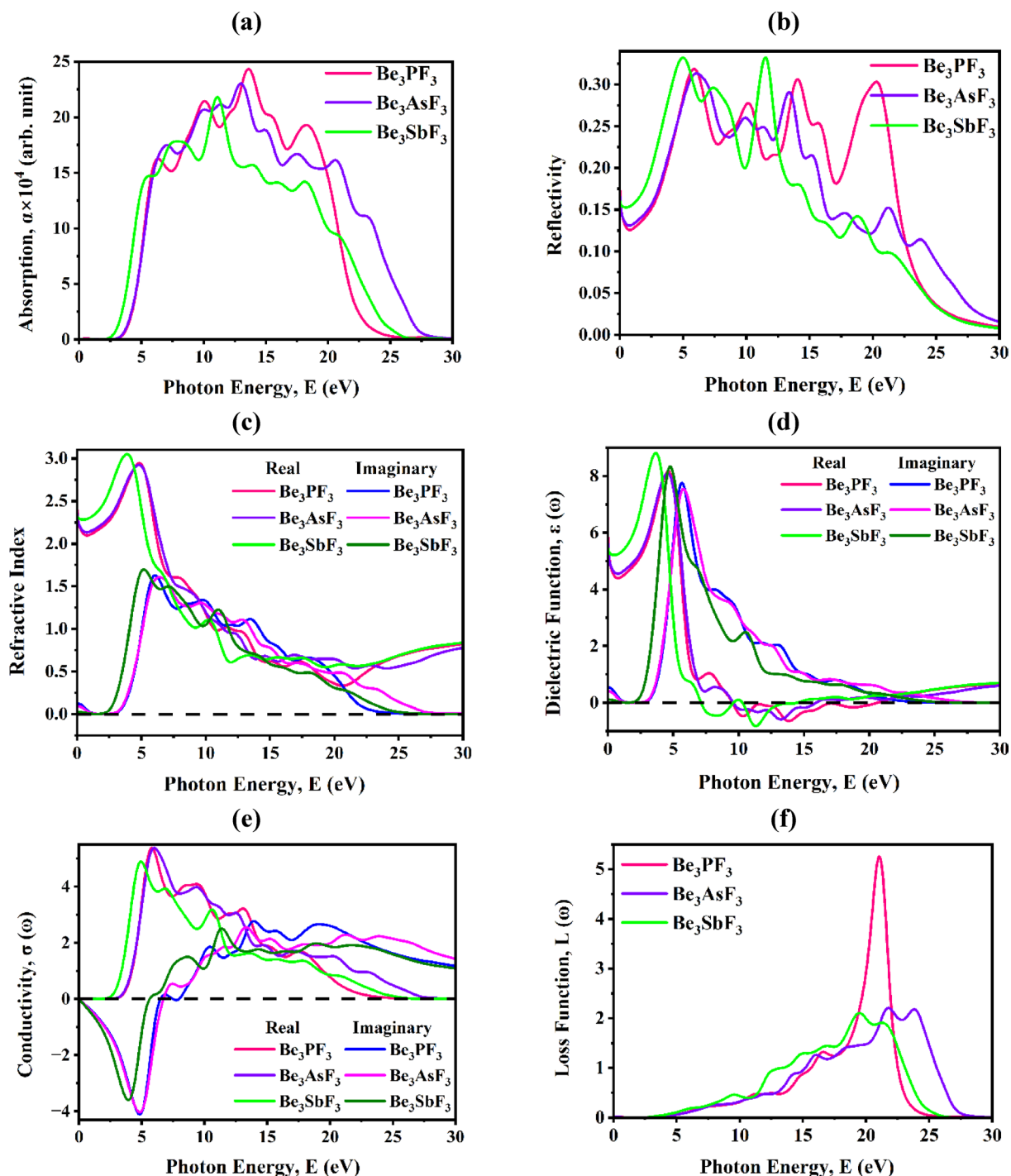
Minor hybridization contributions from the P-3 s, As-4p, and Sb-5 s states are found near the valence band edge, indicating slight differences in electronic interactions across the series as the atomic mass of M increases. Conversely, the conduction band mostly derives from the Be-2 s and Be-2p states, supplemented by contributions from the higher-energy orbitals of group-V elements (P, As, Sb). This indicates a robust covalent interaction between Be-F and Be-M, wherein beryllium orbitals substantially affect electronic transitions and play a role in the conduction process. The extent of orbital overlap between beryllium and fluorine states suggests increased charge delocalization, potentially enhancing electron mobility and overall charge transport characteristics. The prevalence of F-2p states at the valence band maximum and Be-2 s states at the CBM confirms the indirect band gap characteristic revealed in the band structure study. The uniform electronic alignment in all three compounds emphasizes their semiconducting properties and shows the influence of fluorine and beryllium orbitals on electronic transitions. Furthermore, the progressive decline of the conduction band edge from  $\text{Be}_3\text{PF}_3$  to  $\text{Be}_3\text{SbF}_3$  indicates a trend of band gap shrinking, along with the augmented atomic size and diminished electronegativity of the group-V element. This electrical development is essential for optimizing optical absorption and customizing these materials for PV, photodetector, and light-emitting device applications, where precise band alignment improves energy conversion efficiency.

#### Optical characteristics

The optical response of a material offers essential insights into its capacity to interact with, absorb, and transmit electromagnetic radiation, which are crucial for assessing its appropriateness for photonic, optoelectronic, and PV applications. Comprehension of optical behavior is essential for materials designed for solar energy harvesting, light-emitting devices, photodetectors, and nonlinear optical systems. These features are predominantly determined by electronic structure and chemical bonding factors, which affect electron transitions between valence and conduction bands when exposed to incident light. The investigation of halide-based perovskites has proliferated in recent years owing to their exceptional light absorption, adjustable band gaps, and elevated carrier mobility. The optical properties of  $\text{Be}_3\text{MF}_3$  ( $M = \text{P, As, Sb}$ ) were examined by first-principles DFT to evaluate their applicability in advanced energy conversion and optical communication technologies. The analyzed optical functions encompass the absorption coefficient ( $\alpha(\omega)$ ), reflectivity ( $R(\omega)$ ), refractive index ( $n(\omega)$ ), dielectric function ( $\epsilon(\omega)$ ), optical conductivity ( $\sigma(\omega)$ ), and energy loss function ( $L(\omega)$ ), all assessed up to 30 eV photon energy, as seen in Fig. 3a–f. These features jointly delineate the propagation of electromagnetic waves within the material and the efficiency of their interaction with photons across various frequency spectra.

The absorption coefficient ( $\alpha(\omega)$ ) indicates the capacity of  $\text{Be}_3\text{MF}_3$  compounds to absorb photons and generate charge carriers, as illustrated in Fig. 3a, and serves as an essential indicator of their light-harvesting effectiveness. The absorption spectra initiate at specific photon energies that align with their optical band gaps, in accordance with the semiconducting characteristics derived from band structure research. The tight correlation between electronic and optical data validates the trustworthiness and internal consistency of band structure computations<sup>47</sup>. Numerous pronounced absorption peaks are observable over the UV spectrum, indicating effective photon electron interactions at elevated energies. Among the compounds,  $\text{Be}_3\text{PF}_3$  has the most pronounced absorption peak at approximately 13.62 eV with an intensity of 24.27, while  $\text{Be}_3\text{AsF}_3$  and  $\text{Be}_3\text{SbF}_3$  provide peaks around 23.02 and 21.7, respectively. The existence of several maxima and minima throughout the spectrum signifies intricate interband electronic transitions involving Be-F and Be-M orbital hybridizations. Despite the absorption intensity being comparatively low below 2.5 eV because to the indirect bandgap characteristics and fluorine-induced electronic localization, a pronounced rise is noted at elevated photon energies, signifying effective photon–electron coupling in the UV–visible spectrum. The notable optical activity indicates that  $\text{Be}_3\text{MF}_3$  compounds, especially  $\text{Be}_3\text{SbF}_3$ , hold considerable promise for incorporation into UV photodetectors, radiation-resistant coatings, and tandem solar cell designs<sup>48</sup>. Their elevated absorption coefficients in the ultraviolet spectrum render them advantageous for high-power optoelectronic devices, UV filters, and laser protection applications, where wide bandgap materials with resilient optical responses are essential for performance and longevity.

The reflectivity,  $R(\omega)$ , quantifies the efficiency with which a substance reflects incoming electromagnetic radiation. This characteristic directly affects the optical contrast, thermal regulation, and energy retention of materials employed in photonic systems and solar cells. Figure 3b illustrates the reflectivity spectra of  $\text{Be}_3\text{MF}_3$  compounds across an extensive photon energy range. At 0 eV, the reflectivity initiates at moderate levels ranging from 0.16 to 0.175, indicating the intrinsic optical density of the compounds. With increasing photon energy, reflectivity progressively ascends in the infrared and visible spectra, reaching peak values of 0.318 for



**Fig. 3.** The optical characteristics of  $\text{Be}_3\text{MF}_3$  perovskite compounds encompass (a) absorption spectra, (b) reflectivity, (c) variations in refractive index, (d) the real and imaginary components of the dielectric function, (e) optical conductivity, and (f) electron energy loss function, underscoring their applicability in optoelectronic and PV domains.

$\text{Be}_3\text{PF}_3$ , 0.313 for  $\text{Be}_3\text{AsF}_3$ , and 0.332 for  $\text{Be}_3\text{SbF}_3$ . This little difference can be ascribed to the rising atomic mass and the associated alteration in electronic polarizability of the M-site elements (P, As, Sb). At elevated photon energy (exceeding 27 eV), the reflectance of all three substances declines precipitously, nearing zero. The decreasing reflectance signifies a transparent optical window in the UV–visible spectrum, facilitating effective photon transmission through the material<sup>49</sup>. This optical characteristic is essential for anti-reflective coatings, transparent electrodes, and optical encapsulants utilized in PV and photonic devices, as diminished surface reflection improves energy absorption and device performance. Moreover, the prevailing tendency indicates that

the  $\text{Be}_3\text{MF}_3$  perovskites exhibit an intrinsic equilibrium between reflection and absorption a desirable attribute for light-trapping structures in thin-film solar cells.

The refractive index,  $n(\omega)$ , measures the extent of light refraction as it traverses a material and offers insight into optical polarization and dielectric behavior. The real and imaginary components of  $n(\omega)$  were computed, as depicted in Fig. 3c. The static refractive index values,  $n(0)$ , vary from 2.3 to 2.5, indicating a modest optical density characteristic of wide band gap semiconductors. In the infrared and visible spectra, the real part of  $n(\omega)$  progressively ascends, attaining peak values of 2.92 for  $\text{Be}_3\text{PF}_3$ , 2.91 for  $\text{Be}_3\text{AsF}_3$ , and 3.99 for  $\text{Be}_3\text{SbF}_3$  within the energy range of 3.1–4.8 eV. This increase implies robust light–matter interaction and proposes possible applications in nonlinear optical modulation and light confinement systems. As photon energy increase, the refractive index progressively diminishes, signifying a reduction in transparency and a shift towards elevated optical absorption. The imaginary component of  $n(\omega)$ , denoting the extinction coefficient, stays minimal at elevated energies, indicating diminished photon loss due to scattering or absorption. The near resemblance in the refractive index values of  $\text{Be}_3\text{PF}_3$  and  $\text{Be}_3\text{AsF}_3$  suggests analogous bonding properties and electrical behaviors, whereas  $\text{Be}_3\text{SbF}_3$ , with a little elevated peak, may provide improved optical polarization due to its greater atomic size and diminished band gap. The refractive index analysis indicates that  $\text{Be}_3\text{MF}_3$  compounds exhibit excellent optical transparency and moderate refractive properties, making them appropriate for light-guiding components, photonic waveguides, and electro-optic modulators.

The dielectric function,  $\epsilon(\omega) = \epsilon_1(\omega) + i\epsilon_2(\omega)$ , characterizes a material's reaction to an external electromagnetic field and offers essential understanding of polarization effects and optical transition processes. The real component,  $\epsilon_1(\omega)$ , signifies the dispersive (polarizable) aspect, whereas the imaginary component,  $\epsilon_2(\omega)$ , pertains to absorption and energy dissipation resulting from interband transitions. The calculated dielectric spectra are displayed in Fig. 3d. At zero photon energy, the actual dielectric constant  $\epsilon_1(0)$  ranges from 5.4 to 5.9, signifying substantial electronic polarization<sup>50</sup>. The real component displays peak values between 8.1 and 8.9 at around 4.5 eV, indicating significant dipole excitation within this energy spectrum. The imaginary component,  $\epsilon_2(\omega)$ , has pronounced peaks in the visible spectrum and progressively diminishes to zero beyond 25 eV, indicating negligible energy loss at elevated photon energies. The characteristics indicate that the  $\text{Be}_3\text{MF}_3$  perovskites exhibit superior dielectric stability and frequency-dependent polarization regulation, essential for electro-optical modulators, optical sensors, and capacitive energy storage systems. The significant correlation between  $\epsilon_1(\omega)$  and  $n(\omega)$  profiles further validates the self-consistency of the optical analysis and emphasizes the resilience of the materials' optical polarization response.

The optical conductivity ( $\sigma(\omega)$ ), measures a material's capacity to transfer charge carriers when subjected to optical excitation and is essential for evaluating photoconductive efficiency. Figure 3e illustrates the real and imaginary components of  $\sigma(\omega)$  for the  $\text{Be}_3\text{MF}_3$  series. The real component of  $\sigma(\omega)$  displays distinct peaks between 4.7 and 5.3, indicating active interband transitions and effective electron excitation across the optical band gap. The analogous magnitudes among the three chemicals indicate similar electron mobility in optical fields. The imaginary component, linked to phase lag in carrier response, exhibits a gradual decrease over the infrared and visible spectra, signifying diminished reactive losses. The elevated optical conductivity in the visible spectrum indicates robust light-induced carrier excitation and effective current responsiveness when illuminated. Their attributes, along with wide band gap stability, render  $\text{Be}_3\text{MF}_3$  materials exceptional prospects for photoconductive sensors, transparent conductors, and energy-efficient optoelectronic systems.

The energy loss function ( $L(\omega)$ ) indicates the energy dissipated by a high-velocity electron as it moves through a solid and offers direct insights into plasmon resonances collective oscillations of free electrons. Figure 3f depicts the loss function spectra for the  $\text{Be}_3\text{MF}_3$  molecules. With reduced photon energies, the loss function diminishes owing to heightened photon absorption within that spectrum. Distinct plasmon resonance peaks arise at elevated energy, corresponding to plasma frequencies where collective electronic oscillations prevail<sup>51</sup>.

The peak  $L(\omega)$  for  $\text{Be}_3\text{PF}_3$  is around 21.06 eV ( $\approx 5.25$ ), whereas  $\text{Be}_3\text{AsF}_3$  and  $\text{Be}_3\text{SbF}_3$  display lower peak values of 2.2 and 2.1, respectively. The peaks signify robust electron–photon interaction and moderate energy dissipation, validating the promise of  $\text{Be}_3\text{MF}_3$  compounds for plasmonic and photothermal applications, where regulated energy loss and resonance characteristics are beneficial. Furthermore, the comparatively low loss function at visible photon energies indicates negligible damping and superior optical quality, rendering these materials suitable for transparent coatings, UV shielding layers, and light management systems in optoelectronic devices.

#### Mechanical characteristics

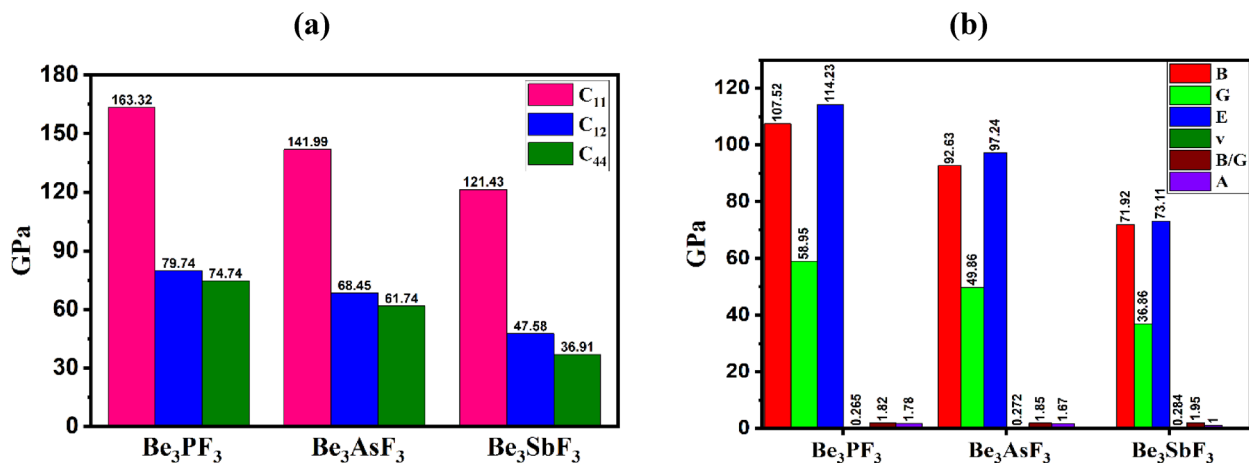
The mechanical properties of crystalline materials are essential for comprehending their stability, deformation behavior, and applicability in industrial and optoelectronic fields. Elastic constants ( $C_{ij}$ ) are essential markers of interatomic bonding strength, measuring a solid's response to external stress and strain. This study comprehensively examined the elastic characteristics of  $\text{Be}_3\text{MF}_3$  ( $M = \text{P, As, Sb}$ ) utilizing the linear finite strain–stress methodology in CASTEP. Figure 4a, b illustrates that the mechanical properties of these cubic systems are determined by three distinct elastic constants  $C_{11}$ ,  $C_{12}$ , and  $C_{44}$  which correspond to resistance against uniaxial, uniform, and shear deformations, respectively. To guarantee mechanical stability, the Born–Huang stability criteria must be fulfilled<sup>52</sup>, articulated as:

$$C_{11} + 2C_{12} > 0 \quad (\text{Spinodal criteria})$$

$$C_{11} - C_{12} > 0 \quad (\text{Born criteria})$$

$$C_{44} > 0 \quad (\text{Shear criteria})$$





**Fig. 4.** Values of calculated (a) elastic and (b) mechanical constants for Be<sub>3</sub>MF<sub>3</sub> (M = P, As, Sb).

Figure 4a demonstrates that all Be<sub>3</sub>MF<sub>3</sub> compounds meet the mechanical stability criterion, with C<sub>11</sub> significantly above C<sub>12</sub> and C<sub>44</sub>, signifying greater resistance to longitudinal compression compared to shear deformation. Among the series, Be<sub>3</sub>PF<sub>3</sub> demonstrates the greatest C<sub>11</sub> value, indicating increased bond strength and diminished compressibility, in accordance with its exceptional mechanical stiffness. The Voigt–Reuss–Hill (VRH) averaging approach was utilized to achieve a comprehensive characterization of polycrystalline mechanical behavior<sup>53</sup>. This approach reconciles the Voigt model (uniform strain) and the Reuss model (uniform stress), yielding a realistic estimation of the effective mechanical moduli. The bulk modulus (B), shear modulus (G), and Young's modulus (Y) were calculated using Hill's Eqs. (3–5)<sup>54,55</sup>.

$$B = \frac{B_v + B_R}{2} \quad (3)$$

$$G = \frac{G_v + G_R}{2} \quad (4)$$

$$Y = \frac{9BG}{3B + G} \quad (5)$$

Consequently, Poisson's ratio (ν) was obtained from the subsequent Eq. (6).

$$\nu = \frac{(3B - 2G)}{2(3B + G)} \quad (6)$$

The B measures a material's resistance to volumetric change under hydrostatic pressure, whereas the G assesses its resistance to deformation under applied shear forces. Young's modulus (Y), a critical indicator of stiffness, denotes the ratio of tensile stress to strain in the elastic range and is inversely related to thermal shock resistance. The computed findings in Fig. 4b, indicate that Be<sub>3</sub>PF<sub>3</sub> demonstrates the greatest values of B, G, and Y, succeeded by Be<sub>3</sub>AsF<sub>3</sub> and Be<sub>3</sub>SbF<sub>3</sub>. This signifies that the Be<sub>3</sub>PF<sub>3</sub> compound is mechanically the most resilient, exhibiting the highest resistance to deformation and compression. The increased moduli of Be<sub>3</sub>PF<sub>3</sub> indicates more robust directional bonding and reduced Be-F and Be-P bond lengths, hence augmenting lattice rigidity. In contrast, when the M-site atom enlarges (from P to when to Sb), the moduli diminish due to bond elongation and diminished charge localization, indicating fewer interatomic interactions and a marginal improvement in flexibility. Elevated values of B, G, and Y indicate that these materials possess structural integrity and mechanical stability, rendering them ideal for high-stress or high-temperature applications, including optoelectronic devices, dielectric substrates, and coating materials in solar energy systems. Pugh's ratio (B/G) and Poisson's ratio (ν), which are critical parameters for differentiating ductile from brittle materials. A pivotal Pugh's ratio threshold of 1.75 distinguishes ductile (B/G > 1.75) from brittle (B/G < 1.75) behavior<sup>56</sup>. The computed B/G ratios for Be<sub>3</sub>PF<sub>3</sub> (1.824), Be<sub>3</sub>AsF<sub>3</sub> (1.858), and Be<sub>3</sub>SbF<sub>3</sub> (1.951) surpass this threshold, so affirming that the complete Be<sub>3</sub>MF<sub>3</sub> series demonstrates ductility. Be<sub>3</sub>SbF<sub>3</sub> exhibits the highest ratio, indicating superior plastic deformability and enhanced resistance to fracture under mechanical stress. This ductility indicates that these materials can experience moderate elastic–plastic deformation prior to failure, providing robustness against stress building during thin-film development and device assembly. This inherent mechanical compliance is beneficial in the manufacturing of perovskite solar cells, where films frequently undergo stress due to thermal annealing, substrate mismatch, and multi-layer stacking. Ductile absorber layers exhibit reduced vulnerability to microcrack development and interfacial delamination, therefore preserving film continuity, grain integrity, and electrical stability throughout prolonged operation<sup>57</sup>. Moreover, the mechanical flexibility inherent in ductility facilitates possible incorporation into flexible or wearable optoelectronic systems, permitting Be<sub>3</sub>MF<sub>3</sub> films to endure bending, rolling, and mechanical vibrations without structural deterioration. Poisson's ratio (ν) offers

additional understanding of the bonding and elastic properties of these materials. Generally,  $\nu > 0.26$  signifies ductility and primarily ionic bonding, while  $\nu < 0.26$  denotes brittleness and covalent bonding. The calculated Poisson's ratios for  $\text{Be}_3\text{PF}_3$ ,  $\text{Be}_3\text{AsF}_3$ , and  $\text{Be}_3\text{SbF}_3$  all surpass 0.26, validating their ductile characteristics and affirming a predominantly ionic bonding structure that enhances elastic flexibility. The simultaneous presence of elevated  $B/G$  and  $\nu$  values confirms that the  $\text{Be}_3\text{MF}_3$  family integrates mechanical strength with processing suitability. This ductility improves compatibility with solution-based and vacuum deposition methods, including spin coating, sputtering, and thermal evaporation, which frequently create localized stress gradients during solvent evaporation and film cooling<sup>58</sup>. A supplementary metric of mechanical significance is the Zener anisotropy index ( $A$ ), defined in Eq. (7).

$$A = \frac{2C_{44}}{C_{11} - C_{12}} \quad (7)$$

The calculated  $A$  values for  $\text{Be}_3\text{MF}_3$  significantly diverge from unity, indicating considerable elastic anisotropy across crystallographic directions. Fig. S1 (Supplementary Information) illustrates that the deformed 2D and 3D elastic contour plots exhibit direction-dependent stiffness resulting from fluctuations in atomic bonding and internal stress distribution within the cubic lattice. Mechanical analysis verifies that all compounds are stable, moderately rigid, and ductile, with  $\text{Be}_3\text{PF}_3$  demonstrating superior stiffness and  $\text{Be}_3\text{SbF}_3$  displaying enhanced ductility. The combination of these features, along with robust ionic bonding and anisotropic elasticity, renders  $\text{Be}_3\text{MF}_3$  materials advantageous for structural, optoelectronic, and thermomechanical applications especially in photovoltaic encapsulation, radiation-resistant optics, and high-frequency device substrates.

### SCAPS 1D result

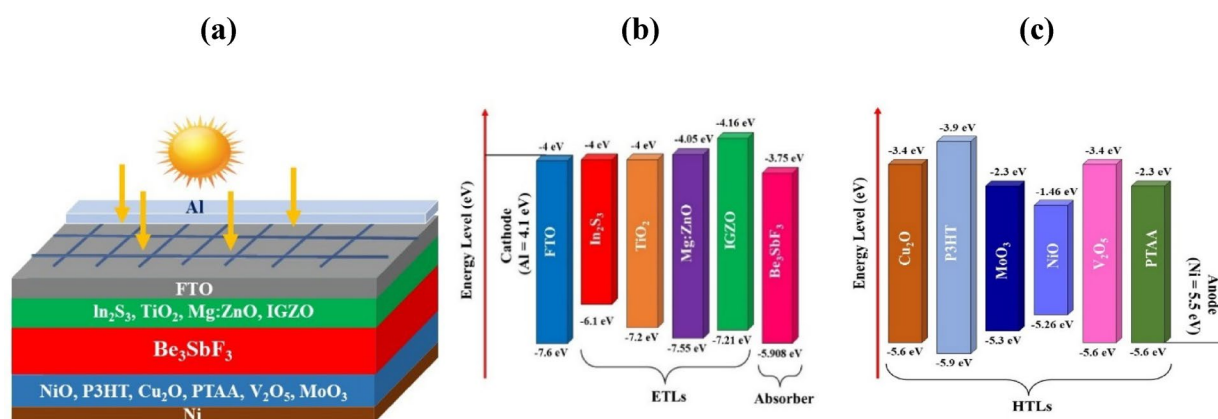
#### *Solar cell performance evolution using $\text{Be}_3\text{SbF}_3$ as absorber material*

Figure 5a–c depicts the fundamental structure and band alignment of a lead-free PSC, which includes ETL, HTL, and a back contact, incorporating a  $\text{Be}_3\text{SbF}_3$  absorber layer. The device utilizes an n-i-p configuration, wherein the intrinsic  $\text{Be}_3\text{SbF}_3$  layer creates an extensive depletion region that improves sensitivity to long-wavelength photons. This facilitates the effective creation and separation of electron–hole pairs both inside and beyond the depletion region, enhancing quantum efficiency and charge collection<sup>59</sup>.

Moreover,  $\text{Be}_3\text{SbF}_3$  proficiently restricts charge carriers and photons, augmenting photon absorption and establishing ohmic contacts with the doped HTL and ETL layers, so improving device performance. This device utilizes  $\text{Cu}_2\text{O}$ , P3HT, NiO,  $\text{V}_2\text{O}_5$ ,  $\text{MoO}_3$ , and PTAA as HTLs, with Ni functioning as the rear metal contact and Al as the front metal contact. The employed ETLs are  $\text{In}_2\text{S}_3$ ,  $\text{TiO}_2$ ,  $\text{Mg:ZnO}$ , and IGZO, with  $\text{Be}_3\text{SbF}_3$  serving as the absorber layer. The comprehensive configuration is  $\text{Al}/\text{FTO}/\text{ETL}/\text{Be}_3\text{SbF}_3/\text{HTL}/\text{Ni}$ . Tables 2 and 3 delineate the simulation input parameters for the absorber layer, ETLs, HTLs, and fluorine-doped tin oxide (FTO), whereas Table 4 presents the values for interfacial defect layers. SCAPS-1D simulations were conducted to assess the efficacy of various PSC topologies under standard operating conditions, which included the AM1.5G solar spectrum, a frequency of 1 MHz, and an ambient temperature of 300 K.

#### *Optimization of ETL and HTL*

The photovoltaic efficacy of  $\text{Be}_3\text{SbF}_3$ -based solar cells is significantly determined by the choice and interaction of appropriate electron transport layers (ETLs) and hole transport layers (HTLs), which directly affect charge extraction, recombination dynamics, and overall device efficiency. This study extensively investigated four ETLs such as  $\text{TiO}_2$ ,  $\text{In}_2\text{S}_3$ ,  $\text{Mg:ZnO}$ , and IGZO alongside six HTLs as  $\text{Cu}_2\text{O}$ , P3HT, NiO,  $\text{V}_2\text{O}_5$ ,  $\text{MoO}_3$ , and PTAA to determine ideal configurations for  $\text{Be}_3\text{SbF}_3$  absorbers. ETL enables electron flow from the absorber to the front electrode while preventing hole recombination, whereas the HTL effectively extracts holes toward the rear contact, providing balanced charge transport and minimizing recombination losses<sup>65</sup>. Figure 6a–d demonstrates that the PV characteristics:  $V_{\text{OC}}$ ,  $J_{\text{SC}}$ , FF, and PCE exhibit considerable variation depending on



**Fig. 5.** Designing  $\text{Be}_3\text{SbF}_3$ -based PSCs: (a) device configuration and band alignment with various (b) ETL and (c) HTL materials.

Parameter (unit)	FTO	In <sub>2</sub> S <sub>3</sub>	TiO <sub>2</sub>	Mg:ZnO	IGZO	Be <sub>3</sub> SbF <sub>3</sub>
N (μm)	0.1	0.05	0.05	0.05	0.05	1.0
E <sub>g</sub> (eV)	3.6	2.1	3.2	3.5	3.05	2.158
χ (eV)	4	4	4	4.05	4.16	3.75
ε	9	13.5	9	8.5	10	5.45
N <sub>v</sub> (cm <sup>-3</sup> )	1.8 × 10 <sup>19</sup>	4 × 10 <sup>13</sup>	1.8 × 10 <sup>19</sup>	1.8 × 10 <sup>19</sup>	5 × 10 <sup>18</sup>	7.18 × 10 <sup>18</sup>
N <sub>c</sub> (cm <sup>-3</sup> )	2.2 × 10 <sup>18</sup>	1.8 × 10 <sup>19</sup>	2.2 × 10 <sup>18</sup>	2.2 × 10 <sup>18</sup>	5 × 10 <sup>18</sup>	4.17 × 10 <sup>18</sup>
μ <sub>n</sub> (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	100	400	20	100	15	80
μ <sub>h</sub> (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	25	210	10	25	0.1	25
N <sub>A</sub> (cm <sup>-3</sup> )	0	0	0	0	0	1 × 10 <sup>17</sup>
N <sub>D</sub> (cm <sup>-3</sup> )	1 × 10 <sup>19</sup>	1 × 10 <sup>18</sup>	1 × 10 <sup>18</sup>	1 × 10 <sup>18</sup>	1 × 10 <sup>18</sup>	0
N <sub>t</sub> (cm <sup>-3</sup> )	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>14</sup>

**Table 2.** Essential input parameters for FTO, ETLs, and Be<sub>3</sub>SbF<sub>3</sub> absorber layers in PSCs<sup>51,60,61,62</sup>.

Parameter (unit)	Cu <sub>2</sub> O	P3HT	NiO	V <sub>2</sub> O <sub>5</sub>	MoO <sub>3</sub>	PTAA
N (μm)	0.1	0.1	0.1	0.1	0.1	0.1
E <sub>g</sub> (eV)	2.2	2	3.8	2.2	3	3.3
χ (eV)	3.4	3.9	1.46	3.4	2.3	2.3
ε	7.5	3	10.7	8	18	9
N <sub>v</sub> (cm <sup>-3</sup> )	1 × 10 <sup>19</sup>	1.8 × 10 <sup>19</sup>	1 × 10 <sup>19</sup>	5 × 10 <sup>20</sup>	2.2 × 10 <sup>18</sup>	1 × 10 <sup>21</sup>
N <sub>c</sub> (cm <sup>-3</sup> )	2 × 10 <sup>19</sup>	2.5 × 10 <sup>18</sup>	2.8 × 10 <sup>19</sup>	9.2 × 10 <sup>19</sup>	1 × 10 <sup>19</sup>	1 × 10 <sup>21</sup>
μ <sub>n</sub> (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	200	10,000	100	150	210	1
μ <sub>h</sub> (cm <sup>2</sup> V <sup>-1</sup> s <sup>-1</sup> )	8600	10,000	43.9	100	210	40
N <sub>A</sub> (cm <sup>-3</sup> )	1 × 10 <sup>20</sup>	1 × 10 <sup>20</sup>	1 × 10 <sup>20</sup>	1 × 10 <sup>20</sup>	1 × 10 <sup>20</sup>	1 × 10 <sup>20</sup>
N <sub>D</sub> (cm <sup>-3</sup> )	0	0	0	0	0	0
N <sub>t</sub> (cm <sup>-3</sup> )	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>	1 × 10 <sup>15</sup>

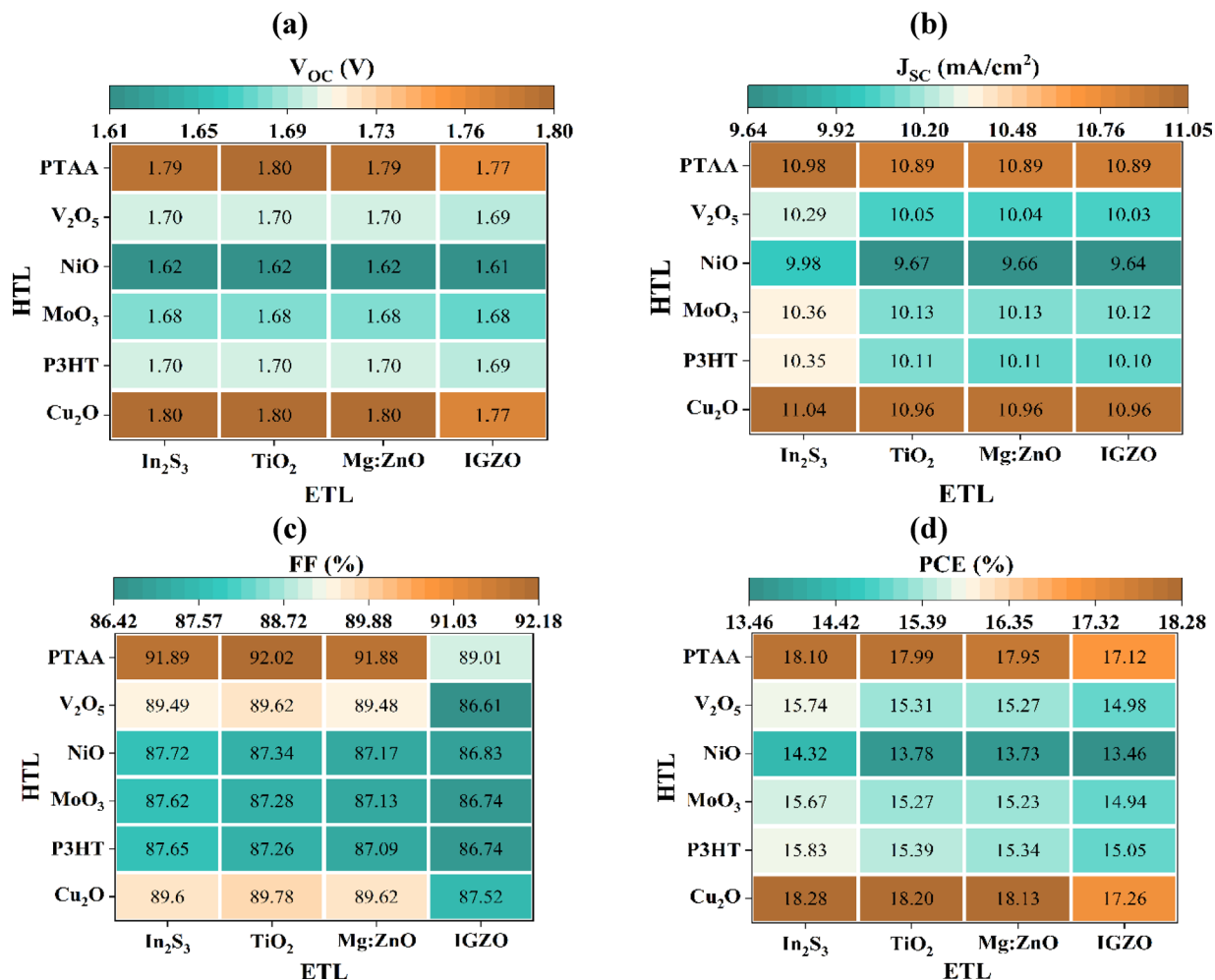
**Table 3.** Key input parameters of diverse HTLs in PSCs<sup>63,64</sup>.

Parameter	Be <sub>3</sub> SbF <sub>3</sub> /ETL	HTL/Be <sub>3</sub> SbF <sub>3</sub>
Type of defect	Neutral	
E <sub>r</sub> (eV)	0.6	
Energetic distribution	Single	
σ <sub>e</sub> /σ <sub>h</sub> (cm <sup>2</sup> )	1 × 10 <sup>-19</sup>	
Total density (cm <sup>-2</sup> )	1 × 10 <sup>11</sup>	

**Table 4.** Investigating primary interface parameters in PSCs.

the combinations of ETL and HTL. Among the ETLs, TiO<sub>2</sub> attained the highest V<sub>OC</sub> (about 1.80 V), whilst Mg:ZnO exhibited the lowest (about 1.61 V). In<sub>2</sub>S<sub>3</sub> and IGZO yielded intermediate values, so validating their stable interfacial characteristics. For J<sub>SC</sub>, most devices displayed values ranging from 9.65 to 11.04 mA/cm<sup>2</sup>, with IGZO demonstrating a little decreased current density attributable to restricted carrier mobility and interface discrepancies. The FF varied from 86.61% (IGZO) to 91.89% (In<sub>2</sub>S<sub>3</sub>), with In<sub>2</sub>S<sub>3</sub> identified as the most efficient electron transport layer because of its enhanced charge extraction and reduced interface recombination. Table S1 (Supplementary Information) illustrates the four most effective device architectures. The Al/FTO/In<sub>2</sub>S<sub>3</sub>/Be<sub>3</sub>SbF<sub>3</sub>/Cu<sub>2</sub>O/Ni design achieved the best PCE of 18.28%, surpassing alternative combinations.

The Cu<sub>2</sub>O HTL demonstrated exceptional efficacy across all ETLs, owing to its good band alignment, superior p-type conductivity, and chemical stability<sup>66</sup>. In contrast, NiO exhibited comparatively lower power conversion efficiencies because of suboptimal energy-level alignment. The findings underscore the necessity of sustaining optimal HTL thickness and band alignment to reduce recombination and improve carrier transmission. The integration of In<sub>2</sub>S<sub>3</sub> as ETL and Cu<sub>2</sub>O as HTL establishes an optimal interface design, delivering a synergistic equilibrium of voltage, current, and efficiency. This underscores Be<sub>3</sub>SbF<sub>3</sub> promise as a durable lead-free absorber for advanced optoelectronic and PV applications.



**Fig. 6.** Assessment of charge transport layers on critical performance metrics: (a)  $V_{OC}$ , (b)  $J_{SC}$ , (c) FF, and (d) PCE in  $Be_3SbF_3$  SCs.

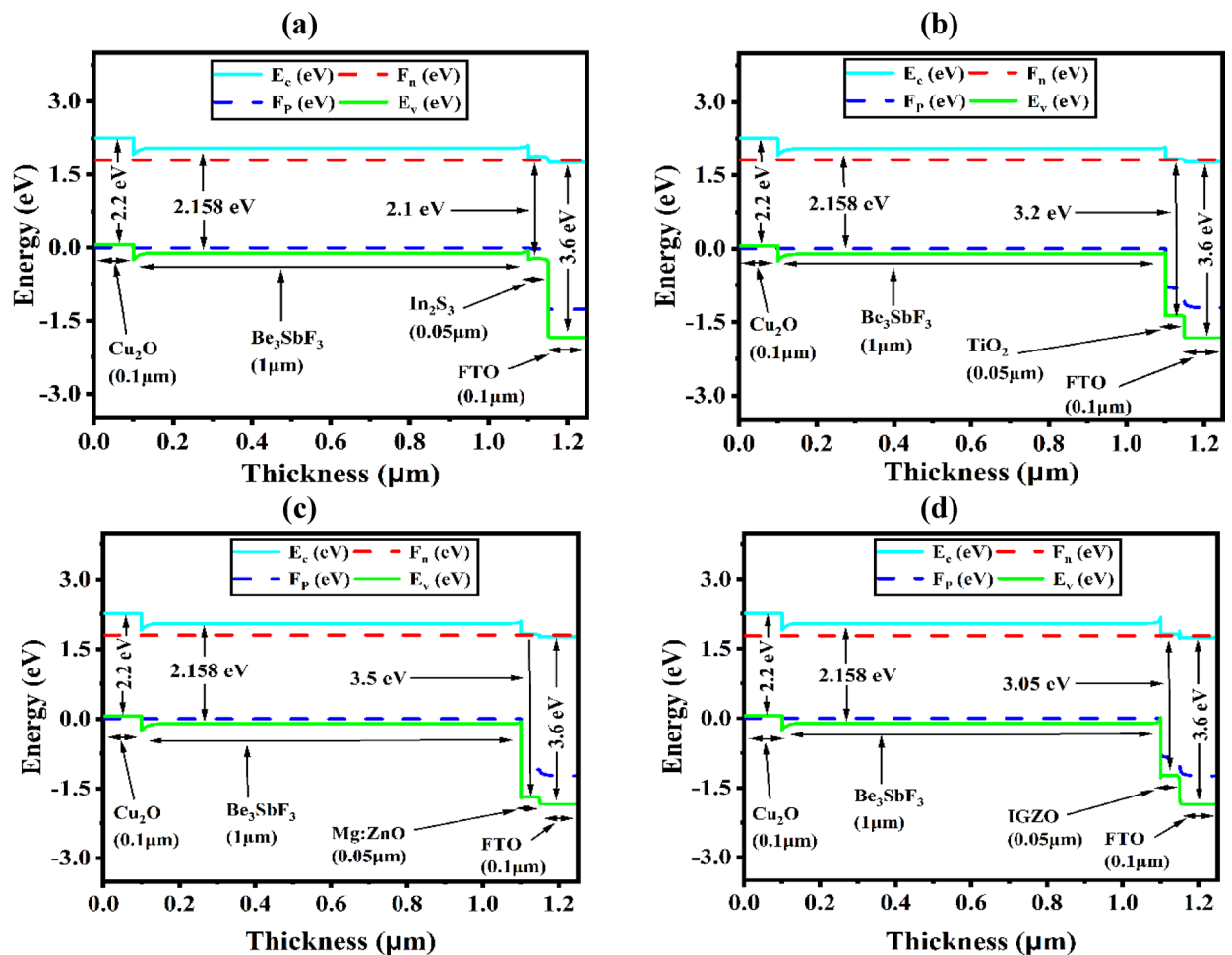
#### Role of back contact work function in determining solar cell characteristics

Back contact is crucial in influencing solar cell efficiency and overall performance. It must minimize resistive losses to ensure efficient current extraction, as inadequate contact elevates series resistance and diminishes power production<sup>67</sup>. Moreover, robust adhesion and chemical stability with the semiconductor are essential for the longevity of the device. Poor design may result in shade and restrict the effective area for energy conversion. An improved back contact enhances charge carrier collection, reduces recombination losses, and consequently improves the performance and dependability of the solar cell. The optimal metal for back contact is primarily determined by its work function (WF), which varies from 4.81 eV to 5.9 eV, as detailed in Table S2. Fig. S2 depicts the PV parameters acquired with diverse back contacts for distinct ETLs. The findings demonstrate that a reduction in WF of the Fe contact corresponds to a decline in the PCE of the SC. An elevation in WFs of metals like Au, Ni, Pd, Pt, and Se demonstrates a consistent positive link with improved PCE. This enhancement is primarily ascribed to the diminution of the Schottky barrier height via improved alignment of the metal work function with the Fermi level of the HTL<sup>68</sup>. Ni is the most appropriate back contact material among these metals, owing to its ideal work function, cost efficiency, plentiful availability, and consistent performance, rendering it a viable option for efficient and scalable device manufacturing.

#### Energy band diagram (EBD) of $Be_3SbF_3$

The energy band diagram (EBD) illustrates the conduction and valence band offsets ( $E_C$  and  $E_V$ ) between the  $Be_3SbF_3$  absorber and the HTL, influenced by various ETLs in conjunction with  $Cu_2O$ . Optimal energy level alignment at the ETL/ $Be_3SbF_3$  and  $Be_3SbF_3$ /HTL interfaces is essential for maximizing device efficiency. Misalignment may result in charge accumulation and interfacial recombination, thereby diminishing performance. For efficient operation, ETL must exhibit a higher electron affinity than  $Be_3SbF_3$  to enable effective electron extraction, whereas HTL should demonstrate a lower ionization energy to promote hole transport<sup>69</sup>. Therefore, optimizing the interfacial electronic properties of both layers is crucial for improving overall PV efficiency. Figure 7a–d illustrates the EBDs of  $Be_3SbF_3$ -based solar cell structures that utilize  $Cu_2O$  as HTL alongside four different ETLs. The configuration employing  $In_2S_3$  as ETL shows a clear intersection of the



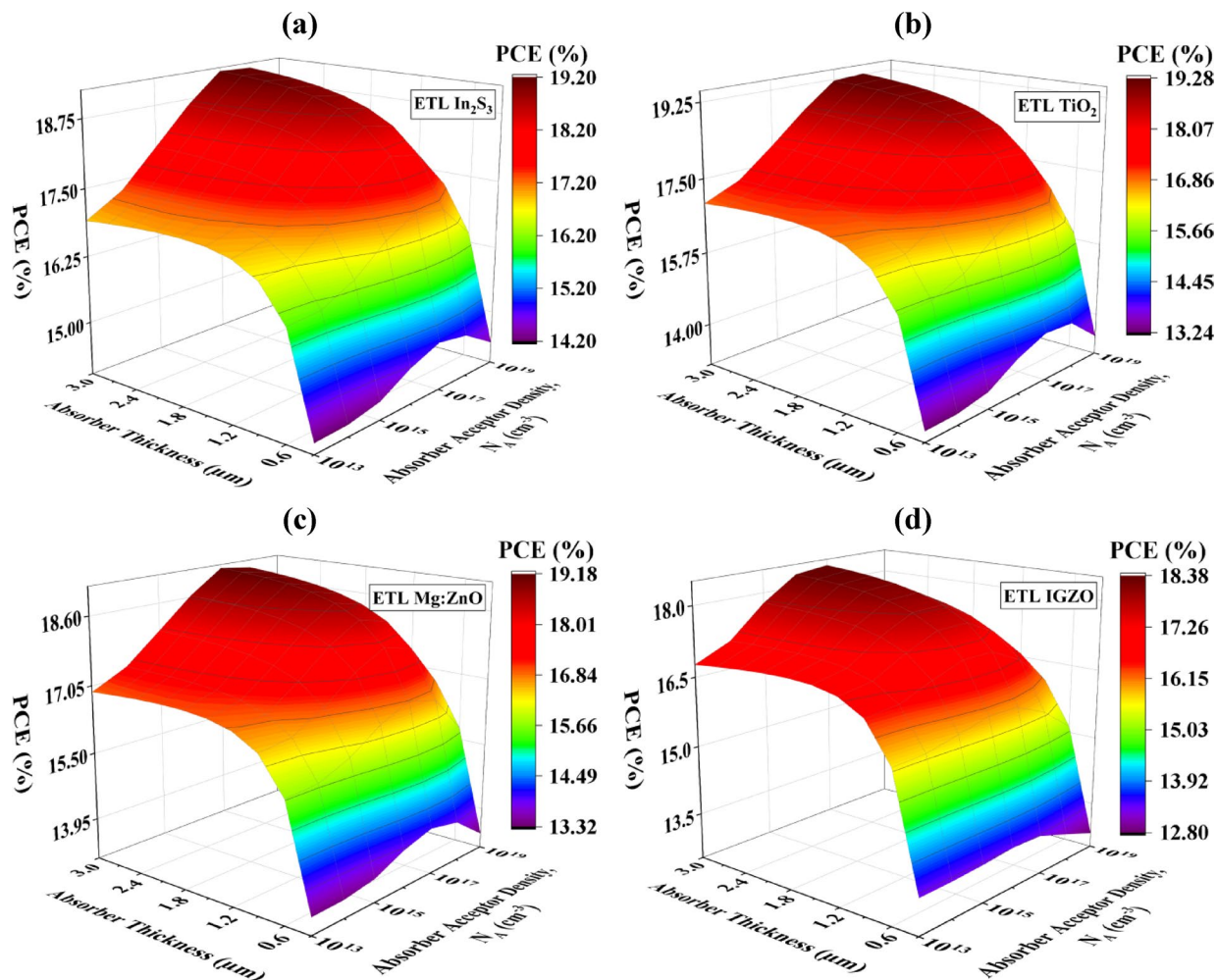


**Fig. 7.** Investigating variations in band structure of SCs with diverse ETLs: (a)  $\text{In}_2\text{S}_3$ , (b)  $\text{TiO}_2$ , (c)  $\text{Mg:ZnO}$ , and (d)  $\text{IGZO}$ .

Fermi level with the  $E_c$ . The ETLs utilized in these devices, ordered by ascending bandgap energy, include  $\text{In}_2\text{S}_3$  (2.1 eV),  $\text{TiO}_2$  (3.2 eV),  $\text{IGZO}$  (3.05 eV), and  $\text{Mg:ZnO}$  (3.5 eV). The quasi-Fermi level for holes ( $F_p$ ) is consistently positioned above  $E_v$ , whereas the quasi-Fermi level for electrons ( $F_n$ ) is closely aligned with  $E_c$ , demonstrating effective charge carrier separation in all configurations.

#### Role of acceptor density and absorber thickness variations in SC performance

The dynamics of charge carriers, recombination rates, and efficiency in PSCs are significantly influenced by the acceptor density ( $N_A$ ) and the thickness of the absorber layer. An optimal  $N_A$  enhances the built-in electric field, facilitating charge separation and minimizing recombination, while values that are too low or too high may impede carrier collection or elevate trap-assisted losses<sup>70</sup>. Thicker absorber layers enhance light absorption but may elevate recombination and resistance, while thinner layers reduce recombination but limit photocurrent<sup>71</sup>. Attaining an appropriate equilibrium of  $N_A$  and thickness is essential for maximizing PCE, achieved by interface engineering, doping regulation, and defect passivation to enhance stability and performance. Figure 8a–d illustrates the synergistic impact of  $N_A$  and absorber thickness on the PCE of the SCs. This study altered  $N_A$  from  $10^{13}$  to  $10^{19} \text{ cm}^{-3}$  and adjusted absorber thickness between 0.3 and 3.0  $\mu\text{m}$ . The findings reveal a notable fluctuation in PCE corresponding to alterations in these two parameters across four distinct SCs including FTO/ETL/ $\text{Be}_3\text{SbF}_3$ / $\text{Cu}_2\text{O}$ /Ni configurations, utilizing ETLs of  $\text{In}_2\text{S}_3$ ,  $\text{TiO}_2$ ,  $\text{Mg:ZnO}$ , and  $\text{IGZO}$ . The PCE is generally reduced when the  $N_A$  is beneath  $10^{17} \text{ cm}^{-3}$ . When the  $N_A$  attains  $10^{17} \text{ cm}^{-3}$ , the PCE hits its zenith, signifying that elevated doping concentrations result in augmented recombination losses. As  $N_A$  surpasses  $10^{17} \text{ cm}^{-3}$ , the PCE progressively diminishes. Likewise, for absorber thicknesses between 0.5 and 1.2  $\mu\text{m}$ , the PCE stays elevated but diminishes for thicker absorbers (1.5 to 3.0  $\mu\text{m}$ ) due to heightened bulk recombination and series resistance. Among the studied ETLs, the  $\text{In}_2\text{S}_3$ -based design (Fig. 8a) exhibited the highest PCE of around 18.28%, indicating effective charge transport and decreased recombination. The  $\text{TiO}_2$ -based (Fig. 8b) and  $\text{Mg:ZnO}$ -based (Fig. 8c) SCs had PCE values of 18.20% and 18.15%, respectively, indicating a close similarity. The  $\text{IGZO}$ -based SC (Fig. 8d) exhibited a marginally reduced efficiency of around 17.26%, presumably attributable to fluctuations in carrier extraction characteristics. The results emphasize the importance of absorber thickness,  $N_A$ , and the selection of ETL in optimizing PCE. Achieving a balance between low  $N_A$  and moderate thickness is



**Fig. 8.** Exploration of PCE fluctuation in relation to absorber layer thickness and  $N_A$  for SCs utilizing various ETLs: (a)  $\text{In}_2\text{S}_3$ , (b)  $\text{TiO}_2$ , (c)  $\text{Mg:ZnO}$ , and (d)  $\text{IGZO}$ .

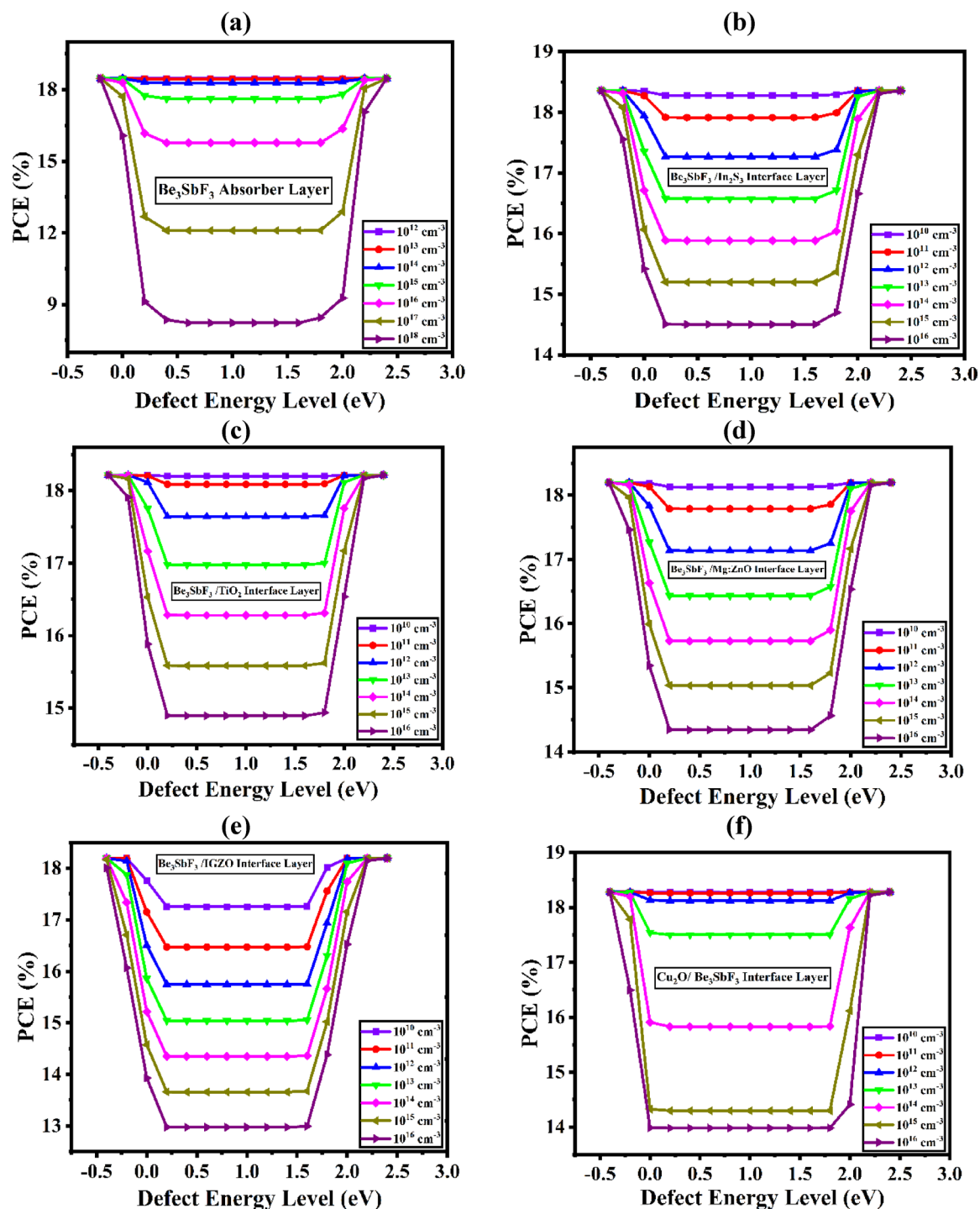
essential for minimizing recombination and enhancing charge collection efficiency. Among the assessed ETLs,  $\text{In}_2\text{S}_3$  exhibited exceptional performance in attaining elevated PCE, establishing it as a viable alternative for the advancement of high-efficiency and robust PSC designs.

#### Bulk and interfacial defect investigation

Organic and inorganic devices are significantly impacted by bulk defects, which are a significant factor contributing to their lower performance. The densities of defect levels are significantly influenced by experimental processing parameters; nonetheless, evaluating device performance at different defect energy levels and concentrations offers a useful comparison framework for optimizing these conditions<sup>72</sup>. Figure 9a–f depicts the fluctuation in solar output parameters (PCE) relative to defect energy and concentration. In the bulk defect study of the absorber layer, the defect energy level for  $\text{Be}_3\text{SbF}_3$  was modified from  $-0.2$  to  $2.4$  eV, and the defect concentration ( $N_t$ ) was altered between  $10^{12}$  and  $10^{18} \text{ cm}^{-3}$ .

The defect energy range ( $-0.2$  to  $2.4$  eV) encompasses both shallow and deep-level trap states, facilitating a thorough examination of defect-assisted recombination in  $\text{Be}_3\text{SbF}_3$ . Shallow traps near the band boundaries (about  $-0.2$  eV) induce minor carrier trapping, whereas deep traps (approximately  $2.4$  eV) function as non-radiative recombination centers that diminish carrier lifespan and voltage. This range aligns with prior research by Reza et al.<sup>73</sup> and others on lead-free perovskites ( $\text{CaRbCl}_3$ ,  $\text{Ca}_3\text{NCl}_3$ ), which indicated that deep-level defects ( $>2$  eV) substantially impair charge collection and solar efficiency<sup>74</sup>. Figure 9a illustrates the associated defect energy levels and  $N_t$  for  $\text{Be}_3\text{SbF}_3$ . The data indicates that elevated  $N_t$  can significantly impact device performance. Decreasing  $N_t$  to  $10^{14} \text{ cm}^{-3}$  yields an enhancement of device performance by as much as 40%. Nonetheless, attaining such a minimal  $N_t$  is still impractical within the limitations of conventional thin-film deposition methods. Consequently,  $N_t$  of  $10^{15} \text{ cm}^{-3}$  is deemed a more plausible concentration for this investigation.

Interface defects significantly affect device performance. Figure 9b–f illustrates the examination of defect concentrations at the ETL/absorber and  $\text{Cu}_2\text{O}$ /absorber interfaces. The device's performance is analyzed at different defect concentrations and energy levels. Defect concentrations were changed from  $10^{10}$  to  $10^{16} \text{ cm}^{-3}$ , whereas defect energy levels spanned from  $-0.4$  to  $2.4$  eV. The findings demonstrate that defects at the HTL/



**Fig. 9.** Performance Analysis of SCs in various layer configurations with varying defect position and density: (a)  $\text{Be}_3\text{SbF}_3$ , (b)  $\text{Be}_3\text{SbF}_3/\text{In}_2\text{S}_3$ , (c)  $\text{Be}_3\text{SbF}_3/\text{TiO}_2$ , (d)  $\text{Be}_3\text{SbF}_3/\text{Mg:ZnO}$ , (e)  $\text{Be}_3\text{SbF}_3/\text{IGZO}$ , and (f)  $\text{Cu}_2\text{O}/\text{Be}_3\text{SbF}_3$ .

absorber interface adversely impact device performance more significantly than those at the ETL/absorber interface. The examination of bulk and interfacial defects elucidates the impact of different defects on overall device performance. In the device bulk layer, defects with energy levels ranging from 0.3 to 2 eV are categorized as deep defects, which adversely impact performance and should be eliminated. Conversely, defects at the absorber/HTL interface, exhibiting energy levels between 0 and 2 eV, are also classified as deep defects. The severity of defects in PV performance is contingent upon the location, concentration, and energy level of each defect, highlighting the necessity of meticulous defect control to enhance device efficiency.

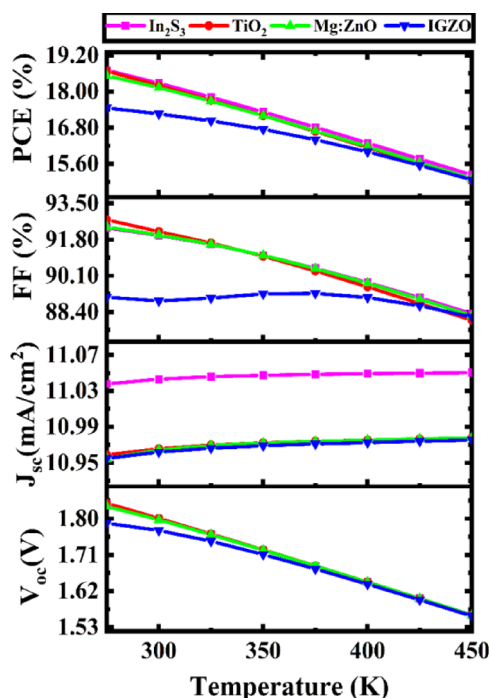
### Impact of temperature on PV efficiency

Temperature is crucial in influencing the operational stability and performance of PSCs. Although PSCs provide high efficiency and manufacturing simplicity, their large-scale production encounters obstacles, especially from environmental conditions including UV radiation, temperature variations, wetness, and humidity. Temperature is particularly crucial, as solar cells often operate in direct sunshine, where temperatures can exceed ambient values by as much as 45 °C<sup>75</sup>. An elevation in temperature results in a reduction of the semiconductor bandgap owing to the expansion of interatomic distances inside the material. The working temperature fluctuates for PSCs extends from 275 to 450 K, allowing for the evaluation of temperature fluctuations on cell performance. Figure 10 illustrates that temperature increase consistently diminishes critical metrics, such as  $V_{OC}$ , FF, and PCE, across all cell configurations. Notably, the  $J_{SC}$  exhibits remarkable stability despite rising temperatures. The noted reduction in  $V_{OC}$  is chiefly due to its inverse correlation with the saturation current density ( $J_0$ ), which demonstrates a significant increase with rising temperature. The temperature-induced increase in  $J_0$  arises from heightened carrier recombination rates in the absorber layer and at the interface regions<sup>76</sup>. The PCE and FF often decrease with elevated temperature, mainly due to the increase in series resistance and augmented non-radiative recombination losses, both of which impede efficient charge transfer and extraction<sup>77</sup>. Conversely, the IGZO-based device exhibits a nearly constant FF, indicating its enhanced thermal and electrical stability compared to alternative ETLs. This stability arises from IGZO enhanced electron mobility, minimal intrinsic defect density, and extensive bandgap (~3.05 eV), which together inhibit thermally induced carrier leakage and maintain the built-in electric field at high temperatures<sup>78</sup>.

Moreover, IGZO offers a properly oriented conduction band and a chemically stable interface with the perovskite absorber, mitigating temperature-induced interfacial deterioration. Nonetheless, as temperature escalates, the concomitant increase in series resistance diminishes overall device efficiency and fill factor, resulting in a significant decline in power conversion efficiency. These findings underscore the significance of heat management in sustaining consistent performance of PSCs under practical operating situations.

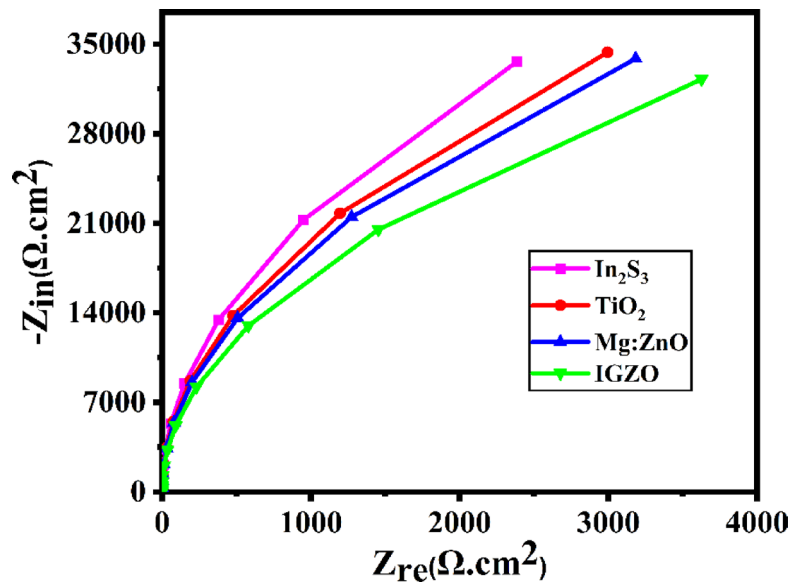
### Nyquist plot analysis

Impedance analysis is an effective method for examining the properties of perovskite materials, especially those containing different halides. This method is recognized as one of the most efficacious techniques for investigating the behavior of perovskite materials with various halide compositions<sup>79</sup>. To enhance comprehension of the effects of different ETLs on critical PV parameters, including  $V_{OC}$ ,  $J_{SC}$ , FF, and PCE obtained from the I-V characteristics, complex impedance ( $Z^*$ ) was utilized to investigate the influence of ion migration and charge carrier dynamics within the perovskite solar cell. Figure 11 displays the Nyquist plot of complex impedance, depicting the correlation between the imaginary component ( $Z_{im}$ ) and the real component ( $Z_{re}$ ) of impedance for various ETL materials. The X-axis denotes recombination resistance, but the Y-axis illustrates geometrical capacitance, indicating charge carrier concentration at the interface layers. Each curve in the graph approximates a semi-circle within the frequency spectrum of 100 Hz to 1 MHz. The real ( $Z_{re}$ ) and imaginary ( $Z_{im}$ ) components of the impedance are graphed against one another, revealing that the impedance spectra for each ETL exhibits a different, approximately semi-circular form. Furthermore, the diameter of the semi-circle differs among various

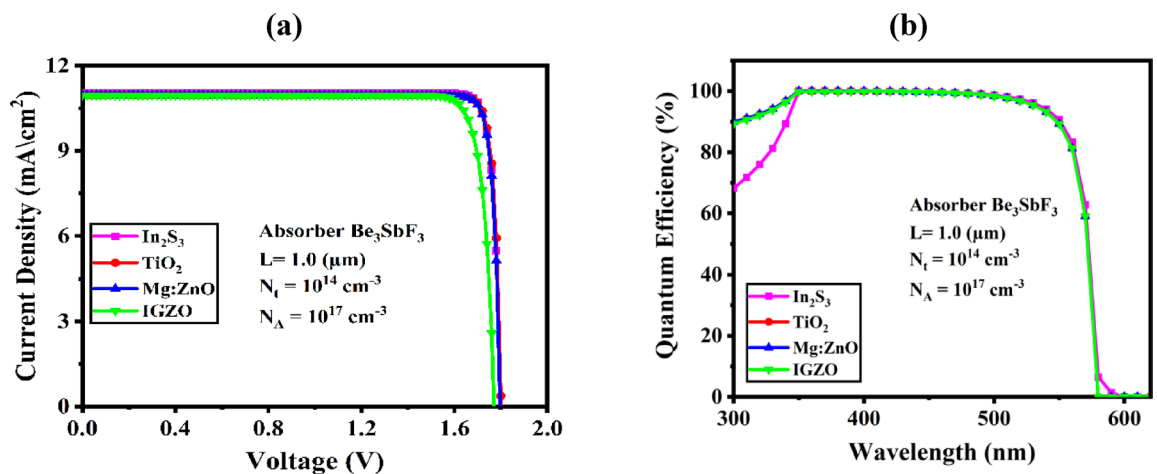


**Fig. 10.** Temperature-driven performance variations in PV cells with distinct ETLs.





**Fig. 11.** Impedance analysis of  $\text{Be}_3\text{SbF}_3$  absorber-based solar cells employing various ETL materials ( $\text{In}_2\text{S}_3$ ,  $\text{TiO}_2$ ,  $\text{Mg:ZnO}$ , and  $\text{IGZO}$ ) via Nyquist plot.



**Fig. 12.** Analysis of (a) J-V and (b) QE characteristics curve of PSCs.

ETL-based structures, signifying variations in electrical properties and potential performance consequences. The diameter of the semi-circle expands with increasing frequencies. The broader semi-circle linked to the  $\text{TiO}_2$  ETL-based PSC indicates an elevated impedance of roughly  $34,366.681 \Omega \text{ cm}^2$ . This indicates increased resistance inside the system, probably attributable to diminished charge transport kinetics or heightened charge transfer resistance at the perovskite- $\text{TiO}_2$  interface. The smaller semi-circle seen in the  $\text{IGZO}$  ETL-based structure suggests reduced resistance, approximately  $32,272.6149 \Omega \text{ cm}^2$ , implying enhanced charge transfer, superior interface qualities, or more charge carrier mobility relative to other ETL materials.

At elevated frequencies, the observed resistance indicates the recombination resistance of the material, but the capacitance at these frequencies denotes the geometric capacitance, which relates to charge accumulation at the interfaces. The low-frequency response is intricate, probably reflecting ionic mobility and hysteresis phenomena<sup>80</sup>. As ions collect and the reduction barriers diminish, the impedance response becomes inductive. The Nyquist plot of the solar cell provides a thorough qualitative insight into the device's capacitance, resistive losses, and the rate of recombination faults present inside the device. This investigation offers critical insights into the electrical characteristics and efficiency of PSCs, enhancing comprehension of the elements affecting their performance.

#### J-V and QE characteristics

Figure 12a presents the J-V characteristics of four configurations of  $\text{Be}_3\text{SbF}_3$ -based solar cells under applied voltages from 0 to 2.0 V. Initially, all devices demonstrate comparable photocurrent values of approximately

11.03 mA/cm<sup>2</sup>, except for the IGZO-based configuration, which exhibits a significantly lower current density. With an increase in voltage beyond 1.6 V, all structures exhibit a reduction in photocurrent, mainly due to increased charge carrier recombination linked to defect states in the perovskite absorber. The observations indicate that improved crystallinity and reduced defect density in the absorber layer enhance PV performance by facilitating efficient charge separation and transport.

Figure 12b illustrates the quantum efficiency (QE) spectra of Be<sub>3</sub>SbF<sub>3</sub>-based solar cells employing different ETLs across the 300–620 nm wavelength range. The In<sub>2</sub>S<sub>3</sub>-based device initiates with a reduced quantum efficiency of approximately 68%, whereas other ETL designs commence at approximately 89.9%, signifying enhanced carrier collection near the absorption edge. The QE increases significantly between 350–530 nm, exceeding 95%, indicating robust optical absorption and effective charge production. This plateau indicates superior interfacial charge transport and negligible recombination losses at optimal photon energy. Nonetheless, QE diminishes beyond 590 nm owing to non-radiative recombination and free-carrier absorption in the doped front and ETLs phenomena characteristic of wide-bandgap perovskite and oxide-based devices<sup>81</sup>. The spectral variations among ETLs result from disparities in band alignment, carrier mobility, and interface defect density, which significantly affect charge extraction performance. The significant quantum efficiency in the visible spectrum validates the robust optoelectronic coupling and effective photo response of Be<sub>3</sub>SbF<sub>3</sub>, underscoring its potential as a lead-free perovskite absorber for high-performance photovoltaic applications.

*An evaluation and comparison of SCAPS-1D findings with prior studies*

This study obtained remarkable performance metrics for the Al/FTO/ETL/Be<sub>3</sub>SbF<sub>3</sub>/HTL/Ni SC configuration. The study concentrated on integrating diverse ETLs, including In<sub>2</sub>S<sub>3</sub>, TiO<sub>2</sub>, Mg:ZnO, and IGZO, into the device architecture, with Cu<sub>2</sub>O functioning as the HTL. The systematic optimization of essential parameters, including layer thickness, doping, defect densities, temperature, J-V, and QE properties, was conducted. After optimizing these variables, the proposed structure exhibited exceptional PV performance, attaining an efficiency of 18.28%, a V<sub>OC</sub> of 1.7991 V, a J<sub>SC</sub> of 11.0423 mA/cm<sup>2</sup>, and an FF of 92%. Although achieving this high PV performance may provide difficulties for experimental researchers, this study offers a framework and guidance for future inquiries. The efficacy of SCs is predominantly determined by structural design, material quality, and production methodologies. Table 5 compares the best configurations documented in prior studies with the PV characteristics of the four device types introduced in this research.

**Machine learning analysis**

Machine learning (ML) has significantly transformed PV research, facilitating swift, data-driven progress in PSCs. Conventional experimental and computational approaches tend to be time-consuming and resource-demanding, while machine learning techniques effectively analyze extensive datasets to identify patterns, forecast performance parameters, and inform material and structural optimization. This capability is essential for PSCs, as efficiency relies on the complex interactions among composition, processing conditions, and device architecture. Machine learning models are essential for predicting PV parameters. Researchers optimize electron and hole transport layers, absorber materials, and interface properties by integrating machine learning

Model	Types of work	PCE (%)	V <sub>OC</sub> (V)	J <sub>SC</sub> (mA/cm <sup>2</sup> )	FF (%)	Refs.
MoSe <sub>2</sub> /Si	E	0.33	0.19	5.71	30.47	<sup>82</sup>
FTO/TiO <sub>2</sub> /Sb <sub>2</sub> Se <sub>3</sub> /Au	E	2.26	0.52	10.3	42.3	<sup>83</sup>
SLG/Mo/Sb <sub>2</sub> Se <sub>3</sub> /CdS/ZnO/ITO/Ag	E	4.22	0.453	19.8	50.6	<sup>84</sup>
TiO <sub>2</sub> /CsPbBr <sub>3</sub> /Spiro-OMETAD	E	6.08	1.307	6.820	68	<sup>85</sup>
CsPbI <sub>3</sub> /FAPbI <sub>3</sub>	E	15.60	1.22	17.26	74	<sup>86</sup>
CsPbI <sub>3</sub> /Cs <sub>0.25</sub> FA <sub>0.75</sub> PbI <sub>3</sub>	E	17.39	1.20	18.91	76.00	<sup>87</sup>
FTO/SnS <sub>2</sub> /Ca <sub>3</sub> PBr <sub>3</sub> /Au	T	13.66	1.25	12.43	87.18	<sup>21</sup>
FTO/TiO <sub>2</sub> /Ca <sub>3</sub> AsCl <sub>3</sub> /Ni	T	17.11	1.34	14.80	86.23	<sup>88</sup>
C <sub>60</sub> /CsGeI <sub>3</sub> /CuSCN	T	17.61	1.0169	19.653	88.13	<sup>89</sup>
FTO/TiO <sub>2</sub> /RbGeI <sub>3</sub> /CuI/Au	T	18.97	0.777	33.969	71.93	<sup>90</sup>
FTO/TiO <sub>2</sub> /Cs <sub>2</sub> SnBr <sub>6</sub> /N-doped TiO <sub>2</sub> /Au	T	38.70	1.29	33.34	90.21	<sup>91</sup>
FTO/TiO <sub>2</sub> /RbGeI <sub>3</sub> /P3HT/Au	T	18.44	0.813	30.84	74.19	<sup>92</sup>
Al/FTO/SnS <sub>2</sub> /Sr <sub>3</sub> NCl <sub>3</sub> /Au	T	18.11	1.24	16.79	86.44	<sup>93</sup>
Al/FTO/SnS <sub>2</sub> /Sr <sub>3</sub> SbCl <sub>3</sub> /Ni	T	16.63	1.364	13.738	88.74	<sup>94</sup>
Al/FTO/SnS <sub>2</sub> /Ca <sub>3</sub> SbI <sub>3</sub> /Au	T	19.47	1.1631	19.07	87.79	<sup>95</sup>
Al/FTO/SnS <sub>2</sub> /Sr <sub>3</sub> SbI <sub>3</sub> /MoO <sub>3</sub> /Ni	T	17.12	0.72	29.57	80.61	<sup>96</sup>
Al/FTO/In <sub>2</sub> S <sub>3</sub> /Be <sub>3</sub> SbF <sub>3</sub> /Cu <sub>2</sub> O/Ni	T	18.28	1.7991	11.042	92	T.W
Al/FTO/TiO <sub>2</sub> /Be <sub>3</sub> SbF <sub>3</sub> /Cu <sub>2</sub> O/Ni	T	18.20	1.8009	10.963	92.18	T.W
Al/FTO/Mg:ZnO/Be <sub>3</sub> SbF <sub>3</sub> /Cu <sub>2</sub> O/Ni	T	18.13	1.7972	10.961	92.02	T.W
Al/FTO/IGZO/Be <sub>3</sub> SbF <sub>3</sub> /Cu <sub>2</sub> O/Ni	T	17.26	1.7707	1.77	88.92	T.W

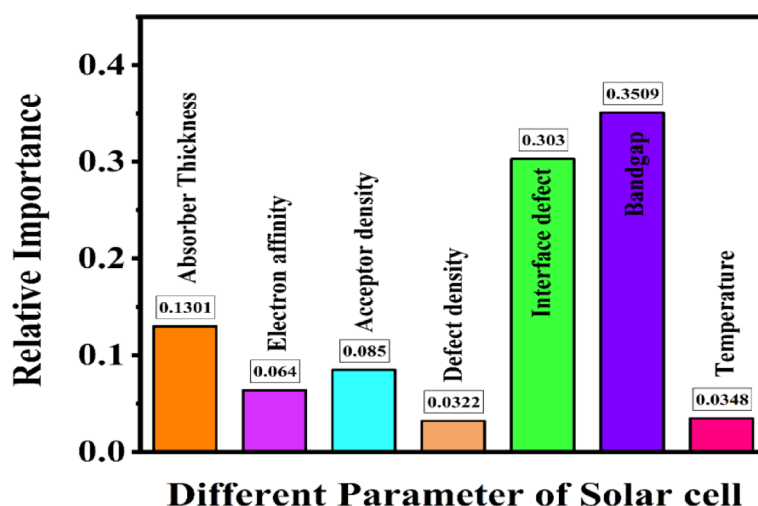
**Table 5.** A comparative assessment of PV performance in Be<sub>3</sub>SbF<sub>3</sub>-based PSCs: insights from previous research (T = Theoretical, E = Experimental, and T.W. = This Work).

with simulation tools such as SCAPS-1D to reduce losses. Research conducted by Shrivastav et al.<sup>97</sup>, Khan et al.<sup>98</sup>, and Li et al.<sup>99</sup> illustrates the efficacy of machine learning in enhancing lead-free and heterojunction PSCs, resulting in elevated predictive accuracy and increased efficiency. The Random Forest (RF) algorithm is notable among machine learning methods for its robustness, interpretability, and capacity to manage nonlinear, multivariate data. RF models demonstrate superior capability in identifying key parameters that influence PSC performance, frequently surpassing linear regression, support vector machines (SVM), and k-nearest neighbors (KNN) models. Advanced techniques, including Gradient Boosting, XGBoost, Artificial Neural Networks, and Deep Learning, enhance the capabilities of machine learning in high-throughput screening and autonomous material discovery<sup>100</sup>. The incorporation of machine learning, particularly random forests, into polymer science and engineering research signifies a transition towards intelligent and automated material optimization. The integration of statistical learning, physical modeling, and experimental insights facilitates the accelerated design of efficient, stable, and sustainable perovskite materials, thereby advancing the development of next-generation lead-free solar technologies.

#### Comprehending the RF algorithm

This study utilized RF technique as the primary machine learning model due to its remarkable proficiency in managing nonlinear interactions, high-dimensional feature spaces, and intricate correlations among PV parameters. The RF approach functions by creating a collection of many decision trees, each trained on randomly chosen subsets of data and input attributes. The ultimate forecast is derived by averaging the outputs of various trees, which successfully reduces overfitting while preserving excellent predictive accuracy. This ensemble-based methodology offers a dependable equilibrium among robustness, precision, and interpretability, rendering it particularly appropriate for evaluating the interaction of material, structural, and operational elements that collectively affect solar cell efficiency<sup>101,102</sup>. The SCAPS-1D simulator was employed to methodically alter essential device parameters, such as absorber layer thickness, electron affinity, defect and doping density, interface defect, bandgap, and operational temperature, to create a comprehensive dataset. This simulation campaign produced 13,113 data points, with 80% (10,490 samples) designated for model training and 20% (2623 samples) for testing and validation. The predictive performance of the trained RF model was assessed using the coefficient of determination ( $R^2$ ) and mean absolute error (MAE) metrics, both indicating the model's robust capacity to reliably predict PV performance based on input data. The RF algorithm's primary advantage is its capacity to quantify feature importance, therefore pinpointing the most significant elements affecting PSC efficiency. This feature analysis allows researchers to prioritize elements for further optimization. The hyperparameter tweaking process was conducted with precision to guarantee optimal model performance. The model utilized 100 decision trees, attaining an equilibrium between computational efficiency and precision. The maximum tree depth was limited to avoid overfitting, and a minimum of two samples per leaf node was mandated to guarantee significant splits. Furthermore, bootstrap sampling and feature subsampling methods were employed to incorporate randomness, improve model stability, and avert the dominance of any feature in the learning process.

Figure 13 depicts the significance of essential physical parameters affecting the performance. Each bar illustrates the contribution of a particular variable to the overall PCE and electrical attributes. Bandgap energy is the most crucial element, highlighting its pivotal function in regulating light absorption, carrier production, and voltage output. An adjusted bandgap guarantees a balance between extensive spectrum absorption and reduced thermalization losses, hence directly improving  $J_{SC}$  and  $V_{OC}$ <sup>103</sup>. The interface defect density significantly affects the system, suggesting that non-radiative recombination at the absorber/transport layer junctions is a major loss mechanism. Mitigating interface defects via surface passivation or lattice matching can significantly improve charge extraction efficiency. Parameters like absorber thickness and acceptor density exhibit mild effects, as they affect optical absorption depth and the intensity of the internal electric field. In contrast, electron



**Fig. 13.** Identifying critical parameters for the evaluation of ML models in Al/FTO/ $\text{In}_2\text{S}_3$ / $\text{Be}_3\text{SbF}_3$ / $\text{Cu}_2\text{O}$ /Ni SCs.

affinity, temperature, and bulk defect density exert relatively negligible effects under typical conditions. The research indicates that optimizing the bandgap and reducing interface defects are the most effective methods for enhancing perovskite solar cell performance.

*An exhaustive evaluation of the model's performance*

The performance of the model was meticulously assessed using a range of statistical indicators, as detailed in Table 6, to guarantee the accuracy and generalizability of the RF predictions. The  $R^2$  values indicated a significant correlation between predicted and experimental results, attaining 0.9894 for  $V_{OC}$ , 0.9881 for PCE, 0.9892 for  $J_{SC}$ , and 0.9814 for FF. The results demonstrate the model's capacity to accurately represent complex physical dependencies and nonlinear connections among PV parameters. Moreover, the reduced values recorded for Mean Squared Error (MSE), Root Mean Squared Error (RMSE), Mean Absolute Error (MAE), and Mean Absolute Percentage Error (MAPE) corroborate the model's stability and minimum divergence from experimental standards. These indications collectively yield an overall prediction accuracy of 98.7%, highlighting the robustness and trustworthiness of the RF method in correlating material attributes with device performance results. The feature selection process was crucial in improving the model's interpretability and computational efficiency by methodically eliminating redundant or weakly correlated variables. This enhancement maintained the model's predictive ability while expediting training convergence and augmenting generalization to novel data. Fig. S3a–d illustrates a nearly perfect correlation between the projected and actual data points along the ideal  $y = x$  line, affirming the model's robust consistency and minimal prediction bias. These findings confirm the RF approach as a robust and interpretable machine learning framework for quantitative prediction, sensitivity analysis, and design optimization of PSCs. Its capacity to amalgamate several performance indicators while preserving precision and stability establishes RF as a highly dependable method for directing the future development of high-efficiency, lead-free perovskite devices.

*Analysis of correlation heatmap*

Figure 14 displays a correlation heatmap that illustrates the linear correlations among several design parameters and the PCE of the PSC. Correlation coefficients between  $-1$  and  $+1$  signify the strength and direction of relationships, with positive values denoting direct associations and negative values representing inverse trends. The analysis indicates that absorber thickness demonstrates a modest positive association with PCE ( $r = 0.25$ ), implying that increased absorber thickness somewhat improves device performance. In contrast, bandgap energy exhibits a significant negative correlation ( $r = -0.65$ ), indicating that bigger bandgaps diminish efficiency. Parameters like acceptor density ( $r = -0.16$ ) and defect density ( $r = -0.11$ ) exhibit weak negative associations, suggesting that excessive doping or structural imperfections marginally impair performance. Likewise, interface imperfections exhibit a moderate negative correlation ( $r = -0.36$ ), substantiating their adverse impact on charge transport and overall efficiency.

Additional parameters such as electron affinity, dielectric permittivity, electron mobility, hole mobility, and temperature demonstrate weak or negligible correlations, indicating minimal direct impact on PCE. The heatmap indicates that bandgap energy and absorber thickness are the primary factors affecting device efficiency. The relatively weak correlations among most variables highlight their relative independence, hence validating the application of nonlinear machine learning models such as Random Forest to elucidate the intricate interdependencies influencing solar cell performance.

*SHAP-based importance analysis for comprehending feature contributions*

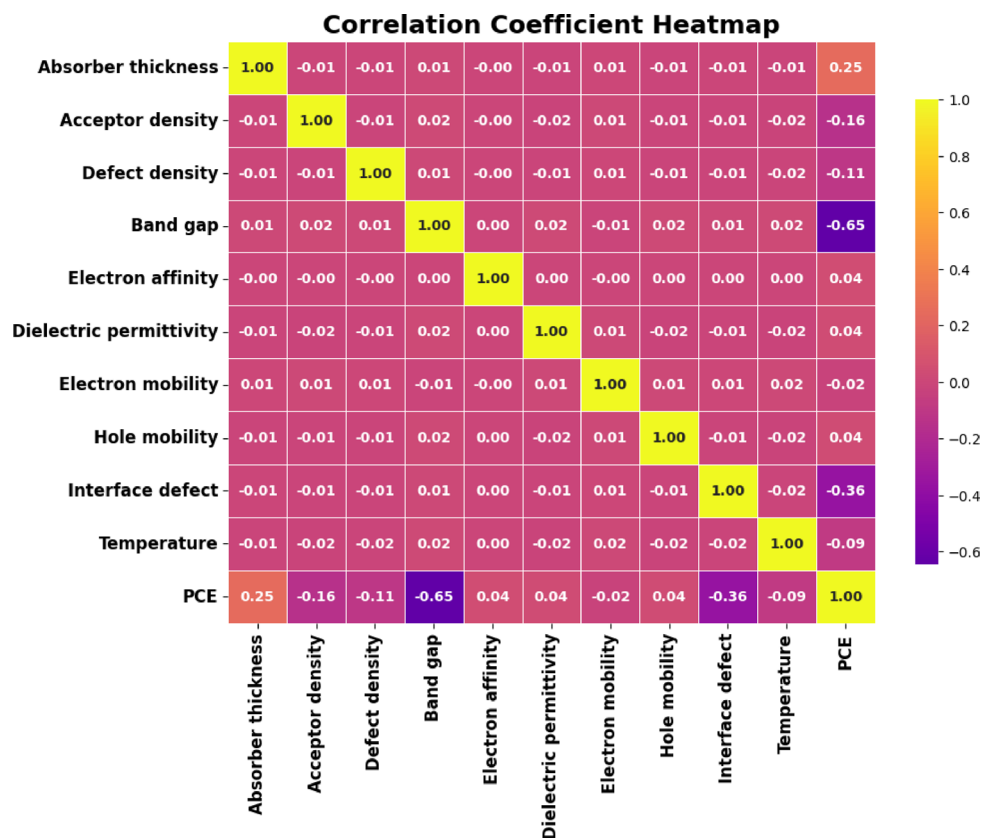
Figure 15 presents the SHAP (Shapley Additive Explanations) summary graphic, which delineates the contribution of each design parameter to PCE as predicted by the Random Forest model. SHAP values on the x-axis measure the impact of each feature positive values denote an enhancement in PCE, whilst negative values suggest a decline in performance. The research finds bandgap energy and interface defect density as the most significant detrimental factors, demonstrating that elevated values of both parameters substantially decrease PCE.

These data align with the correlation heatmap, validating their significant detrimental effect on charge transport and recombination. Defect density exerts a slight negative impact, underscoring the necessity of reducing structural flaws. Conversely, absorber thickness exhibits a significant positive SHAP effect, indicating that an increase in thickness enhances light absorption and device efficiency. Electron mobility has a minor positive contribution, whereas temperature, electron affinity, acceptor density, hole mobility, and dielectric permittivity show minimal impact, with SHAP values around zero. The SHAP analysis identifies bandgap energy and interface defect density as primary factors limiting efficiency, whereas absorber thickness and carrier

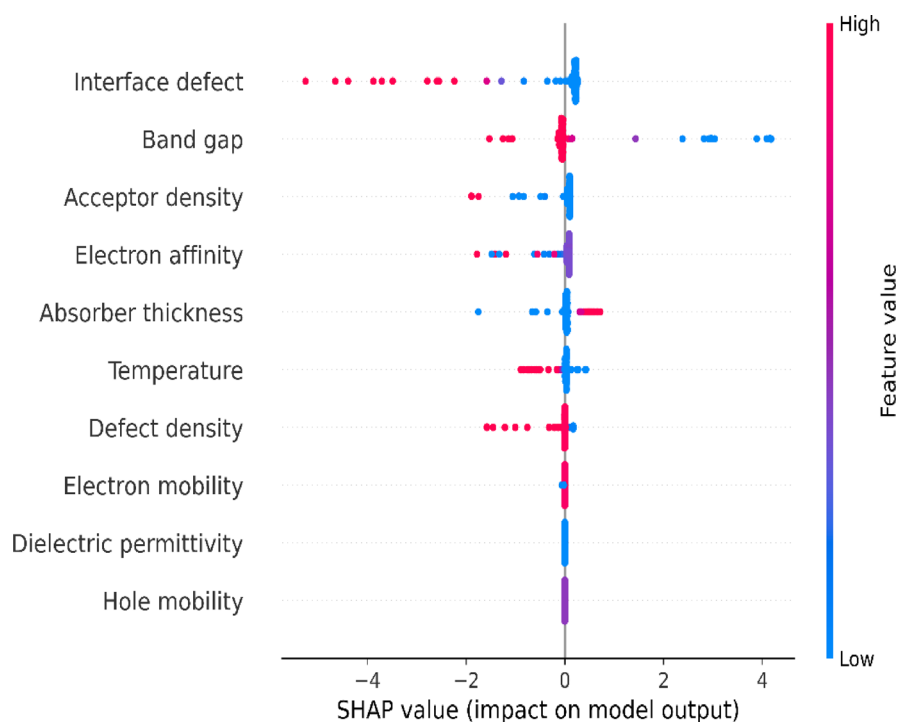
Parameter	MSE	RMSE	MAE	$R^2$	MAPE
$V_{OC}$	0.0012	0.0022	0.0011	0.9894	0.06%
$J_{SC}$	0.0016	0.0245	0.0084	0.9892	0.08%
FF	0.0057	0.0755	0.0224	0.9814	0.02%
PCE	0.0037	0.0606	0.0235	0.9881	0.14%
Average	0.00305	0.0407	0.01385	0.987	0.075

**Table 6.** Evaluating forecast accuracy for key PV parameters through various types of measures.

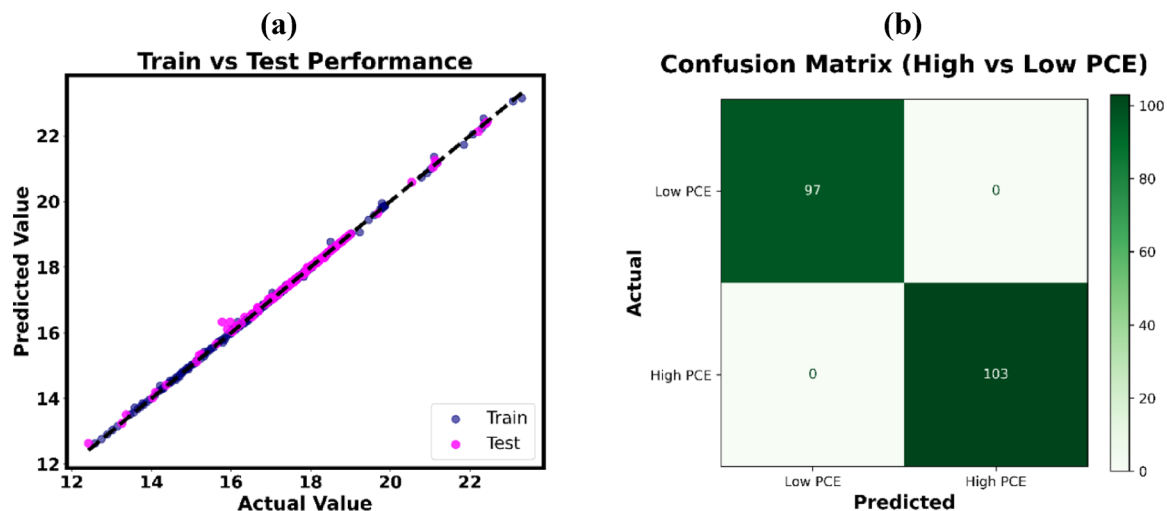




**Fig. 14.** Utilizing a correlation heatmap to illustrate the interactions between factors in the proposed SC configuration.



**Fig. 15.** Thorough examination of variable influences on PCE using SHAP values, with rankings determined by their significance in the rf regression model.



**Fig. 16.** Performance evaluation of the proposed solar cell design as demonstrated by (a) parity plot and (b) confusion matrix.

mobility serve as secondary enhancers. These insights offer essential recommendations for enhancing PSC architecture and attaining superior performance.

#### Evaluating model performance through parity plot and confusion matrix

The predictive efficacy of RF model was assessed by regression and classification studies, yielding a thorough comprehension of its correctness in modeling PSCs. Figure 16a illustrates that the scatter plot predicted against actual values for both the training and testing datasets demonstrates a robust correlation along the diagonal line, signifying exceptional concordance with experimental data. The substantial overlap between training (blue) and testing (magenta) points validates the model's superior prediction accuracy and strong generalization to novel data, effectively reducing overfitting. The results underscore the RF model's capability in identifying nonlinear relationships among essential PV parameters, including bandgap, defect density, and carrier mobility.

Figure 16b presents a confusion matrix that further corroborates the model's classification efficacy in differentiating between high- and low-efficiency PSCs. The matrix demonstrates flawless classification accuracy, comprising 97 true negatives (Low PCE) and 103 genuine positives (High PCE), with no erroneous predictions. This illustrates the model's dependability in accurately categorizing device efficiency levels without error. The regression and classification outcomes collectively validate that the RF model can properly forecast PV performance and effectively distinguish across device quality categories. This dual predictive capability enhances data-driven decision-making in material selection, structural design, and process optimization. The results highlight the revolutionary impact of machine learning in PV research, facilitating the methodical and intelligent design of next-generation PSCs with improved efficiency and stability.

#### Feasibility of device fabrication

To actualize the computationally predicted performance, the fabrication feasibility of  $\text{Be}_3\text{SbF}_3$ -based perovskite solar cells (PSCs) necessitates a meticulously organized experimental framework that includes material synthesis, device assembly, stability assessment, advanced characterization, and strategies for efficiency enhancement. This methodical methodology not only corroborates theoretical conclusions but also establishes a solid basis for enhancing performance in practical scenarios.

##### • Synthesis and preparation of materials

1. *Synthesis of  $\text{Be}_3\text{SbF}_3$* : High-purity  $\text{Be}_3\text{SbF}_3$  can be produced using solid-state reaction, chemical vapor deposition (CVD), or solution-based processes, contingent upon the required film thickness and scalability. Maintaining stoichiometric accuracy and phase consistency is essential, as these elements directly affect charge transfer and optical absorption efficiency.
2. *Structural and morphological characterization*: X-ray diffraction (XRD) will ascertain crystal structure and phase purity, while scanning electron microscopy (SEM) and transmission electron microscopy (TEM) will offer comprehensive insights into surface morphology, grain connectivity, and defect density.
3. *Compositional engineering and doping*: Strategic doping and alloying techniques can be employed to precisely adjust electrical conductivity and enhance stability. Inductively coupled plasma mass spectrometry (ICP-MS) and energy-dispersive X-ray spectroscopy (EDS) guarantee compositional uniformity and identify trace contaminants.

##### • Fabrication of devices

1. *Solar cell construction*: Layer deposition can be accomplished using spin coating, thermal evaporation, or sputtering, resulting in a planar heterojunction device structure, exemplified by Al/FTO/In<sub>2</sub>S<sub>3</sub>/Be<sub>3</sub>SbF<sub>3</sub>/Cu<sub>2</sub>O/Ni. To avert deterioration from oxygen and moisture, manufacture must take place within an inert environment glovebox.
2. *Quality control*: The uniformity of layer production and the accuracy of thickness control can be assessed by atomic force microscopy (AFM) and profilometry. Elevated interfacial uniformity guarantees efficient charge extraction and reduces recombination losses.

#### • Optimization of efficiency using optical engineering

Considering that Be<sub>3</sub>SbF<sub>3</sub> possesses a moderate indirect bandgap, the use of anti-reflective coatings (ARCs) can markedly improve light absorption, especially in the ultraviolet (UV) and visible spectrum regions where reflectivity losses are substantial. Materials like MgF<sub>2</sub>, SiO<sub>2</sub>, Si<sub>3</sub>N<sub>4</sub>, or TiO<sub>2</sub> nanolayers can serve as anti-reflective coatings to diminish surface reflection, enhance photon absorption, and elevate overall PCE<sup>104</sup>. Additionally, textured transparent conducting oxides (TCOs) or graded refractive index coatings may be investigated to prolong optical path lengths and improve light trapping within the absorber layer. The integration of ARCs is particularly advantageous for low-absorption materials, as it efficiently mitigates the constraints of optical thickness without elevating recombination concerns.

#### • Assessment of stability

1. *Environmental stability*: Devices must undergo humidity and oxygen exposure assessments, along with regular J-V measurements to evaluate degradation kinetics. Accelerated aging experiments under constant illumination and elevated temperature circumstances can replicate long-term operating stability.
2. *Thermal and mechanical stability*: Thermal cycling and mechanical bending assessments are crucial for evaluating the durability of flexible or temperature-sensitive setups. These results confirm that Be<sub>3</sub>SbF<sub>3</sub> preserves both structural integrity and efficiency under varying operational settings.

#### • Comparative and benchmarking evaluation

1. *Reference cells*: Performance benchmarks should be established by fabricating reference solar cells utilizing silicon and lead-based perovskites under identical conditions. A comparative assessment of J<sub>SC</sub>, V<sub>OC</sub>, FF, and PCE yields essential insights into the competitiveness of Be<sub>3</sub>SbF<sub>3</sub> absorbers.
2. *Simulation correlation*: Experimental results must be juxtaposed with SCAPS-1D simulation data, emphasizing critical parameters such as carrier lifespan, recombination rate, and defect density. Discrepancies between experimental and simulated data will inform iterative enhancements in both modeling and manufacturing processes.

#### • Sophisticated characterization methods

1. *Interface and defect analysis*: Interfaces and defect states can be examined by X-ray photoelectron spectroscopy (XPS) and Fourier-transform infrared spectroscopy (FTIR) to evaluate chemical bonding and energy-level alignment. Deep-level transient spectroscopy (DLTS) can provide additional quantification of trap states and charge carrier dynamics.
2. *Microscopic and electrical mapping*: Techniques such as Kelvin probe force microscopy (KPFM) and conductive atomic force microscopy (c-AFM) provide the imaging of localized potential changes and charge transport dynamics at the nanoscale, providing enhanced understanding of interfacial quality and electronic uniformity.

#### • Scalability and market potential

Commercializing Be<sub>3</sub>SbF<sub>3</sub>-based solar cells necessitates addressing problems pertaining to stability, scalable production, interface optimization, and cost efficiency. Utilizing encapsulation methods and moisture barriers can enhance environmental resistance. Concurrently, solution-processable large-area deposition techniques, like slot-die coating and blade coating, might enhance industrial scalability. Moreover, including ARCs in large-scale production could significantly enhance module efficiency without modifying absorber chemistry. Alongside economical synthesis and robust interfacial engineering, these advancements establish Be<sub>3</sub>SbF<sub>3</sub> as a feasible lead-free option for next-generation, high-performance optoelectronic and photovoltaic devices. This comprehensive method connects theoretical modeling with practical verification, thereby progressing the development of commercially viable Be<sub>3</sub>SbF<sub>3</sub> perovskite solar systems.

## Conclusion

This study conducts a thorough examination of the structural, optoelectronic, and PV characteristics of Be<sub>3</sub>MF<sub>3</sub> (M = P, As, Sb) perovskites utilizing an integrated methodology that merges DFT with SCAPS-1D simulations. First-principles of DFT computations were utilized to investigate the structural stability, electronic band structure, charge transport, and optical characteristics of these materials. All compounds demonstrated a persistent cubic phase, in accordance with theoretical predictions, and an indirect semiconducting band structure. Optical study demonstrated a wide absorption spectrum and a tight band gap, designating Be<sub>3</sub>SbF<sub>3</sub> as a viable solar absorber with significant light-harvesting capability. Photovoltaic simulations utilizing Be<sub>3</sub>SbF<sub>3</sub> as the absorber across four device designs revealed that the Al/FTO/In<sub>2</sub>S<sub>3</sub>/Be<sub>3</sub>SbF<sub>3</sub>/Cu<sub>2</sub>O/Ni configuration attained exceptional

performance, with  $V_{OC}$  of 1.79 V,  $J_{SC}$  of 11.04 mA/cm<sup>2</sup>, FF of 89.6%, and PCE of 18.28%. The improved performance is due to advantageous band alignment, effective charge extraction, and minimal interfacial recombination. The influence of absorber and ETL thickness, defect density, acceptor density, and temperature was examined, indicating best efficiency at an absorber thickness of 1  $\mu$ m, although increased temperature diminished  $V_{OC}$ , FF, and PCE. Further assessments of QE and J-V characteristics enhanced understanding of device functionality. The RF machine learning model was employed to enhance performance utilizing SCAPS-1D data. connection heatmap and SHAP studies revealed that bandgap energy and interface defect density are significant negative determinants of PCE, but absorber thickness exhibited a positive connection. The parity plot demonstrated exceptional prediction accuracy ( $R^2=0.987$ ), while the confusion matrix attained classification accuracy. These findings collectively establish a data-driven paradigm for the design and optimization of Be<sub>3</sub>MF<sub>3</sub>-based PSCs, providing vital insights for next-generation, high-efficiency PV systems.

## Data availability

Data will be made available on request.

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## Declarations

### Competing interests

The authors declare no competing interests.

## Additional information

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